



FINITE ELEMENT METHODS

B.Tech V semester (Autonomous) IARE R-16

By

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S.No

COURSE OBJECTIVES

I	Introduce basic concepts of finite element methods including domain discretization, polynomial interpolation and application of boundary conditions.
II	Understand the theoretical basics of governing equations and convergence criteria of finite element method.
III	Develop of mathematical model for physical problems and concept of discretization of continuum.
IV	Discuss the accurate Finite Element Solutions for the various field problems
V	Use the commercial Finite Element packages to build Finite Element models and solve a selected range of engineering problems

COs	COURSE OUTCOMES
CO1	Describe the concept of FEM and difference between the FEM with other methods and problems based on 1-D bar elements and shape functions.
CO2	Derive elemental properties and shape functions for truss and beam elements and related problems.
CO3	Understand the concept deriving the elemental matrix and solving the basic problems of CST and axi-symmetric solids
CO4	Explore the concept of steady state heat transfer in fin and composite slab
CO5	Understand the concept of consistent and lumped mass models and solve the dynamic analysis of all types of elements.

UNIT I

INTRODUCTION

CLOs	COURSE LEARNING OUTCOMES
CLO1	Describe the basic concepts of FEM and steps involved in it.
CLO2	Understand the difference between the FEM and Other methods.
CLO3	Understand the stress-strain relation for 2-D and their field problem.
CLO4	Understand the concepts of shape functions for one dimensional and quadratic elements, stiffness matrix and boundary conditions
CLO5	Apply numerical methods for solving one dimensional bar problems

Introduction to FEM:

- ⦿ Stiffness equations for a axial bar element in local co-ordinates using Potential Energy approach and Virtual energy principle.
- ⦿ Finite element analysis of uniform, stepped and tapered bars subjected to mechanical and thermal loads.
- ⦿ Assembly of Global stiffness matrix and load vector.
- ⦿ Quadratic shape functions.
- ⦿ Properties of stiffness matrix

Axially Loaded Bar – Governing Equations and Boundary Conditions



- Differential Equation

$$\frac{d}{dx} \left[EA(x) \frac{du}{dx} \right] + f(x) = 0 \quad 0 < x < L$$

- Boundary Condition Types

- Prescribed displacement (essential BC)
- Prescribed force/derivative of displacement (natural BC)

Axially Loaded Bar –Boundary Conditions



Examples

☐ Fixed end

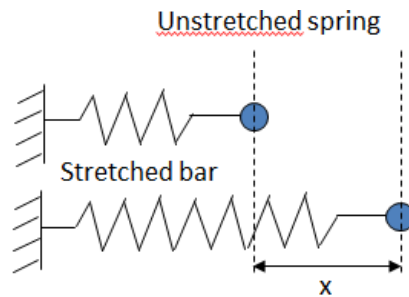
☐ Simple support

☐ Free end

Potential Energy

Elastic Potential Energy (PE)

Spring case



$$PE = 0$$

$$PE = \frac{1}{2} kx^2$$

Axially loaded bar

Unreformed: $PE = 0$

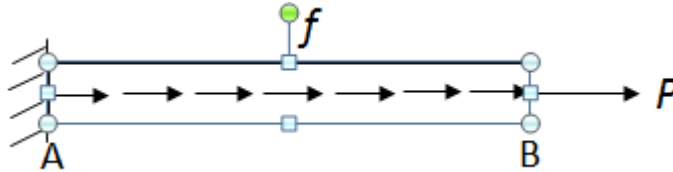
deformed: $PE = \frac{1}{2} \int_0^L \sigma \epsilon A dx$

Elastic body

$$PE = \frac{1}{2} \int_V \boldsymbol{\sigma}^T \boldsymbol{\epsilon} dv$$

Potential Energy

Work Potential (WE)



f : distributed force over a line

P : point force

u : displacement

Total Potential Energy

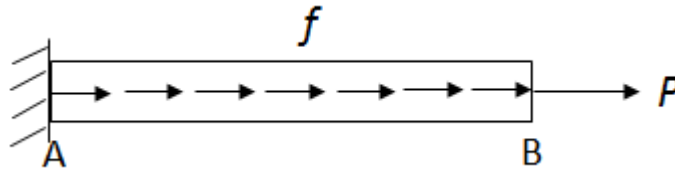
$$\Pi = \frac{1}{2} \int_0^L \sigma \varepsilon A dx - \int_0^L u \cdot f dx - P \cdot u_B$$

Principle of Minimum Potential Energy

For conservative systems, of all the kinematically admissible displacement fields, those corresponding to equilibrium extremize the total potential energy. If the extremum condition is a minimum, the equilibrium state is stable.

Potential Energy + Rayleigh-Ritz Approach

Example:



Step 1: assume a displacement field

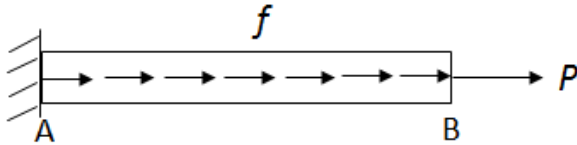
$$u = \sum_i a_i \phi_i(x) \quad i = 1 \text{ to } n$$

- ϕ is shape function / basis function
- n is the order of approximation

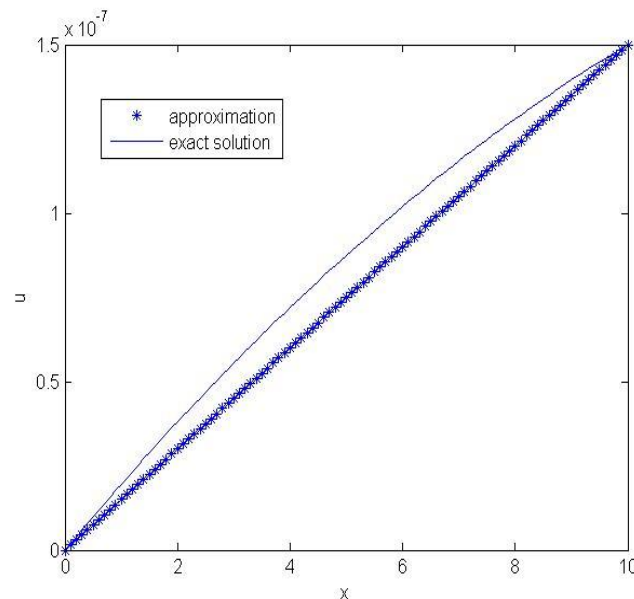
Step 2: calculate total potential energy

Potential Energy + Rayleigh-Ritz Approach

Example:

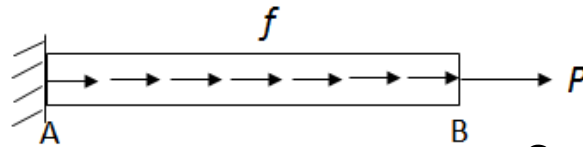


Step 3: select a_i so that the total potential energy is minimum



Galerkin's Method

Example:



Seek an approximation so

$$\begin{cases} \frac{d}{dx} \left[EA(x) \frac{du}{dx} \right] + f(x) = 0 \\ u(x=0) = 0 \\ EA(x) \frac{du}{dx} \Big|_{x=L} = P \end{cases}$$

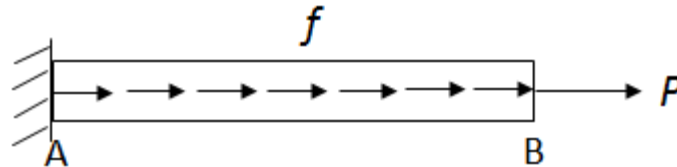


$$\begin{cases} \int_V w_i \left(\frac{d}{dx} \left[EA(x) \frac{d\tilde{u}}{dx} \right] + f(x) \right) dV = 0 \\ \tilde{u}(x=0) = 0 \\ EA(x) \frac{d\tilde{u}}{dx} \Big|_{x=L} = P \end{cases}$$

- In the Galerkin's method, the weight function is chosen to be the same as the shape function.

Galerkin's Method

Example:



$$\int_V w_i \left(\frac{d}{dx} \left[EA(x) \frac{d\tilde{u}}{dx} \right] + f(x) \right) dV = 0 \quad \Rightarrow \quad \overset{\textcircled{1}}{-\int_0^L EA(x) \frac{d\tilde{u}}{dx} \frac{dw_i}{dx} dx} + \overset{\textcircled{2}}{\int_0^L w_i f dx} + \overset{\textcircled{3}}{w_i EA(x) \frac{d\tilde{u}}{dx} \Big|_0^L} = 0$$

⦿ Differential Equation

$$\frac{d}{dx} \left[EA(x) \frac{du}{dx} \right] + f(x) = 0 \quad 0 < x < L$$



⦿ Weighted-Integral Formulation

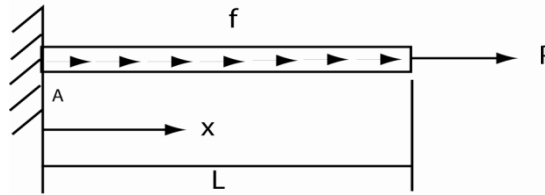
$$\int_0^L w \left(\frac{d}{dx} \left[EA(x) \frac{du}{dx} \right] + f(x) \right) dx = 0$$



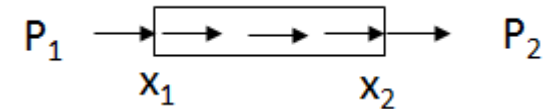
⦿ Weak Form

$$0 = \int_0^L \left[\frac{dw}{dx} \left(EA(x) \frac{du}{dx} \right) - w f(x) \right] dx - w \left(EA(x) \frac{du}{dx} \right) \Big|_0^L$$

Example:



Step 1: Discretization



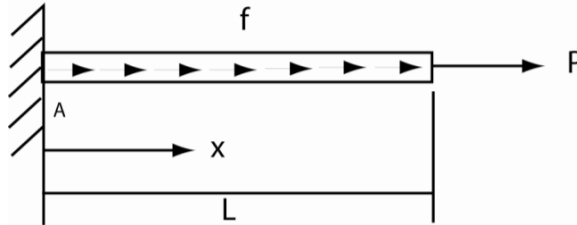
Step 2: Weak form of one element

$$\int_{x_1}^{x_2} \left[\frac{dw}{dx} \left(EA(x) \frac{du}{dx} \right) - w(x) f(x) \right] dx - w(x) \left(EA(x) \frac{du}{dx} \right) \Big|_{x_1}^{x_2} = 0$$

$$\int_{x_1}^{x_2} \left[\frac{dw}{dx} \left(EA(x) \frac{du}{dx} \right) - w(x) f(x) \right] dx - w(x_2) P_2 - w(x_1) P_1 = 0$$

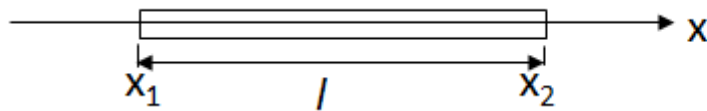
Approximation Methods – Finite Element Method

- Example (*cont*):

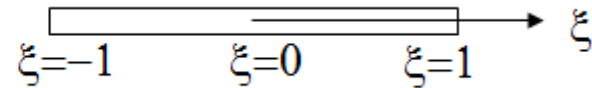


- Step 3: Choosing shape functions- linear shape functions

$$u = \phi_1 u_1 + \phi_2 u_2$$



$$\phi_1 = \frac{x_2 - x}{l}; \quad \phi_2 = \frac{x - x_1}{l}$$

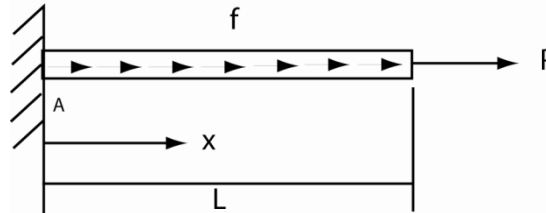


$$\phi_1 = \frac{1 - \xi}{2}; \quad \phi_2 = \frac{1 + \xi}{2}$$

$$\xi = \frac{2}{l}(x - x_1) - 1; \quad x = \frac{(\xi + 1)l}{2} + x_1$$

Approximation Methods – Finite Element Method

Example (cont):

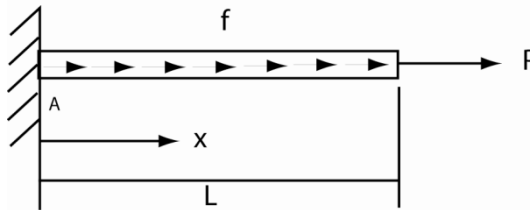


Step 4: Forming element equation

- Let $w = \phi_1$ weak form becomes $\int_{x_1}^{x_2} -\frac{1}{l} \left(EA \cdot \frac{u_2 - u_1}{l} \right) dx - \int_{x_1}^{x_2} \phi_1 f dx - \phi_1 P_2 - \phi_1 P_1 = 0 \longrightarrow \frac{EA}{l} u_1 - \frac{EA}{l} u_2 = \int_{x_1}^{x_2} \phi_1 f dx + P_1$
- Let $w = \phi_2$ weak form becomes $\int_{x_1}^{x_2} \frac{1}{l} \left(EA \cdot \frac{u_2 - u_1}{l} \right) dx - \int_{x_1}^{x_2} \phi_2 f dx - \phi_2 P_2 - \phi_2 P_1 = 0 \longrightarrow -\frac{EA}{l} u_1 + \frac{EA}{l} u_2 = \int_{x_1}^{x_2} \phi_2 f dx + P_2$

$$\longrightarrow \frac{EA}{l} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix} = \begin{Bmatrix} \int_{x_1}^{x_2} \phi_1 f dx \\ \int_{x_1}^{x_2} \phi_2 f dx \end{Bmatrix} + \begin{Bmatrix} P_1 \\ P_2 \end{Bmatrix} = \begin{Bmatrix} f_1 \\ f_2 \end{Bmatrix} + \begin{Bmatrix} P_1 \\ P_2 \end{Bmatrix}$$

Approximation Methods – Finite Element Method



- Example (*cont*):

- Step 5: Assembling to form system equation

Approach 1:

Element 1:

$$\frac{E^I A^I}{l^I} \begin{bmatrix} 1 & -1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} u_1^I \\ u_2^I \\ 0 \\ 0 \end{Bmatrix} = \begin{Bmatrix} f_1^I \\ f_2^I \\ 0 \\ 0 \end{Bmatrix} + \begin{Bmatrix} P_1^I \\ P_2^I \\ 0 \\ 0 \end{Bmatrix}$$

Element 2:

$$\frac{E^{II} A^{II}}{l^{II}} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & -1 & 0 \\ 0 & -1 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} 0 \\ u_1^{II} \\ u_2^{II} \\ 0 \end{Bmatrix} = \begin{Bmatrix} 0 \\ f_1^{II} \\ f_2^{II} \\ 0 \end{Bmatrix} + \begin{Bmatrix} 0 \\ P_1^{II} \\ P_2^{II} \\ 0 \end{Bmatrix}$$

Element 3:

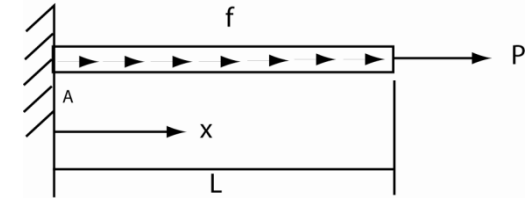
$$\frac{E^{III} A^{III}}{l^{III}} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ 0 & 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} 0 \\ 0 \\ u_1^{III} \\ u_2^{III} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ f_1^{III} \\ f_2^{III} \end{Bmatrix} + \begin{Bmatrix} 0 \\ 0 \\ P_1^{III} \\ P_2^{III} \end{Bmatrix}$$



Example (*cont*):

Step 5: Assembling to form system equation

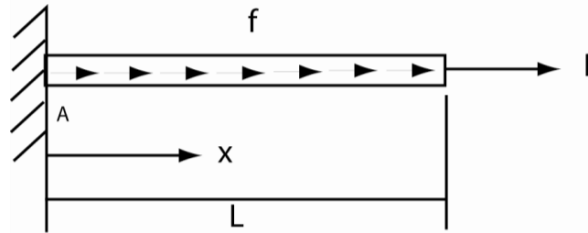
Assembled System:



$$\begin{bmatrix} \frac{E^I A^I}{l^I} & -\frac{E^I A^I}{l^I} & 0 & 0 \\ -\frac{E^I A^I}{l^I} & \frac{E^I A^I}{l^I} + \frac{E^{II} A^{II}}{l^{II}} & -\frac{E^{II} A^{II}}{l^{II}} & 0 \\ 0 & -\frac{E^{II} A^{II}}{l^{II}} & \frac{E^{II} A^{II}}{l^{II}} + \frac{E^{III} A^{III}}{l^{III}} & -\frac{E^{III} A^{III}}{l^{III}} \\ 0 & 0 & -\frac{E^{III} A^{III}}{l^{III}} & \frac{E^{III} A^{III}}{l^{III}} \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{Bmatrix} = \begin{Bmatrix} f_1 \\ f_2 \\ f_3 \\ f_4 \end{Bmatrix} + \begin{Bmatrix} P_1 \\ P_2 \\ P_3 \\ P_4 \end{Bmatrix} = \begin{Bmatrix} f_1^I \\ f_2^I + f_1^{II} \\ f_2^{II} + f_1^{III} \\ f_2^{III} \end{Bmatrix} + \begin{Bmatrix} P_1^I \\ P_2^I + P_1^{II} \\ P_2^{II} + P_1^{III} \\ P_2^{III} \end{Bmatrix}$$

Approximation Methods

Example (cont):



Step 5: Assembling to form system equation

Approach 2: Element connectivity table

$$k_{ij}^e \rightarrow K_{IJ}$$

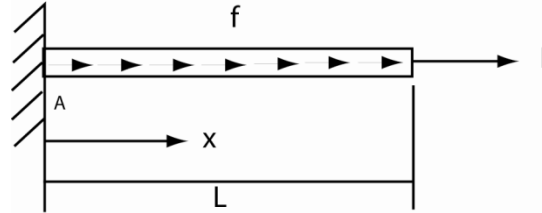
	Element 1	Element 2	Element 3
1	1	2	3
2	2	3	4

↓
local node
(i,j)

global node index
(I,J)

Approximation Methods

Example (cont):



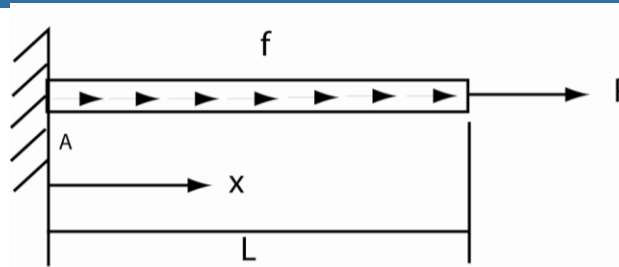
Step 6: Imposing boundary conditions and forming condense

Condensed system:

$$\begin{pmatrix} \frac{E^I A^I}{l^I} + \frac{E^{II} A^{II}}{l^{II}} & -\frac{E^{II} A^{II}}{l^{II}} & 0 \\ -\frac{E^{II} A^{II}}{l^{II}} & \frac{E^{II} A^{II}}{l^{II}} + \frac{E^{III} A^{III}}{l^{III}} & -\frac{E^{III} A^{III}}{l^{III}} \\ 0 & -\frac{E^{III} A^{III}}{l^{III}} & \frac{E^{III} A^{III}}{l^{III}} \end{pmatrix} \begin{Bmatrix} u_2 \\ u_3 \\ u_4 \end{Bmatrix} = \begin{Bmatrix} f_2 \\ f_3 \\ f_4 \end{Bmatrix} + \begin{Bmatrix} 0 \\ 0 \\ P \end{Bmatrix}$$

Approximation Methods

- Example (*cont*):



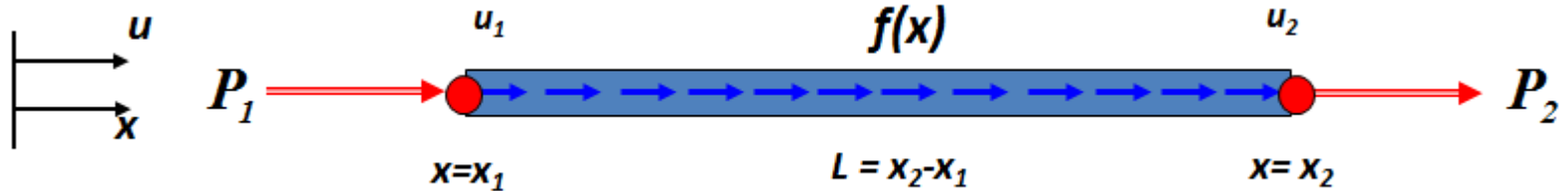
- Step 7: solution
- Step 8: post calculation

$$u = u_1\phi_1 + u_2\phi_2 \longrightarrow \varepsilon = \frac{du}{dx} = u_1 \frac{d\phi_1}{dx} + u_2 \frac{d\phi_2}{dx} \longrightarrow \sigma = E\varepsilon = Eu_1 \frac{d\phi_1}{dx} + Eu_2 \frac{d\phi_2}{dx}$$

Summary - Major Steps in FEM

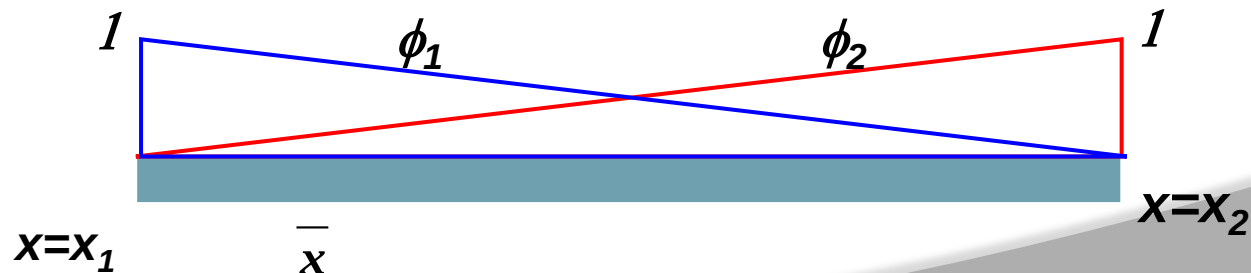
- ⦿ Discretization
- ⦿ Derivation of element equation
 - weak form
 - construct form of approximation solution over one element
 - derive finite element model
- ⦿ Assembling – putting elements together
- ⦿ Imposing boundary conditions
- ⦿ Solving equations
- ⦿ Post computation

Linear Formulation for Bar Element

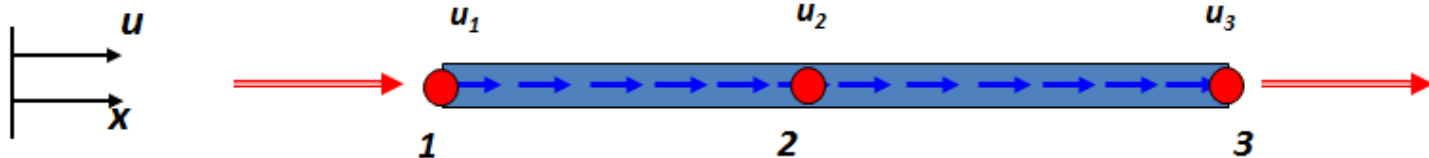


$$\begin{Bmatrix} P_1 \\ P_2 \end{Bmatrix} + \begin{Bmatrix} f_1 \\ f_2 \end{Bmatrix} = \begin{bmatrix} K_{11} & K_{12} \\ K_{12} & K_{22} \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix}$$

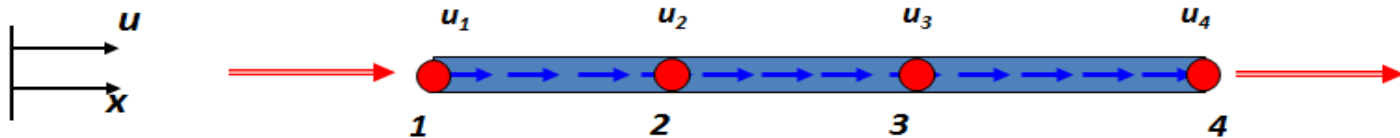
$$\text{where } K_{ij} = \int_{x_1}^{x_2} EA \left(\frac{d\phi_i}{dx} \frac{d\phi_j}{dx} \right) dx = K_{ji}, \quad f_i = \int_{x_1}^{x_2} (\phi_i f) dx$$



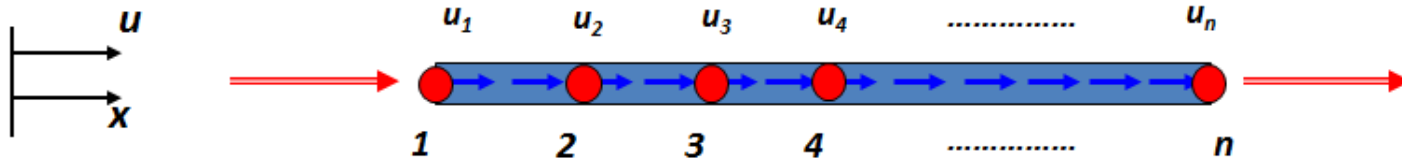
Higher Order Formulation for Bar Element



$$u(x) = u_1\phi_1(x) + u_2\phi_2(x) + u_3\phi_3(x)$$

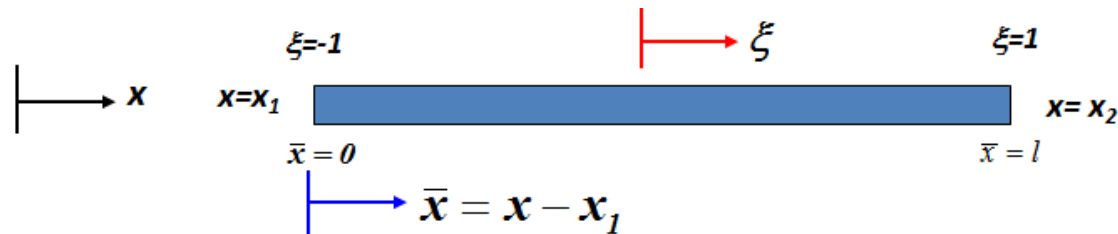


$$u(x) = u_1\phi_1(x) + u_2\phi_2(x) + u_3\phi_3(x) + u_4\phi_4(x)$$

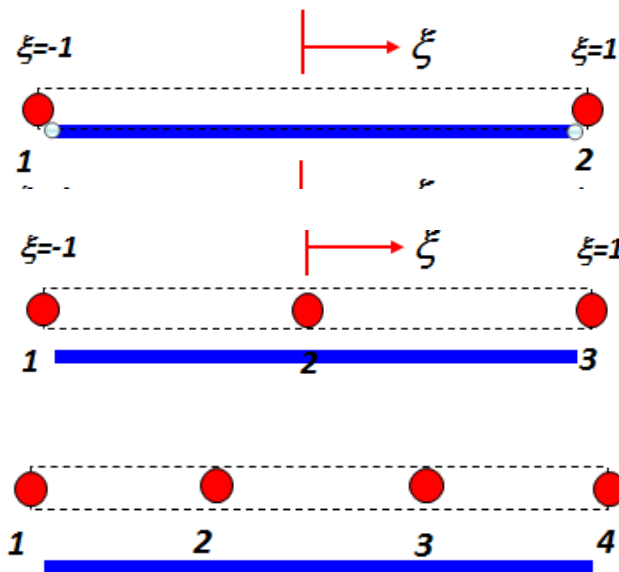


$$u(x) = u_1\phi_1(x) + u_2\phi_2(x) + u_3\phi_3(x) + u_4\phi_4(x) + \dots + u_n\phi_n(x)$$

Natural Coordinates and Interpolation Functions



Natural (or Normal) Coordinate: $\xi = \frac{x - \frac{x_1 + x_2}{2}}{l/2}$



$$\phi_1 = -\frac{\xi - 1}{2}, \quad \phi_2 = \frac{\xi + 1}{2}$$

$$\phi_1 = \frac{\xi(\xi - 1)}{2}, \quad \phi_2 = -(\xi + 1)(\xi - 1), \quad \phi_3 = \frac{(\xi + 1)\xi}{2}$$

$$\phi_1 = -\frac{9}{16} \left(\xi + \frac{1}{3} \right) \left(\xi - \frac{1}{3} \right) (\xi - 1), \quad \phi_2 = \frac{27}{16} (\xi + 1) \left(\xi - \frac{1}{3} \right) (\xi - 1)$$

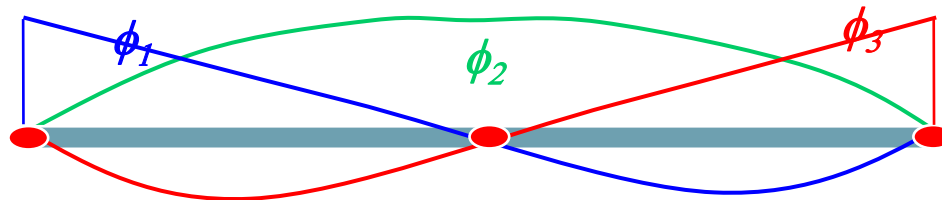
$$\phi_3 = -\frac{27}{16} (\xi + 1) \left(\xi + \frac{1}{3} \right) (\xi - 1), \quad \phi_4 = \frac{9}{16} (\xi + 1) \left(\xi + \frac{1}{3} \right) \left(\xi - \frac{1}{3} \right)$$

Quadratic Formulation for Bar Element

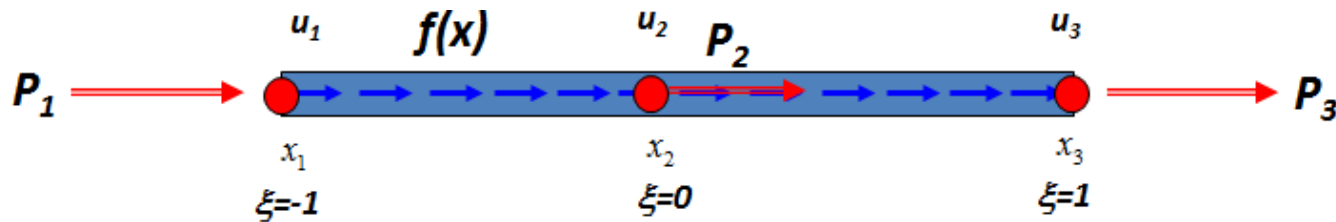
$$\begin{Bmatrix} P_1 \\ P_2 \\ P_3 \end{Bmatrix} + \begin{Bmatrix} f_1 \\ f_2 \\ f_3 \end{Bmatrix} = \begin{bmatrix} K_{11} & K_{12} & K_{13} \\ K_{12} & K_{22} & K_{23} \\ K_{13} & K_{23} & K_{33} \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix}$$

where $K_{ij} = \int_{x_1}^{x_2} EA \left(\frac{d\phi_i}{dx} \frac{d\phi_j}{dx} \right) dx = \int_{-1}^1 EA \left(\frac{d\phi_i}{d\xi} \frac{d\phi_j}{d\xi} \right) \frac{2}{l} d\xi = K_{ji}$

and $f_i = \int_{x_1}^{x_2} (\phi_i f) dx = \int_{-1}^1 (\phi_i f) \frac{l}{2} d\xi, \quad i, j = 1, 2, 3$



Quadratic Formulation for Bar Element



$$u(\xi) = u_1 \phi_1(\xi) + u_2 \phi_2(\xi) + u_3 \phi_3(\xi) = u_1 \frac{\xi(\xi-1)}{2} - u_2 (\xi+1)(\xi-1) + u_3 \frac{(\xi+1)\xi}{2}$$

$$\phi_1 = \frac{\xi(\xi-1)}{2}, \quad \phi_2 = -(\xi+1)(\xi-1), \quad \phi_3 = \frac{(\xi+1)\xi}{2}$$

$$\xi = \frac{x - \frac{x_1 + x_2}{2}}{l/2}$$

$$\frac{l}{2} d\xi = dx$$

$$\frac{d\xi}{dx} = \frac{2}{l}$$

$$\frac{d\phi_1}{dx} = \frac{2}{l} \frac{d\phi_1}{d\xi} = \frac{2\xi-1}{l}, \quad \frac{d\phi_2}{dx} = \frac{2}{l} \frac{d\phi_2}{d\xi} = -\frac{4\xi}{l}, \quad \frac{d\phi_3}{dx} = \frac{2}{l} \frac{d\phi_3}{d\xi} = \frac{2\xi+1}{l}$$

UNIT-2

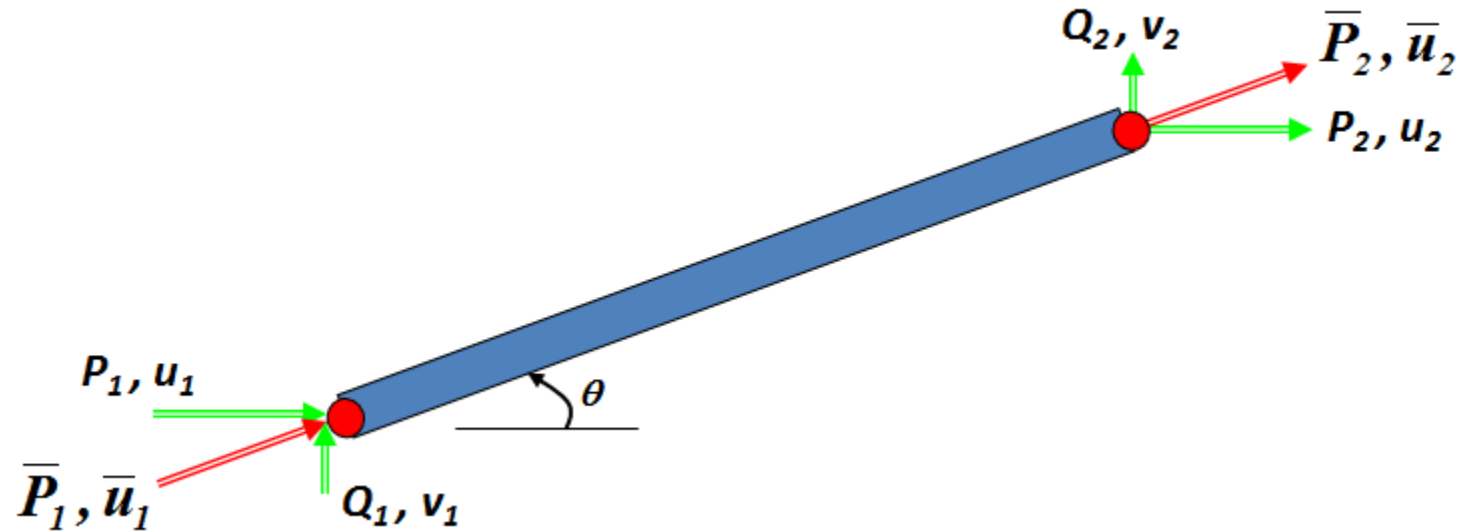
ANALYSIS OF TRUSSES AND BEAMS

CLOs	Course Learning Outcomes
CLO 1	Derive the elemental property matrix for beam and bar elements.
CLO 2	Solve the equations of truss and beam elements
CLO 3	Understand the concepts of shape functions for beam element.
CLO 4	Apply the numerical methods for solving truss and beam problems

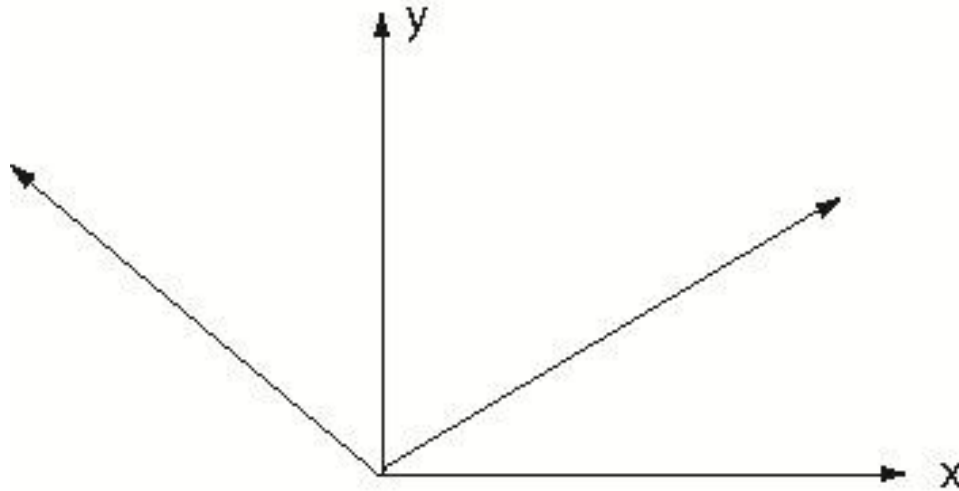
Finite Element Analysis of Trusses:

- ① Stiffness equations for a truss bar element oriented in 2D plane
- ① Finite Element Analysis of Trusses
- ① Plane Truss and Space Truss elements
- ① Methods of assembly

Arbitrarily Oriented 1-D Bar Element on 2-D Plane



Relationship Between Local Coordinates and Global Coordinates

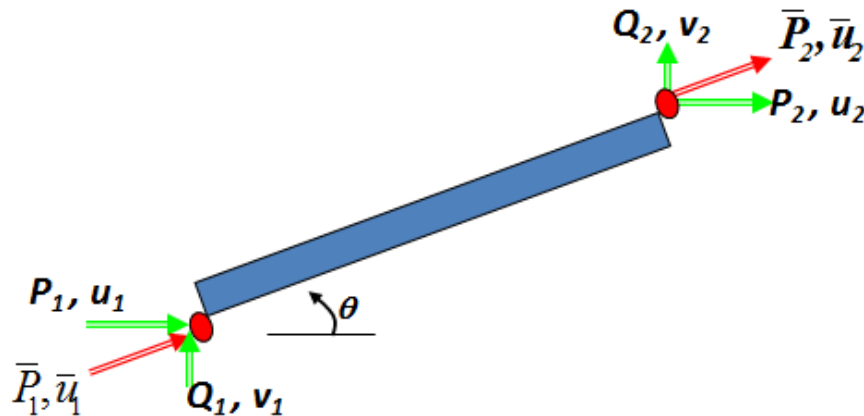


$$\begin{Bmatrix} \bar{u}_1 \\ \bar{v}_1 = 0 \\ \bar{u}_2 \\ \bar{v}_2 = 0 \end{Bmatrix} = \begin{bmatrix} \cos\theta & \sin\theta & 0 & 0 \\ -\sin\theta & \cos\theta & 0 & 0 \\ \hline 0 & 0 & \cos\theta & \sin\theta \\ 0 & 0 & -\sin\theta & \cos\theta \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \end{Bmatrix}$$

Relationship Between Local Coordinates and Global Coordinates

$$\begin{Bmatrix} \bar{P}_1 \\ 0 \\ \bar{P}_2 \\ 0 \end{Bmatrix} = \begin{bmatrix} \cos\theta & \sin\theta & 0 & 0 \\ -\sin\theta & \cos\theta & 0 & 0 \\ 0 & 0 & \cos\theta & \sin\theta \\ 0 & 0 & -\sin\theta & \cos\theta \end{bmatrix} \begin{Bmatrix} P_1 \\ Q_1 \\ P_2 \\ Q_2 \end{Bmatrix}$$

Stiffness Matrix of 1-D Bar Element on 2-D Plane

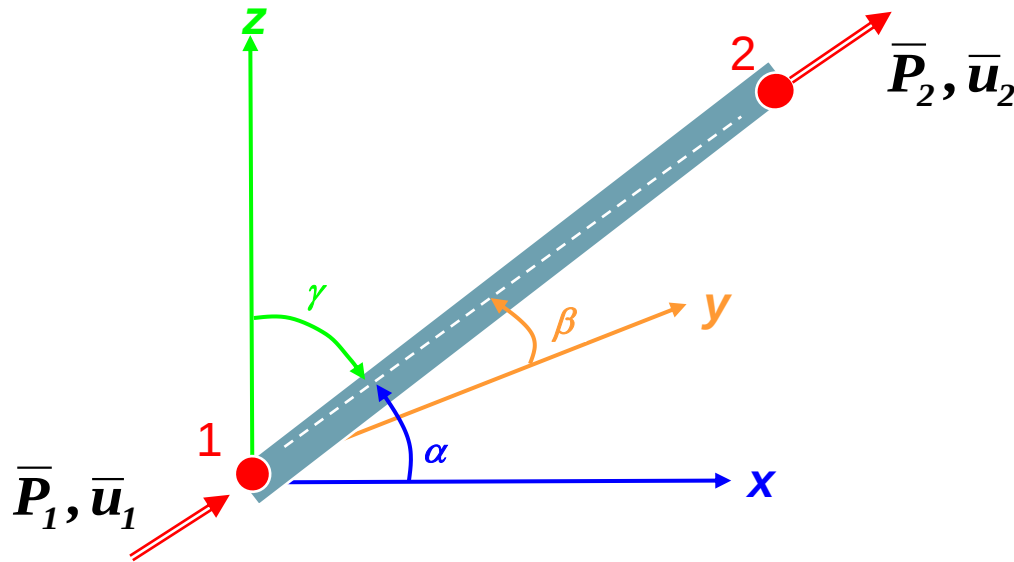


$$[\bar{K}_{ij}] = \frac{AE}{L} \begin{bmatrix} 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_2 \\ Q_2 \end{Bmatrix} = \begin{bmatrix} \cos \theta & -\sin \theta & 0 & 0 \\ \sin \theta & \cos \theta & 0 & 0 \\ 0 & 0 & \cos \theta & -\sin \theta \\ 0 & 0 & \sin \theta & \cos \theta \end{bmatrix} [\bar{K}_{ij}] \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \end{Bmatrix}$$

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_2 \\ Q_2 \end{Bmatrix} = \frac{AE}{L} \begin{bmatrix} \cos^2 \theta & \sin \theta \cos \theta & -\cos^2 \theta & -\sin \theta \cos \theta \\ \sin \theta \cos \theta & \sin^2 \theta & -\sin \theta \cos \theta & -\sin^2 \theta \\ -\cos^2 \theta & -\sin \theta \cos \theta & \cos^2 \theta & \sin \theta \cos \theta \\ -\sin \theta \cos \theta & -\sin^2 \theta & \sin \theta \cos \theta & \sin^2 \theta \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \end{Bmatrix}$$

Arbitrarily Oriented 1-D Bar Element in 3-D Space

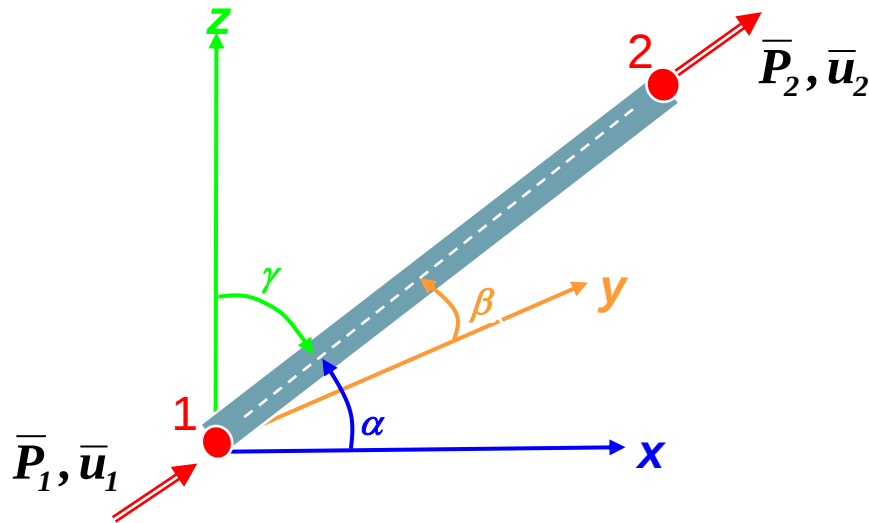


α , β , γ are the
Direction Cosines of the
bar in the x-y-z
coordinate system

$$\begin{Bmatrix} \bar{u}_1 \\ \bar{v}_1 = 0 \\ \bar{w}_1 = 0 \\ \bar{u}_2 \\ \bar{v}_2 = 0 \\ \bar{w}_2 = 0 \end{Bmatrix} = \begin{bmatrix} \alpha_{\bar{x}} & \beta_{\bar{x}} & \gamma_{\bar{x}} & 0 & 0 & 0 \\ \alpha_{\bar{y}} & \beta_{\bar{y}} & \gamma_{\bar{y}} & 0 & 0 & 0 \\ \alpha_{\bar{z}} & \beta_{\bar{z}} & \gamma_{\bar{z}} & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & \alpha_{\bar{x}} & \beta_{\bar{x}} & \gamma_{\bar{x}} \\ 0 & 0 & 0 & \alpha_{\bar{y}} & \beta_{\bar{y}} & \gamma_{\bar{y}} \\ 0 & 0 & 0 & \alpha_{\bar{z}} & \beta_{\bar{z}} & \gamma_{\bar{z}} \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ w_1 \\ u_2 \\ v_2 \\ w_2 \end{Bmatrix}$$

$$\begin{Bmatrix} \bar{P}_1 \\ \bar{Q}_1 = 0 \\ \bar{R}_1 = 0 \\ \bar{P}_2 \\ \bar{Q}_2 = 0 \\ \bar{R}_2 = 0 \end{Bmatrix} = \begin{bmatrix} \alpha_{\bar{x}} & \beta_{\bar{x}} & \gamma_{\bar{x}} & 0 & 0 & 0 \\ \alpha_{\bar{y}} & \beta_{\bar{y}} & \gamma_{\bar{y}} & 0 & 0 & 0 \\ \alpha_{\bar{z}} & \beta_{\bar{z}} & \gamma_{\bar{z}} & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & \alpha_{\bar{x}} & \beta_{\bar{x}} & \gamma_{\bar{x}} \\ 0 & 0 & 0 & \alpha_{\bar{y}} & \beta_{\bar{y}} & \gamma_{\bar{y}} \\ 0 & 0 & 0 & \alpha_{\bar{z}} & \beta_{\bar{z}} & \gamma_{\bar{z}} \end{bmatrix} \begin{Bmatrix} P_1 \\ Q_1 \\ R_1 \\ P_2 \\ Q_2 \\ R_2 \end{Bmatrix}$$

Stiffness Matrix of 1-D Bar Element in 3-D Space



$$\begin{Bmatrix} \bar{P}_1 \\ \bar{Q}_1 = 0 \\ \bar{R}_1 = 0 \\ \bar{P}_2 \\ \bar{Q}_2 = 0 \\ \bar{R}_2 = 0 \end{Bmatrix} = \frac{AE}{L} \begin{bmatrix} 1 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} \bar{u}_1 \\ \bar{v}_1 = 0 \\ \bar{w}_1 = 0 \\ \bar{u}_2 \\ \bar{v}_2 = 0 \\ \bar{w}_2 = 0 \end{Bmatrix}$$

$$\begin{Bmatrix} P_1 \\ Q_1 \\ R_1 \\ P_2 \\ Q_2 \\ R_2 \end{Bmatrix} = \frac{AE}{L} \begin{bmatrix} \alpha_{\bar{x}}^2 & \alpha_{\bar{x}}\beta_{\bar{x}} & \alpha_{\bar{x}}\gamma_{\bar{x}} & -\alpha_{\bar{x}}^2 & -\alpha_{\bar{x}}\beta_{\bar{x}} & -\alpha_{\bar{x}}\gamma_{\bar{x}} \\ \alpha_{\bar{x}}\beta_{\bar{x}} & \beta_{\bar{x}}^2 & \beta_{\bar{x}}\gamma_{\bar{x}} & -\alpha_{\bar{x}}\beta_{\bar{x}} & -\beta_{\bar{x}}^2 & -\beta_{\bar{x}}\gamma_{\bar{x}} \\ \alpha_{\bar{x}}\gamma_{\bar{x}} & \beta_{\bar{x}}\gamma_{\bar{x}} & \gamma_{\bar{x}}^2 & -\alpha_{\bar{x}}\gamma_{\bar{x}} & -\beta_{\bar{x}}\gamma_{\bar{x}} & -\gamma_{\bar{x}}^2 \\ -\alpha_{\bar{x}}^2 & -\alpha_{\bar{x}}\beta_{\bar{x}} & -\alpha_{\bar{x}}\gamma_{\bar{x}} & \alpha_{\bar{x}}^2 & \alpha_{\bar{x}}\beta_{\bar{x}} & \alpha_{\bar{x}}\gamma_{\bar{x}} \\ -\alpha_{\bar{x}}\beta_{\bar{x}} & -\beta_{\bar{x}}^2 & -\beta_{\bar{x}}\gamma_{\bar{x}} & \alpha_{\bar{x}}\beta_{\bar{x}} & \beta_{\bar{x}}^2 & \beta_{\bar{x}}\gamma_{\bar{x}} \\ -\alpha_{\bar{x}}\gamma_{\bar{x}} & -\beta_{\bar{x}}\gamma_{\bar{x}} & -\gamma_{\bar{x}}^2 & \alpha_{\bar{x}}\gamma_{\bar{x}} & \beta_{\bar{x}}\gamma_{\bar{x}} & \gamma_{\bar{x}}^2 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ w_1 \\ u_2 \\ v_2 \\ w_2 \end{Bmatrix}$$

Matrix Assembly of Multiple Bar Elements

Element I

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_2 \\ Q_2 \end{Bmatrix} = \frac{AE}{L} \begin{bmatrix} 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \end{Bmatrix}$$

Element II

$$\begin{Bmatrix} P_2 \\ Q_2 \\ P_3 \\ Q_3 \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 1 & -\sqrt{3} & -1 & \sqrt{3} \\ -\sqrt{3} & 3 & \sqrt{3} & -3 \\ -1 & \sqrt{3} & 1 & -\sqrt{3} \\ \sqrt{3} & -3 & -\sqrt{3} & 3 \end{bmatrix} \begin{Bmatrix} u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}$$

Element III

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_3 \\ Q_3 \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 1 & \sqrt{3} & -1 & -\sqrt{3} \\ \sqrt{3} & 3 & -\sqrt{3} & -3 \\ -1 & -\sqrt{3} & 1 & \sqrt{3} \\ -\sqrt{3} & -3 & \sqrt{3} & 3 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_3 \\ v_3 \end{Bmatrix}$$

Matrix Assembly of Multiple Bar Elements

Element I

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_2 \\ Q_2 \\ P_3 \\ Q_3 \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 4 & 0 & -4 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ -4 & 0 & 4 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}$$

Element II

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_2 \\ Q_2 \\ P_3 \\ Q_3 \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -\sqrt{3} & -1 & \sqrt{3} \\ 0 & 0 & -\sqrt{3} & 3 & \sqrt{3} & -3 \\ 0 & 0 & -1 & \sqrt{3} & 1 & -\sqrt{3} \\ 0 & 0 & \sqrt{3} & -3 & -\sqrt{3} & 3 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}$$

Element III

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_2 \\ Q_2 \\ P_3 \\ Q_3 \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 1 & \sqrt{3} & 0 & 0 & -1 & -\sqrt{3} \\ \sqrt{3} & 3 & 0 & 0 & -\sqrt{3} & -3 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & -\sqrt{3} & 0 & 0 & 1 & \sqrt{3} \\ -\sqrt{3} & -3 & 0 & 0 & \sqrt{3} & 3 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}$$

Matrix Assembly of Multiple Bar Elements

$$\begin{Bmatrix} R_1 \\ S_1 \\ R_2 \\ S_2 \\ R_3 \\ S_3 \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 4+1 & 0+\sqrt{3} & -4 & 0 & -1 & -\sqrt{3} \\ 0+\sqrt{3} & 0+3 & 0 & 0 & -\sqrt{3} & -3 \\ -4 & 0 & 4+1 & 0-\sqrt{3} & -1 & \sqrt{3} \\ 0 & 0 & 0-\sqrt{3} & 0+3 & \sqrt{3} & -3 \\ -1 & -\sqrt{3} & -1 & \sqrt{3} & 1+1 & \sqrt{3}-\sqrt{3} \\ -\sqrt{3} & -3 & \sqrt{3} & -3 & \sqrt{3}-\sqrt{3} & 3+3 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}$$

Apply known boundary conditions

$$\begin{Bmatrix} R_1=? \\ S_1=0 \\ R_2=F \\ S_2=? \\ R_3=? \\ S_3=? \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 5 & \sqrt{3} & -4 & 0 & -1 & -\sqrt{3} \\ \sqrt{3} & 3 & 0 & 0 & -\sqrt{3} & -3 \\ -4 & 0 & 5 & -\sqrt{3} & -1 & \sqrt{3} \\ 0 & 0 & -\sqrt{3} & 3 & \sqrt{3} & -3 \\ -1 & -\sqrt{3} & -1 & \sqrt{3} & 2 & 0 \\ -\sqrt{3} & -3 & \sqrt{3} & -3 & 0 & 6 \end{bmatrix} \begin{Bmatrix} u_1=0 \\ v_1=? \\ u_2=? \\ v_2=0 \\ u_3=0 \\ v_3=0 \end{Bmatrix}$$

Solution Procedures

$$\begin{Bmatrix} R_2 = F \\ S_1 = 0 \\ R_1 = ? \\ S_2 = ? \\ R_3 = ? \\ S_3 = ? \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} -4 & 0 & 5 & -\sqrt{3} & -1 & \sqrt{3} \\ \sqrt{3} & 3 & 0 & 0 & -\sqrt{3} & -3 \\ \hline 5 & \sqrt{3} & -4 & 0 & -1 & -\sqrt{3} \\ 0 & 0 & -\sqrt{3} & 3 & \sqrt{3} & -3 \\ -1 & -\sqrt{3} & -1 & \sqrt{3} & 2 & 0 \\ -\sqrt{3} & -3 & \sqrt{3} & -3 & 0 & 6 \end{bmatrix} \begin{Bmatrix} u_1 = 0 \\ v_1 = ? \\ u_2 = ? \\ v_2 = 0 \\ u_3 = 0 \\ v_3 = 0 \end{Bmatrix}$$

$\Rightarrow u_2 = 4FL/5AE, v_1 = 0$

$$\begin{Bmatrix} R_2 = F \\ S_1 = 0 \\ R_1 = ? \\ S_2 = ? \\ R_3 = ? \\ S_3 = ? \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} -4 & 0 & 5 & -\sqrt{3} & -1 & \sqrt{3} \\ \sqrt{3} & 3 & 0 & 0 & -\sqrt{3} & -3 \\ \hline 5 & \sqrt{3} & -4 & 0 & -1 & -\sqrt{3} \\ 0 & 0 & -\sqrt{3} & 3 & \sqrt{3} & -3 \\ -1 & -\sqrt{3} & -1 & \sqrt{3} & 2 & 0 \\ -\sqrt{3} & -3 & \sqrt{3} & -3 & 0 & 6 \end{bmatrix} \begin{Bmatrix} u_1 = 0 \\ v_1 = 0 \\ u_2 = \frac{4FL}{5AE} \\ v_2 = 0 \\ u_3 = 0 \\ v_3 = 0 \end{Bmatrix}$$

Recovery of axial forces

Element ①

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_2 \\ Q_2 \end{Bmatrix} = \frac{AE}{L} \begin{bmatrix} 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} u_1 = 0 \\ v_1 = 0 \\ u_2 = \frac{4FL}{5AE} \\ v_2 = 0 \end{Bmatrix} = F \begin{Bmatrix} -4/5 \\ 0 \\ 4/5 \\ 0 \end{Bmatrix}$$

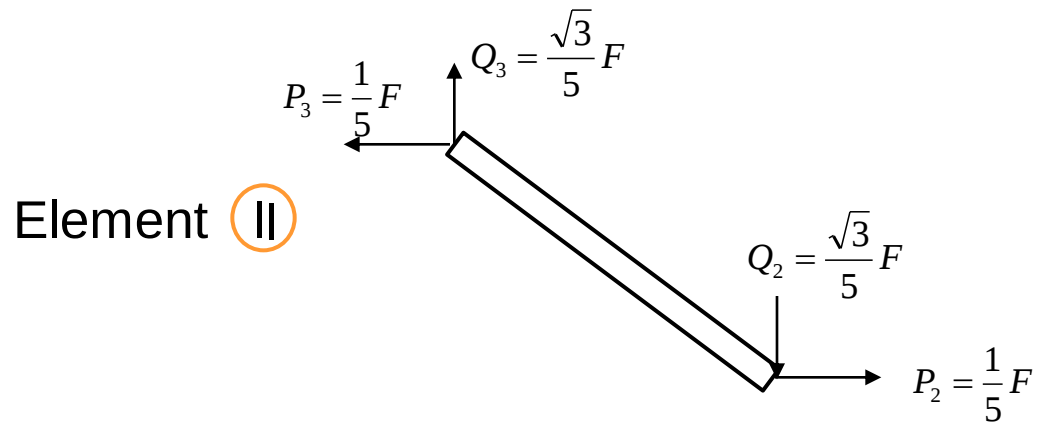
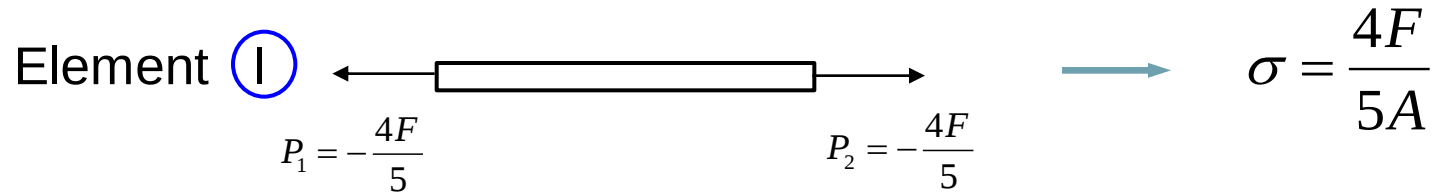
Element ②

$$\begin{Bmatrix} P_2 \\ Q_2 \\ P_3 \\ Q_3 \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 1 & -\sqrt{3} & -1 & \sqrt{3} \\ -\sqrt{3} & 3 & \sqrt{3} & -3 \\ -1 & \sqrt{3} & 1 & -\sqrt{3} \\ \sqrt{3} & -3 & -\sqrt{3} & 3 \end{bmatrix} \begin{Bmatrix} u_2 = \frac{4FL}{5AE} \\ v_2 = 0 \\ u_3 = 0 \\ v_3 = 0 \end{Bmatrix} = F \begin{Bmatrix} 1/5 \\ -\sqrt{3}/5 \\ -1/5 \\ \sqrt{3}/5 \end{Bmatrix}$$

Element ③

$$\begin{Bmatrix} P_1 \\ Q_1 \\ P_3 \\ Q_3 \end{Bmatrix} = \frac{AE}{4L} \begin{bmatrix} 1 & \sqrt{3} & -1 & -\sqrt{3} \\ \sqrt{3} & 3 & -\sqrt{3} & -3 \\ -1 & -\sqrt{3} & 1 & \sqrt{3} \\ -\sqrt{3} & -3 & \sqrt{3} & 3 \end{bmatrix} \begin{Bmatrix} u_1 = 0 \\ v_1 = 0 \\ u_3 = 0 \\ v_3 = 0 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{Bmatrix}$$

Stresses inside members

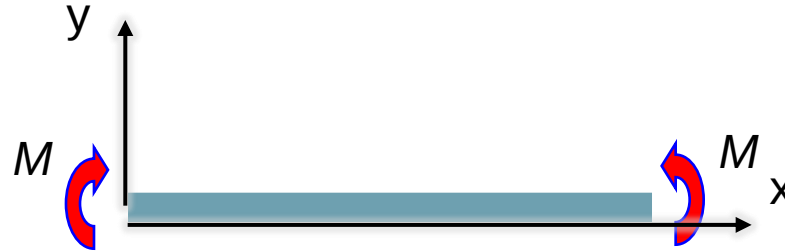


Finite Element Analysis of Beams



- Hermite shape functions
- Element stiffness matrix
- Load vector
- Problems

Review



Pure bending problems

➤ Normal strain: $\epsilon_x = -\frac{y}{\rho}$

➤ Normal stress: $\sigma_x = -\frac{Ey}{\rho}$

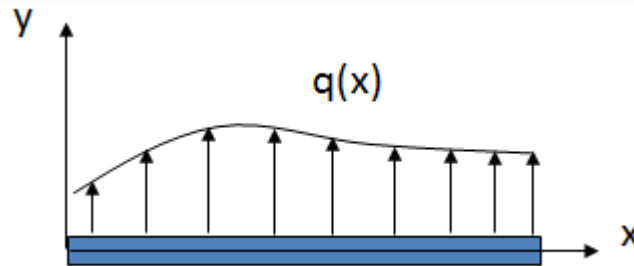
➤ Normal stress with bending moment: $\int -\sigma_x y dA = M$

➤ Moment-curvature relationship: $\frac{1}{\rho} = \frac{M}{EI} \rightarrow M = EI \frac{1}{\rho} \approx EI \frac{d^2 y}{dx^2}$

➤ Flexure formula: $\sigma_x = -\frac{My}{I} \quad I = \int y^2 dA$

Bending Beam

Review



Relationship between shear force, bending moment and transverse load:

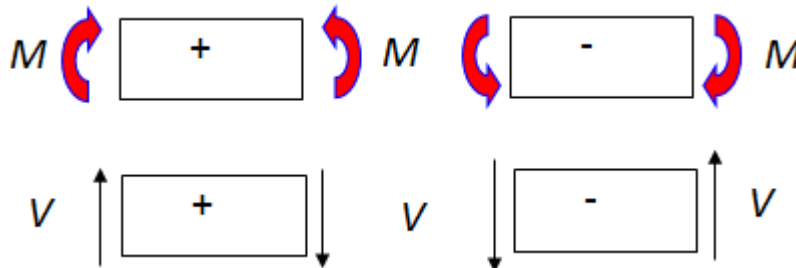
Deflection

$$\frac{dV}{dx} = q$$

$$\frac{dM}{dx} = V$$

$$EI \frac{d^4 y}{dx^4} = q$$

Sign convention:



Governing Equation and Boundary Condition

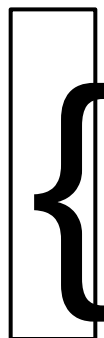
➤ Governing Equation

$$\frac{d^2}{dx^2} \left(EI \frac{d^2 v(x)}{dx^2} \right) - q(x) = 0, \quad 0 < x < L$$

➤ Boundary Conditions

$$v = ? \ \& \ \frac{dv}{dx} = ? \ \& \ EI \frac{d^2 v}{dx^2} = ? \ \& \ \frac{d}{dx} \left(EI \frac{d^2 v}{dx^2} \right) = ?, \quad \text{at } x = 0$$

$$v = ? \ \& \ \frac{dv}{dx} = ? \ \& \ EI \frac{d^2 v}{dx^2} = ? \ \& \ \frac{d}{dx} \left(EI \frac{d^2 v}{dx^2} \right) = ?, \quad \text{at } x = L \quad \frac{dv}{dx}$$



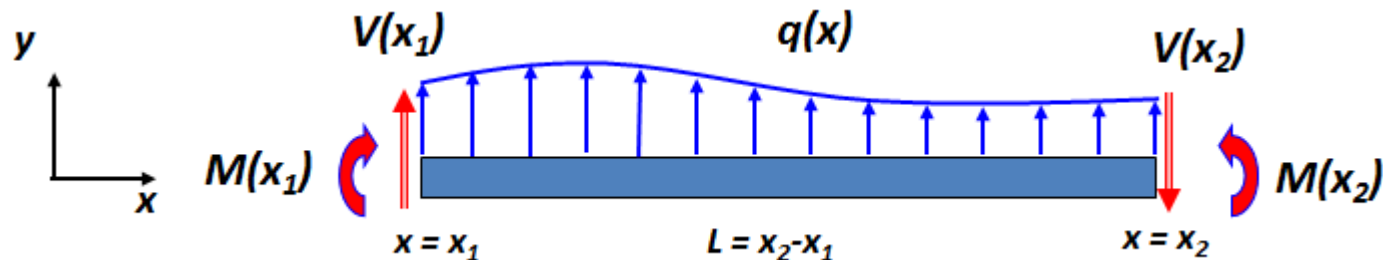
Essential BCs – if v or $\frac{dv}{dx}$ is specified at the boundary.

Natural BCs – if $EI \frac{d^2 v}{dx^2}$ or $\frac{d}{dx} \left(EI \frac{d^2 v}{dx^2} \right)$ is specified at the boundary.

$$EI \frac{d^2 v}{dx^2} \quad \frac{d}{dx} \left(EI \frac{d^2 v}{dx^2} \right)$$

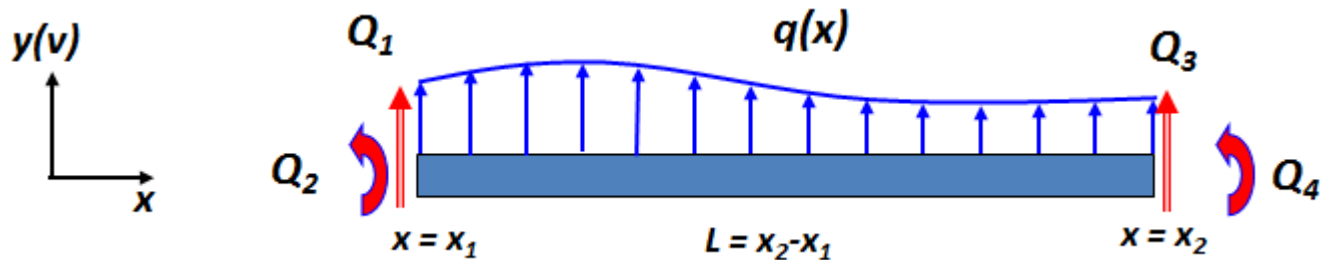
Weak Form from Integration-by-Parts ----- (2nd time)

$$0 = \int_{x_1}^{x_2} \left[\frac{d^2 w}{dx^2} \left(EI \frac{d^2 v}{dx^2} \right) - wq \right] dx + w \frac{d}{dx} \left(EI \frac{d^2 v}{dx^2} \right) \Big|_{x_1}^{x_2} - \frac{dw}{dx} \left(EI \frac{d^2 v}{dx^2} \right) \Big|_{x_1}^{x_2}$$



$$0 = \int_{x_1}^{x_2} \left[\frac{d^2 w}{dx^2} \left(EI \frac{d^2 v}{dx^2} \right) - wq \right] dx + \left[wV - \frac{dw}{dx} M \right] \Big|_{x_1}^{x_2}$$

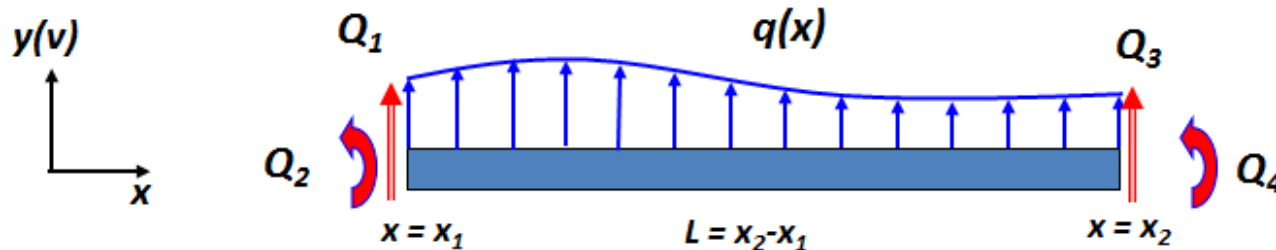
$$0 = \int_{x_1}^{x_2} \left[\frac{d^2 w}{dx^2} \left(EI \frac{d^2 v}{dx^2} \right) - wq \right] dx + \left[wV - \frac{dw}{dx} M \right]_{x_1}^{x_2}$$



$$Q_1 = V(x_1), \quad Q_2 = -M(x_1), \quad Q_3 = -V(x_2), \quad Q_4 = M(x_2)$$

$$\int_{x_1}^{x_2} \left[\frac{d^2 w}{dx^2} \left(EI \frac{d^2 v}{dx^2} \right) - wq \right] dx = w(x_1)Q_1 + w(x_2)Q_3 + \frac{dw}{dx} \Big|_1 Q_2 + \frac{dw}{dx} \Big|_2 Q_4$$

Ritz Method for Approximation



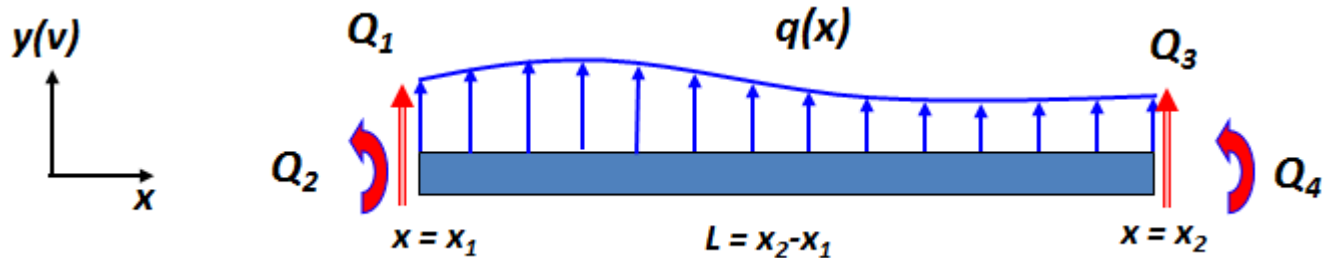
$$\text{Let } v(x) = \sum_{j=1}^n u_j \phi_j(x) \text{ and } n = 4$$

$$u_1 = v(x_1); \quad u_2 = \left. \frac{dv}{dx} \right|_{x=x_1}; \quad u_3 = v(x_2); \quad u_4 = \left. \frac{dv}{dx} \right|_{x=x_2};$$

$$\int_{x_1}^{x_2} \left[\frac{d^2 w}{dx^2} \left(EI \sum_{j=1}^4 u_j \frac{d^2 \phi_j}{dx^2} \right) - wq \right] dx = w(x_1)Q_1 + w(x_2)Q_3 + \left. \frac{dw}{dx} \right|_1 Q_2 + \left. \frac{dw}{dx} \right|_2 Q_4$$

⊙ Let $w(x) = f_i(x)$, $i = 1, 2, 3, 4$

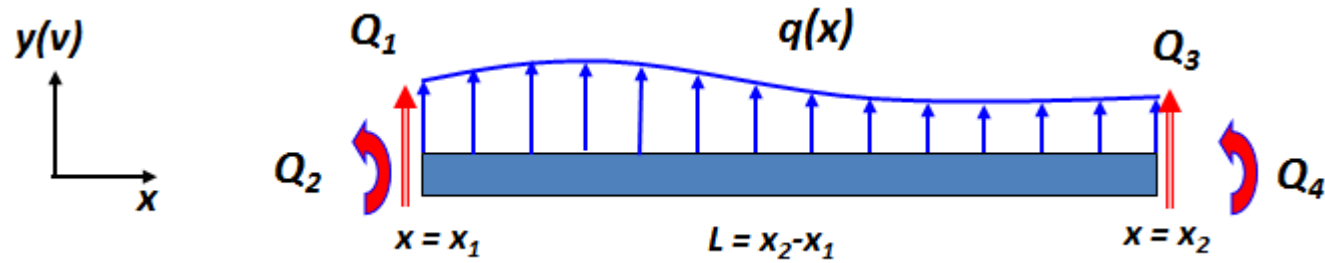
$$\int_{x_1}^{x_2} \left[\frac{d^2 \phi_i}{dx^2} \left(EI \sum_{j=1}^4 u_j \frac{d^2 \phi_j}{dx^2} \right) - \phi_i q \right] dx = \phi_i(x_1)Q_1 + \phi_i(x_2)Q_3 + \left. \frac{d\phi_i}{dx} \right|_1 Q_2 + \left. \frac{d\phi_i}{dx} \right|_2 Q_4$$



$$\left[\left(\phi_i \Big|_{x_1} \right) Q_1 + \left(\frac{d\phi_i}{dx} \Big|_{x_1} \right) Q_2 + \left(\phi_i \Big|_{x_2} \right) Q_3 + \left(\frac{d\phi_i}{dx} \Big|_{x_2} \right) Q_4 \right] = \sum_{j=1}^4 K_{ij} u_j - q_i$$

$$\text{where } K_{ij} = \int_{x_1}^{x_2} EI \left(\frac{d^2 \phi_i}{dx^2} \frac{d^2 \phi_j}{dx^2} \right) dx \text{ and } q_i = \int_{x_1}^{x_2} \phi_i q dx$$

Ritz Method for Approximation



$$\begin{bmatrix} \left(\phi_1 \Big|_{x_1} \right) & \left(\frac{d\phi_1}{dx} \Big|_{x_1} \right) & \left(\phi_1 \Big|_{x_2} \right) & \left(\frac{d\phi_1}{dx} \Big|_{x_2} \right) \\ \left(\phi_2 \Big|_{x_1} \right) & \left(\frac{d\phi_2}{dx} \Big|_{x_1} \right) & \left(\phi_2 \Big|_{x_2} \right) & \left(\frac{d\phi_2}{dx} \Big|_{x_2} \right) \\ \left(\phi_3 \Big|_{x_1} \right) & \left(\frac{d\phi_3}{dx} \Big|_{x_1} \right) & \left(\phi_3 \Big|_{x_2} \right) & \left(\frac{d\phi_3}{dx} \Big|_{x_2} \right) \\ \left(\phi_4 \Big|_{x_1} \right) & \left(\frac{d\phi_4}{dx} \Big|_{x_1} \right) & \left(\phi_4 \Big|_{x_2} \right) & \left(\frac{d\phi_4}{dx} \Big|_{x_2} \right) \end{bmatrix} \begin{Bmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \end{Bmatrix} = \begin{bmatrix} K_{11} & K_{12} & K_{13} & K_{14} \\ K_{21} & K_{22} & K_{23} & K_{24} \\ K_{31} & K_{32} & K_{33} & K_{34} \\ K_{41} & K_{42} & K_{43} & K_{44} \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{Bmatrix} - \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{Bmatrix}$$

Selection of Shape Function

$$\begin{bmatrix} \left(\phi_1\right|_{x_1}) & \left(\frac{d\phi_1}{dx}\right|_{x_1}) & \left(\phi_1\right|_{x_2}) & \left(\frac{d\phi_1}{dx}\right|_{x_2}) \\ \left(\phi_2\right|_{x_1}) & \left(\frac{d\phi_2}{dx}\right|_{x_1}) & \left(\phi_2\right|_{x_2}) & \left(\frac{d\phi_2}{dx}\right|_{x_2}) \\ \left(\phi_3\right|_{x_1}) & \left(\frac{d\phi_3}{dx}\right|_{x_1}) & \left(\phi_3\right|_{x_2}) & \left(\frac{d\phi_3}{dx}\right|_{x_2}) \\ \left(\phi_4\right|_{x_1}) & \left(\frac{d\phi_4}{dx}\right|_{x_1}) & \left(\phi_4\right|_{x_2}) & \left(\frac{d\phi_4}{dx}\right|_{x_2}) \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$



Interpolation
Properties



$$\begin{Bmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \end{Bmatrix} = \begin{bmatrix} K_{11} & K_{12} & K_{13} & K_{14} \\ K_{12} & K_{22} & K_{23} & K_{24} \\ K_{13} & K_{23} & K_{33} & K_{34} \\ K_{14} & K_{24} & K_{34} & K_{44} \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{Bmatrix} - \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{Bmatrix}$$

Derivation of Shape Function for Beam Element

Local Coordinates

$$v(\xi) = \tilde{u}_1 \phi_1 + \tilde{u}_2 \phi_2 + \tilde{u}_3 \phi_3 + \tilde{u}_4 \phi_4$$

and

$$\frac{dv(\xi)}{d\xi} = \tilde{u}_1 \frac{d\phi_1}{d\xi} + \tilde{u}_2 \frac{d\phi_2}{d\xi} + \tilde{u}_3 \frac{d\phi_3}{d\xi} + \tilde{u}_4 \frac{d\phi_4}{d\xi}$$

where

$$\tilde{u}_1 = v_1 \quad \tilde{u}_2 = \frac{dv_1}{d\xi} \quad \tilde{u}_3 = v_2 \quad \tilde{u}_4 = \frac{dv_2}{d\xi}$$

Let

$$\phi_i = a_i + b_i \xi + c_i \xi^2 + d_i \xi^3$$

Find coefficients to satisfy the interpolation properties.

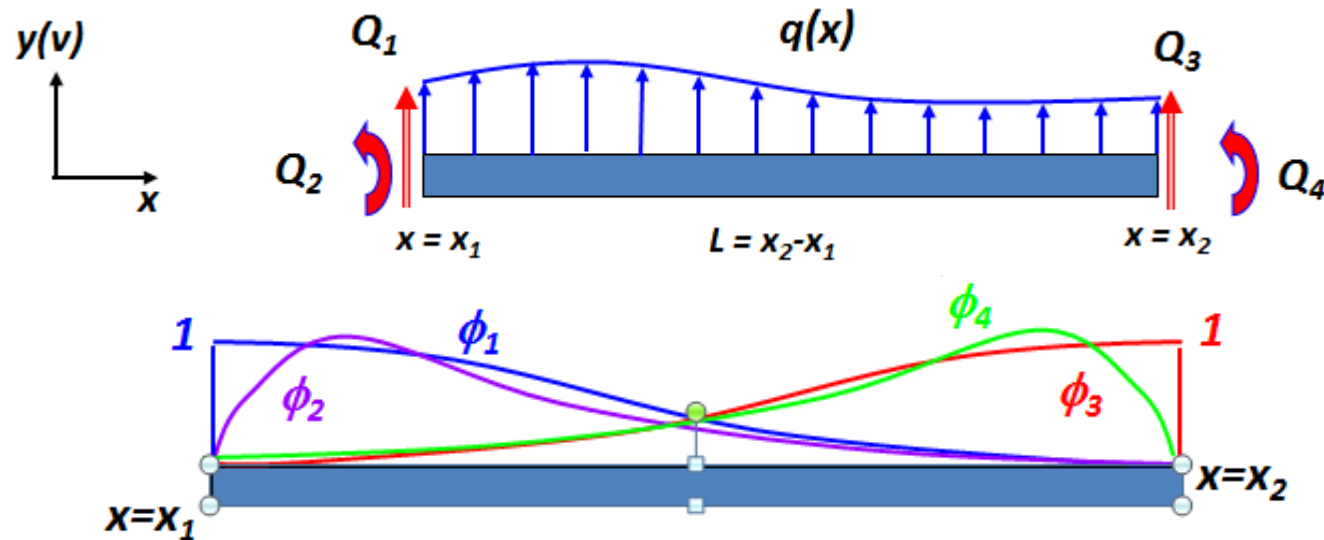
Derivation of Shape Function for Beam Element

In the global coordinates:

$$v(x) = v_1 \phi_1(x) + \frac{l}{2} \frac{dv_1}{dx} \phi_2(x) + v_2 \phi_3(x) + \frac{l}{2} \frac{dv_2}{dx} \phi_4(x)$$

$$\begin{Bmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_4 \end{Bmatrix} = \begin{Bmatrix} 1 - 3\left(\frac{x - x_1}{x_2 - x_1}\right)^2 + 2\left(\frac{x - x_1}{x_2 - x_1}\right)^3 \\ \frac{2}{l}(x - x_1)\left(1 - \frac{x - x_1}{x_2 - x_1}\right)^2 \\ 3\left(\frac{x - x_1}{x_2 - x_1}\right)^2 - 2\left(\frac{x - x_1}{x_2 - x_1}\right)^3 \\ \frac{2}{l}(x - x_1)\left[\left(\frac{x - x_1}{x_2 - x_1}\right)^2 - \frac{x - x_1}{x_2 - x_1}\right] \end{Bmatrix}$$

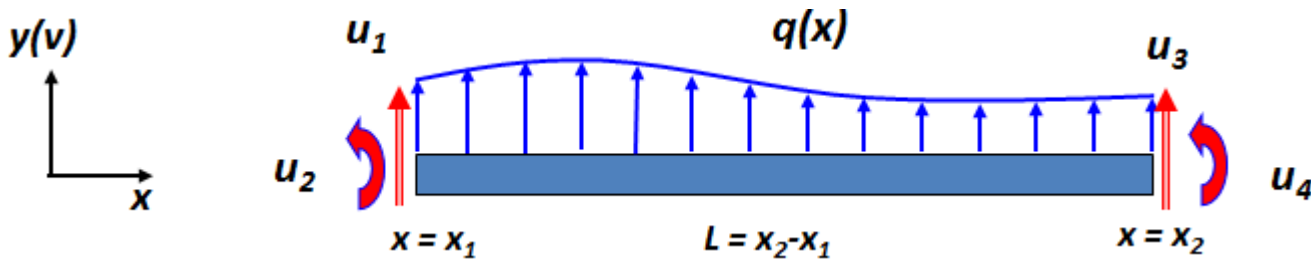
Element Equations of 4th Order 1-D Model



$$\begin{Bmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \end{Bmatrix} + \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{Bmatrix} = \begin{bmatrix} K_{11} & K_{12} & K_{13} & K_{14} \\ K_{12} & K_{22} & K_{23} & K_{24} \\ K_{13} & K_{23} & K_{33} & K_{34} \\ K_{14} & K_{24} & K_{34} & K_{44} \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{Bmatrix}$$

where $K_{ij} = \int_{x_1}^{x_2} EI \left(\frac{d^2 \phi_i}{dx^2} \frac{d^2 \phi_j}{dx^2} \right) dx = K_{ji}$ and $q_i = \int_{x_1}^{x_2} \phi_i q dx$

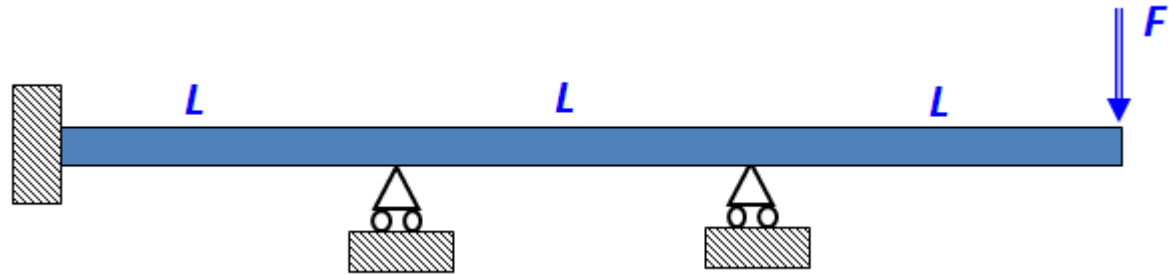
Element Equations of 4th Order 1-D Model



$$\begin{Bmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \end{Bmatrix} + \begin{Bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 6 & 3L & -6 & 3L \\ 3L & 2L^2 & -3L & L^2 \\ -6 & -3L & 6 & -3L \\ 3L & L^2 & -3L & 2L^2 \end{bmatrix} \begin{Bmatrix} u_1 = v_1 \\ u_2 = \theta_1 \\ u_3 = v_2 \\ u_4 = \theta_2 \end{Bmatrix}$$

$$\text{where } q_i = \int_{x_1}^{x_2} \phi_i q dx$$

Example 1.



Governing equation:

$$\frac{d^2}{dx^2} \left(EI \frac{d^2 v}{dx^2} \right) - q(x) = 0 \quad 0 < x < L$$

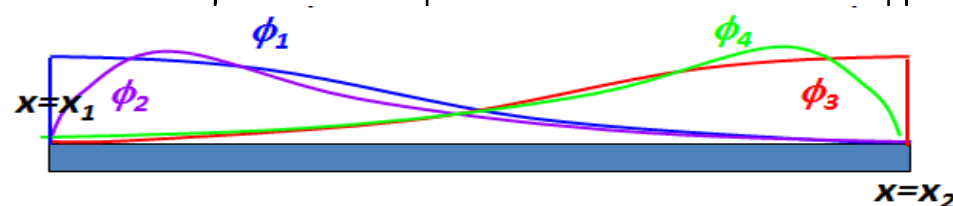
Weak form for one element

$$\text{where } \int_{x_1}^{x_2} \left(EI \frac{d^2 w}{dx^2} \frac{d^2 v}{dx^2} - wq \right) dx - w(x_1)Q_1 - \left. \frac{dw}{dx} \right|_{x_1} Q_2 - w(x_2)Q_3 - \left. \frac{dw}{dx} \right|_{x_2} Q_4 = 0$$

$$Q_1 = V(x_1) \quad Q_2 = -M(x_1) \quad Q_3 = -V(x_2) \quad Q_4 = M(x_2)$$

Finite Element Analysis of 1-D Problems

⊙ Approximation function: $v(x) = v_1 \phi_1(x) + \frac{l}{2} \frac{dv_1}{dx} \phi_2(x) + v_2 \phi_3(x) + \frac{l}{2} \frac{dv_2}{dx} \phi_4(x)$

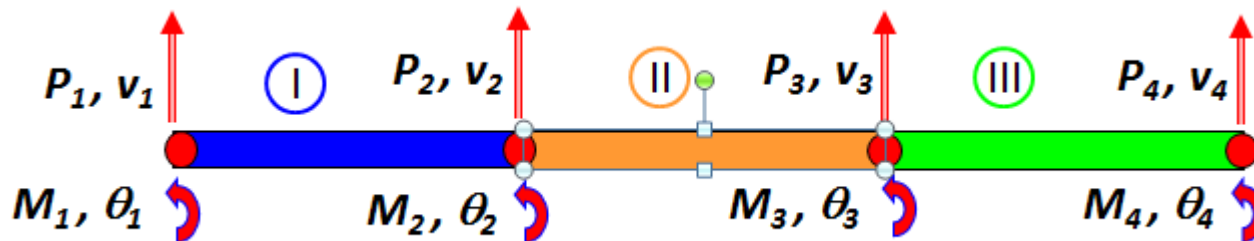
$$\begin{Bmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_4 \end{Bmatrix} = \begin{Bmatrix} 1 - 3\left(\frac{x-x_1}{x_2-x_1}\right)^2 + 2\left(\frac{x-x_1}{x_2-x_1}\right)^3 \\ \frac{2}{l}(x-x_1)\left(1 - \frac{x-x_1}{x_2-x_1}\right)^2 \\ 3\left(\frac{x-x_1}{x_2-x_1}\right)^2 - 2\left(\frac{x-x_1}{x_2-x_1}\right)^3 \\ \frac{2}{l}(x-x_1)\left[\left(\frac{x-x_1}{x_2-x_1}\right)^2 - \frac{x-x_1}{x_2-x_1}\right] \end{Bmatrix}$$


Finite Element Analysis of 1-D Problems

Finite element model:

$$\begin{Bmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 6 & 3L & -6 & 3L \\ 3L & 2L^2 & -3L & L^2 \\ -6 & -3L & 6 & -3L \\ 3L & L^2 & -3L & 2L^2 \end{bmatrix} \begin{Bmatrix} v_1 \\ \theta_1 \\ v_2 \\ \theta_2 \end{Bmatrix}$$

Discretization:



Matrix Assembly of Multiple Beam Elements

Element I

$$\begin{Bmatrix} Q_1^I \\ Q_2^I \\ Q_3^I \\ Q_4^I \\ 0 \\ 0 \\ 0 \\ 0 \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 6 & 3L & -6 & 3L & 0 & 0 & 0 & 0 \\ 3L & 2L^2 & -3L & L^2 & 0 & 0 & 0 & 0 \\ -6 & -3L & 6 & -3L & 0 & 0 & 0 & 0 \\ 3L & L^2 & -3L & 2L^2 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} v_1 \\ \theta_1 \\ v_2 \\ \theta_2 \\ v_3 \\ \theta_3 \\ v_4 \\ \theta_4 \end{Bmatrix}$$

Element II

$$\begin{Bmatrix} 0 \\ 0 \\ Q_1^{II} \\ Q_2^{II} \\ Q_3^{II} \\ Q_4^{II} \\ 0 \\ 0 \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 6 & 3L & -6 & 3L & 0 & 0 \\ 0 & 0 & 3L & 2L^2 & -3L & L^2 & 0 & 0 \\ 0 & 0 & -6 & -3L & 6 & -3L & 0 & 0 \\ 0 & 0 & 3L & L^2 & -3L & 2L^2 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} v_1 \\ \theta_1 \\ v_2 \\ \theta_2 \\ v_3 \\ \theta_3 \\ v_4 \\ \theta_4 \end{Bmatrix}$$

Matrix Assembly of Multiple Beam Elements

Element III

$$\begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ Q_1^{III} \\ Q_2^{III} \\ Q_3^{III} \\ Q_4^{III} \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 6 & 3L & -6 & 3L \\ 0 & 0 & 0 & 0 & 3L & 2L^2 & -3L & L^2 \\ 0 & 0 & 0 & 0 & -6 & -3L & 6 & -3L \\ 0 & 0 & 0 & 0 & 3L & L^2 & -3L & 2L^2 \end{bmatrix} \begin{Bmatrix} v_1 \\ \theta_1 \\ v_2 \\ \theta_2 \\ v_3 \\ \theta_3 \\ v_4 \\ \theta_4 \end{Bmatrix}$$

$$\begin{Bmatrix} P_1 \\ M_1 \\ P_2 \\ M_2 \\ P_3 \\ M_3 \\ P_4 \\ M_4 \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 6 & 3L & -6 & 3L & 0 & 0 & 0 & 0 \\ 3L & 2L^2 & -3L & L^2 & 0 & 0 & 0 & 0 \\ -6 & -3L & 6+6 & -3L+3L & -6 & 3L & 0 & 0 \\ \hline 3L & L^2 & -3L+3L & 2L^2+2L^2 & -3L & L^2 & 0 & 0 \\ 0 & 0 & -6 & -3L & 6+6 & -3L+3L & -6 & 3L \\ 0 & 0 & 3L & L^2 & -3L+3L & 2L^2+2L^2 & -3L & L^2 \\ \hline 0 & 0 & 0 & 0 & -6 & -3L & 6 & -3L \\ 0 & 0 & 0 & 0 & 3L & L^2 & -3L & 2L^2 \end{bmatrix} \begin{Bmatrix} v_1 \\ \theta_1 \\ v_2 \\ \theta_2 \\ v_3 \\ \theta_3 \\ v_4 \\ \theta_4 \end{Bmatrix}$$

Solution Procedures

Apply known boundary conditions

$$\begin{Bmatrix} P_1 = ? \\ M_1 = ? \\ P_2 = ? \\ M_2 = 0 \\ P_3 = ? \\ M_3 = 0 \\ P_4 = -F \\ M_4 = 0 \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 6 & 3L & -6 & 3L & 0 & 0 & 0 & 0 \\ 3L & 2L^2 & -3L & L^2 & 0 & 0 & 0 & 0 \\ -6 & -3L & 12 & 0 & -6 & 3L & 0 & 0 \\ 3L & L^2 & 0 & 4L^2 & -3L & L^2 & 0 & 0 \\ 0 & 0 & -6 & -3L & 12 & 0 & -6 & 3L \\ 0 & 0 & 3L & L^2 & 0 & 4L^2 & -3L & L^2 \\ 0 & 0 & 0 & 0 & -6 & -3L & 6 & -3L \\ 0 & 0 & 0 & 0 & 3L & L^2 & -3L & 2L^2 \end{bmatrix} \begin{Bmatrix} v_1 = 0 \\ \theta_1 = 0 \\ v_2 = 0 \\ \theta_2 = ? \\ v_3 = 0 \\ \theta_3 = ? \\ v_4 = ? \\ \theta_4 = ? \end{Bmatrix}$$

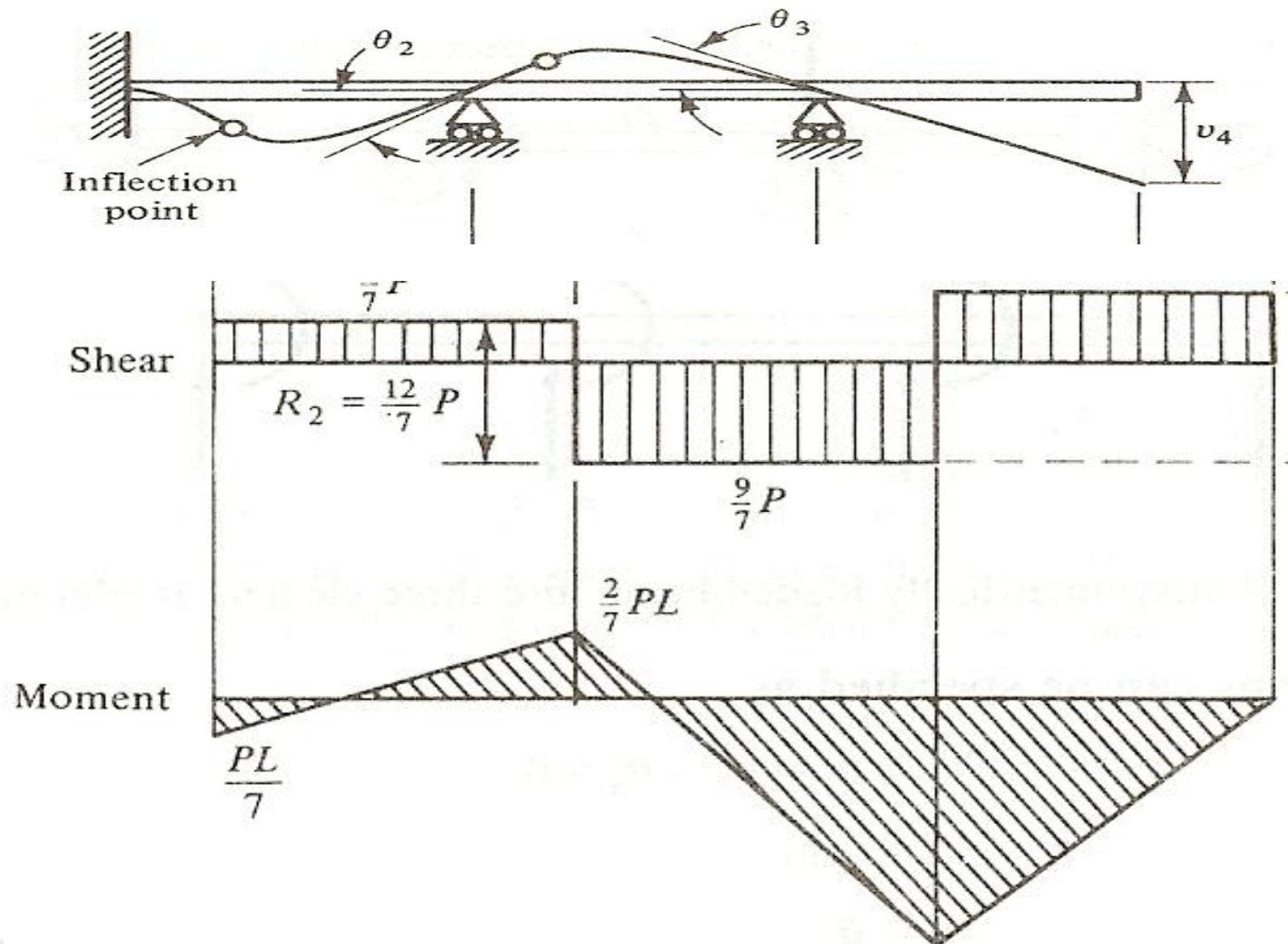
$$\begin{Bmatrix} M_2 = 0 \\ M_3 = 0 \\ P_4 = -F \\ M_4 = 0 \\ P_1 = ? \\ M_1 = ? \\ P_2 = ? \\ P_3 = ? \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 3L & L^2 & 0 & 4L^2 & -3L & L^2 & 0 & 0 \\ 0 & 0 & 3L & L^2 & 0 & 4L^2 & -3L & L^2 \\ 0 & 0 & 0 & 0 & -6 & -3L & 6 & -3L \\ 0 & 0 & 0 & 0 & 3L & L^2 & -3L & 2L^2 \\ 6 & 3L & -6 & 3L & 0 & 0 & 0 & 0 \\ 3L & 2L^2 & -3L & L^2 & 0 & 0 & 0 & 0 \\ -6 & -3L & 12 & 0 & -6 & 3L & 0 & 0 \\ 0 & 0 & -6 & -3L & 12 & 0 & -6 & 3L \end{bmatrix} \begin{Bmatrix} v_1 = 0 \\ \theta_1 = 0 \\ v_2 = 0 \\ \theta_2 = ? \\ v_3 = 0 \\ \theta_3 = ? \\ v_4 = ? \\ \theta_4 = ? \end{Bmatrix}$$

Solution Procedures

$$\begin{Bmatrix} M_2 = 0 \\ M_3 = 0 \\ P_4 = -F \\ M_4 = 0 \\ P_1 = ? \\ M_1 = ? \\ P_2 = ? \\ P_3 = ? \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 3L & L^2 & 0 & -3L & 4L^2 & L^2 & 0 & 0 \\ 0 & 0 & 3L & 0 & L^2 & 4L^2 & -3L & L^2 \\ 0 & 0 & 0 & -6 & 0 & -3L & 6 & -3L \\ 0 & 0 & 0 & 3L & 0 & L^2 & -3L & 2L^2 \\ \hline 6 & 3L & -6 & 0 & 3L & 0 & 0 & 0 \\ 3L & 2L^2 & -3L & 0 & L^2 & 0 & 0 & 0 \\ -6 & -3L & 12 & -6 & 0 & 3L & 0 & 0 \\ 0 & 0 & -6 & -12 & -3L & 0 & -6 & 3L \end{bmatrix} \begin{Bmatrix} v_1 = 0 \\ \theta_1 = 0 \\ v_2 = 0 \\ v_3 = 0 \\ \theta_2 = ? \\ \theta_3 = ? \\ v_4 = ? \\ \theta_4 = ? \end{Bmatrix}$$

$$\begin{Bmatrix} M_2 = 0 \\ M_3 = 0 \\ P_4 = -F \\ M_4 = 0 \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 4L^2 & L^2 & 0 & 0 \\ L^2 & 4L^2 & -3L & L^2 \\ 0 & -3L & 6 & -3L \\ 0 & L^2 & -3L & 2L^2 \end{bmatrix} \begin{Bmatrix} \theta_2 = ? \\ \theta_3 = ? \\ v_4 = ? \\ \theta_4 = ? \end{Bmatrix} \quad \begin{Bmatrix} P_1 = ? \\ M_1 = ? \\ P_2 = ? \\ P_3 = ? \end{Bmatrix} = \frac{2EI}{L^3} \begin{bmatrix} 3L & 0 & 0 & 0 \\ L^2 & 0 & 0 & 0 \\ 0 & 3L & 0 & 0 \\ -3L & 0 & -6 & 3L \end{bmatrix} \begin{Bmatrix} \theta_2 \\ \theta_3 \\ v_4 \\ \theta_4 \end{Bmatrix}$$

Shear Resultant & Bending Moment Diagram



UNIT-3

CONTINUUM ELEMENTS

Course Learning Outcomes

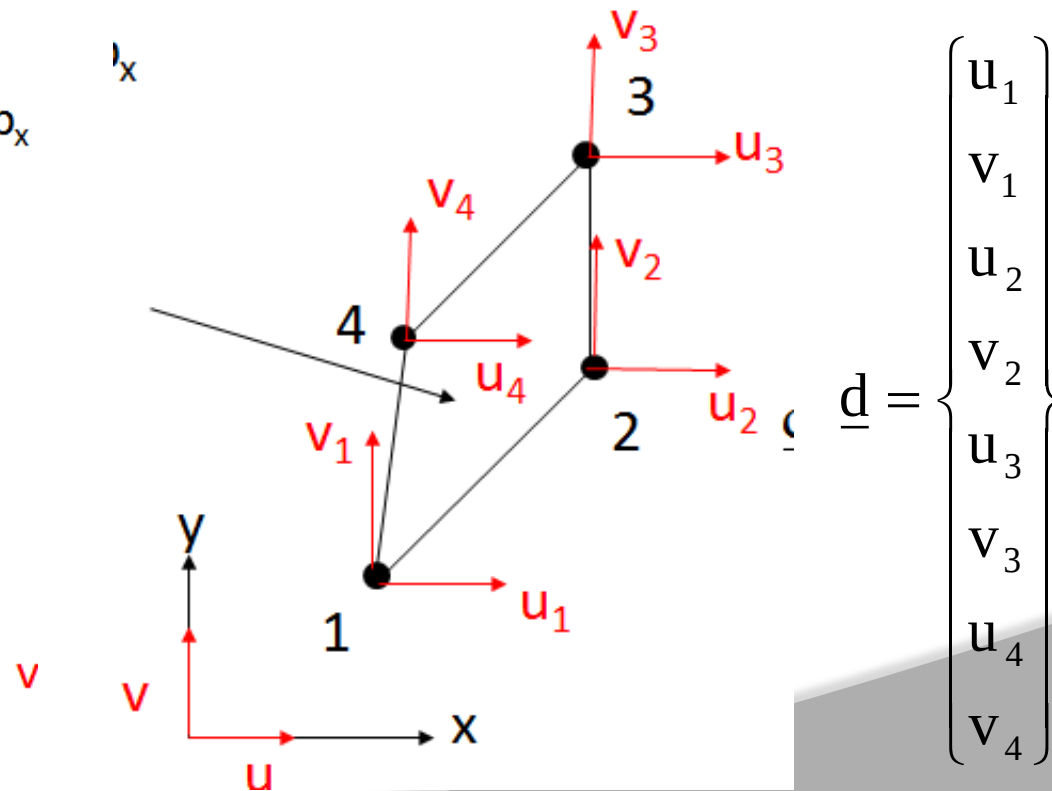
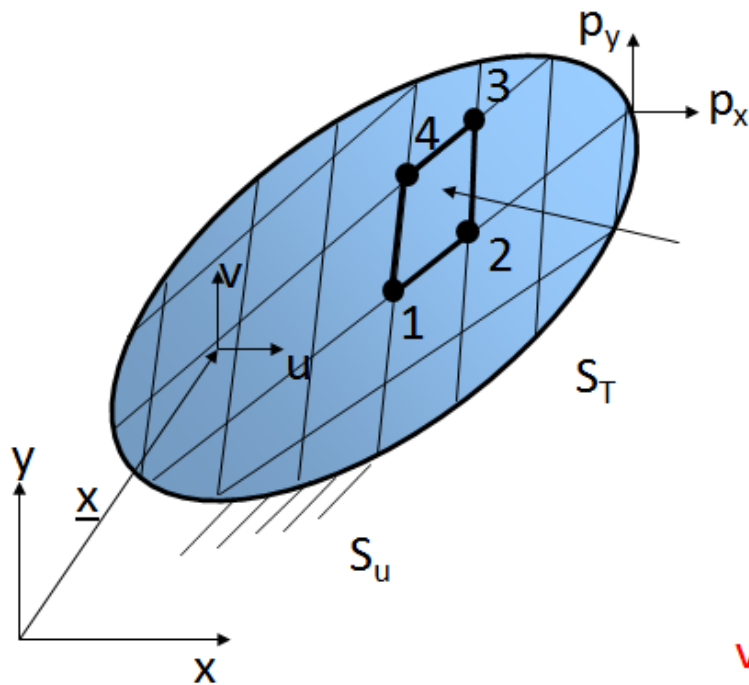
CLOs	Course Learning Outcomes
CLO 1	Derive the element stiffness matrices for triangular elements and axi- symmetric solids and estimate the load vector and stresses.
CLO 2	Formulate simple and complex problems into finite elements and solve structural and thermal problems
CLO 3	Understand the concept of CST and LST and their shape functions.

Introduction

- Computation of shape functions for constant strain triangle
- Properties of the shape functions
- Computation of strain-displacement matrix
- Computation of element stiffness matrix
- Computation of nodal loads due to body forces
- Computation of nodal loads due to traction
- Recommendations for use
- Example problems

Finite element formulation for 2D

- Divide the body into connected to each other through special points (“nodes”)



$$u(x, y) \approx N_1(x, y) u_1 + N_2(x, y) u_2 + N_3(x, y) u_3 + N_4(x, y) u_4$$

$$v(x, y) \approx N_1(x, y) v_1 + N_2(x, y) v_2 + N_3(x, y) v_3 + N_4(x, y) v_4$$

$$\underline{u} = \begin{Bmatrix} u(x, y) \\ v(x, y) \end{Bmatrix} = \begin{bmatrix} N_1 & 0 & N_2 & 0 & N_3 & 0 & N_4 & 0 \\ 0 & N_1 & 0 & N_2 & 0 & N_3 & 0 & N_4 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{Bmatrix}$$

$$\underline{u} = \underline{N} \underline{d}$$

TASK 2: APPROXIMATE THE STRAIN and STRESS WITHIN EACH ELEMENT

Approximation of the strain in element 'e'

$$\varepsilon_x = \frac{\partial u(x, y)}{\partial x} \approx \frac{\partial N_1(x, y)}{\partial x} u_1 + \frac{\partial N_2(x, y)}{\partial x} u_2 + \frac{\partial N_3(x, y)}{\partial x} u_3 + \frac{\partial N_4(x, y)}{\partial x} u_4$$

$$\varepsilon_y = \frac{\partial v(x, y)}{\partial y} \approx \frac{\partial N_1(x, y)}{\partial y} v_1 + \frac{\partial N_2(x, y)}{\partial y} v_2 + \frac{\partial N_3(x, y)}{\partial y} v_3 + \frac{\partial N_4(x, y)}{\partial y} v_4$$

$$\gamma_{xy} = \frac{\partial u(x, y)}{\partial y} + \frac{\partial v(x, y)}{\partial x} \approx \frac{\partial N_1(x, y)}{\partial y} u_1 + \frac{\partial N_1(x, y)}{\partial x} v_1 + \dots$$

$$\underline{\varepsilon} = \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix}$$

$$= \underbrace{\begin{bmatrix} \frac{\partial N_1(x, y)}{\partial x} & 0 & \frac{\partial N_2(x, y)}{\partial x} & 0 & \frac{\partial N_3(x, y)}{\partial x} & 0 & \frac{\partial N_4(x, y)}{\partial x} & 0 \\ 0 & \frac{\partial N_1(x, y)}{\partial y} & 0 & \frac{\partial N_2(x, y)}{\partial y} & 0 & \frac{\partial N_3(x, y)}{\partial y} & 0 & \frac{\partial N_4(x, y)}{\partial y} \\ \frac{\partial N_1(x, y)}{\partial y} & \frac{\partial N_1(x, y)}{\partial x} & \frac{\partial N_2(x, y)}{\partial y} & \frac{\partial N_2(x, y)}{\partial x} & \frac{\partial N_3(x, y)}{\partial y} & \frac{\partial N_3(x, y)}{\partial x} & \frac{\partial N_4(x, y)}{\partial y} & \frac{\partial N_4(x, y)}{\partial x} \end{bmatrix}}_{\underline{B}} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{Bmatrix}$$

$$\underline{\varepsilon} = \underline{B} \underline{d}$$

Summary: For each element

- Displacement approximation in terms of shape functions

$$\underline{u} = \underline{N} \underline{d}$$

- Strain approximation in terms of strain-displacement

$$\underline{\varepsilon} = \underline{B} \underline{d}$$

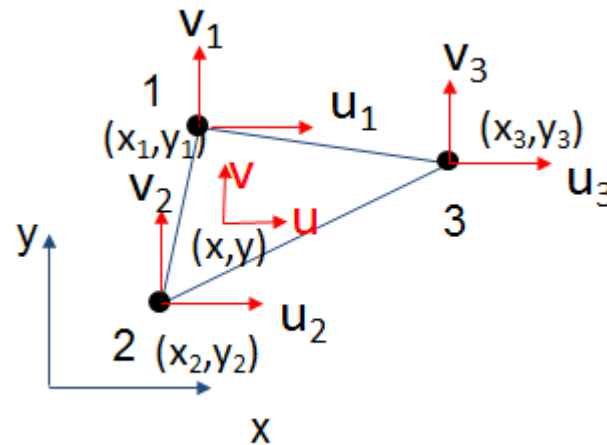
- Stress approximation $\underline{\sigma} = \underline{D} \underline{B} \underline{d}$

- Element stiffness matrix

$$\underline{k} = \int_{V^e} \underline{B}^T \underline{D} \underline{B} dV$$

- Element nodal load vector $\underline{f} = \underbrace{\int_{V^e} \underline{N}^T \underline{X} dV}_{\underline{f}_b} + \underbrace{\int_{S_T^e} \underline{N}^T \underline{T}_S dS}_{\underline{f}_s}$

⦿ **Constant Strain Triangle (CST)** : Simplest 2D finite element



- ⦿ 3 nodes per element
- ⦿ 2 dofs per node (each node can move in x- and y- directions)
- ⦿ Hence 6 dofs per element

- The displacement approximation in terms of shape functions is

$$u(x,y) \approx N_1 u_1 + N_2 u_2 + N_3 u_3$$

$$v(x,y) \approx N_1 v_1 + N_2 v_2 + N_3 v_3$$

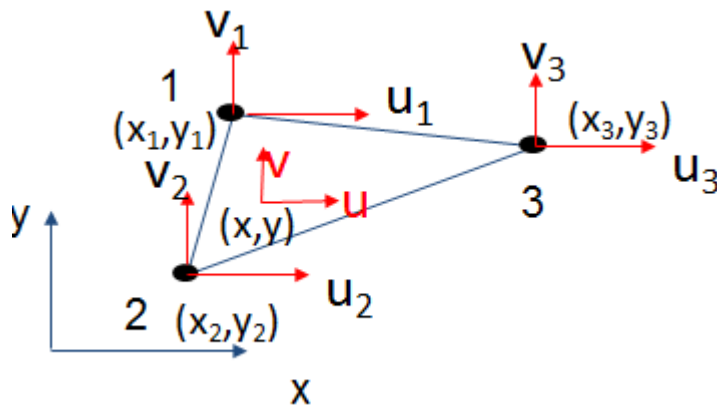
$$\underline{u} = \begin{Bmatrix} u(x,y) \\ v(x,y) \end{Bmatrix} = \begin{bmatrix} N_1 & 0 & N_2 & 0 & N_3 & 0 \\ 0 & N_1 & 0 & N_2 & 0 & N_3 \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}$$

$$\underline{u}_{2 \times 1} = \underline{N}_{2 \times 6} \underline{d}_{6 \times 1}$$

$$\underline{N} = \begin{bmatrix} N_1 & 0 & N_2 & 0 & N_3 & 0 \\ 0 & N_1 & 0 & N_2 & 0 & N_3 \end{bmatrix}$$

- Formula for the shape functions are

where



$$N_1 = \frac{a_1 + b_1x + c_1y}{2A}$$

$$N_2 = \frac{a_2 + b_2x + c_2y}{2A}$$

$$N_3 = \frac{a_3 + b_3x + c_3y}{2A}$$

$$A = \text{area of triangle} = \frac{1}{2} \det \begin{bmatrix} 1 & x_1 & y_1 \\ 1 & x_2 & y_2 \\ 1 & x_3 & y_3 \end{bmatrix}$$

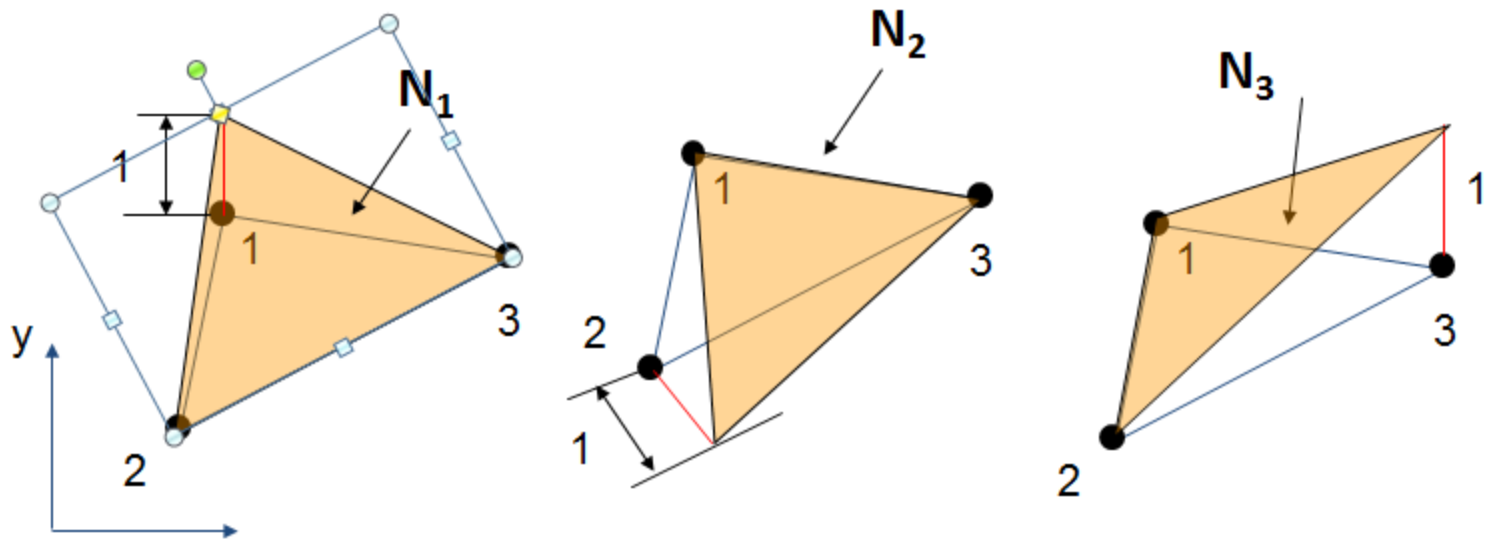
$$a_1 = x_2y_3 - x_3y_2 \quad b_1 = y_2 - y_3 \quad c_1 = x_3 - x_2$$

$$a_2 = x_3y_1 - x_1y_3 \quad b_2 = y_3 - y_1 \quad c_2 = x_1 - x_3$$

$$a_3 = x_1y_2 - x_2y_1 \quad b_3 = y_1 - y_2 \quad c_3 = x_2 - x_1$$

Properties of the shape functions

- ⦿ The shape functions N_1 , N_2 and N_3 are linear functions of x and y



$$N_i = \begin{cases} 1 & \text{at node 'i'} \\ 0 & \text{at other nodes} \end{cases}$$

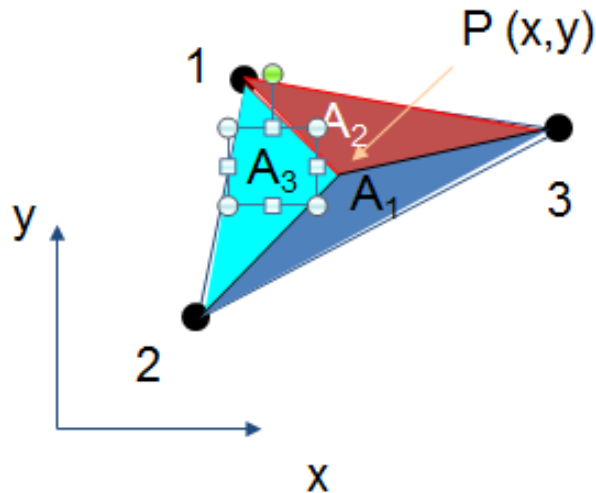
- At every point in the domain

$$\sum_{i=1}^3 N_i = 1$$

$$\sum_{i=1}^3 N_i x_i = x$$

$$\sum_{i=1}^3 N_i y_i = y$$

- Geometric interpretation of the shape functions, at any point $P(x,y)$ that the shape functions are evaluated



$$N_1 = \frac{A_1}{A}$$

$$N_2 = \frac{A_2}{A}$$

$$N_3 = \frac{A_3}{A}$$

Approximation of the strains

$$\underline{\varepsilon} = \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = \begin{Bmatrix} \frac{\partial u}{\partial x} \\ \frac{\partial v}{\partial y} \\ \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \end{Bmatrix} \approx \underline{B} \underline{d}$$

$$\underline{B} = \begin{bmatrix} \frac{\partial N_1(x, y)}{\partial x} & 0 & \frac{\partial N_2(x, y)}{\partial x} & 0 & \frac{\partial N_3(x, y)}{\partial x} & 0 \\ 0 & \frac{\partial N_1(x, y)}{\partial y} & 0 & \frac{\partial N_2(x, y)}{\partial y} & 0 & \frac{\partial N_3(x, y)}{\partial y} \\ \frac{\partial N_1(x, y)}{\partial y} & \frac{\partial N_1(x, y)}{\partial x} & \frac{\partial N_2(x, y)}{\partial y} & \frac{\partial N_2(x, y)}{\partial x} & \frac{\partial N_3(x, y)}{\partial y} & \frac{\partial N_3(x, y)}{\partial x} \end{bmatrix}$$

$$= \frac{1}{2A} \begin{bmatrix} b_1 & 0 & b_2 & 0 & b_3 & 0 \\ 0 & c_1 & 0 & c_2 & 0 & c_3 \\ c_1 & b_1 & c_2 & b_2 & c_3 & b_3 \end{bmatrix}$$

- ⦿ Inside each element, all components of strain are constant: hence the name Constant Strain Triangle.
- ⦿ Element stresses (constant inside each element).

$$\underline{\sigma} = \underline{DB} \underline{d}$$

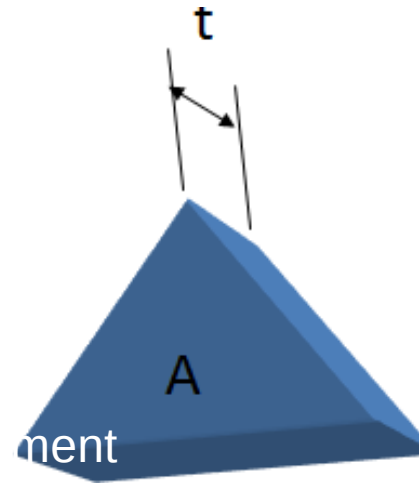
IMPORTANT NOTE:

- The displacement field is continuous across element boundaries
- The strains and stresses are NOT continuous across element boundaries

Element stiffness matrix

$$\underline{k} = \int_{V^e} \underline{B}^T \underline{D} \underline{B} dV$$

Since $\underline{B}$ is constant



$$\underline{k} = \underline{B}^T \underline{D} \underline{B} \int_{V^e} dV = \underline{B}^T \underline{D} \underline{B} A t$$

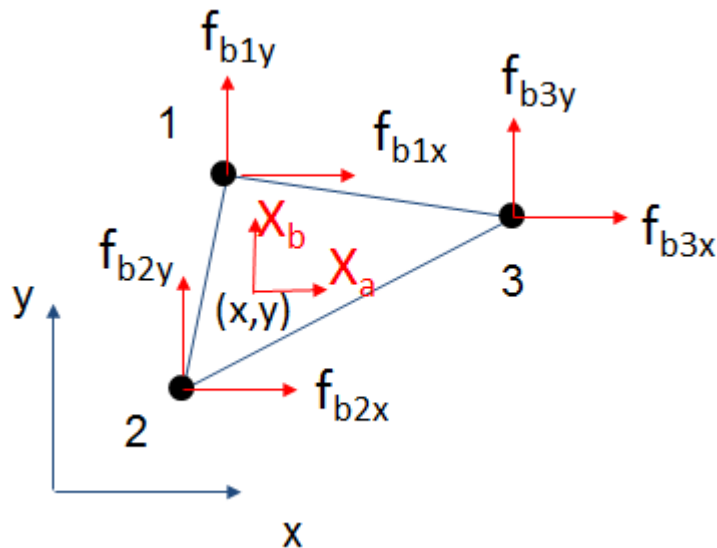
t=thickness of the element
A=surface area of the element

Element nodal load vector

$$\underline{f} = \underbrace{\int_{V^e} \underline{N}^T \underline{X} dV}_{\underline{f}_b} + \underbrace{\int_{S_T^e} \underline{N}^T \underline{T}_s dS}_{\underline{f}_s}$$

Element nodal load vector due to body forces

$$\underline{f}_b = \int_{V^e} \underline{N}^T \underline{X} dV = t \int_{A^e} \underline{N}^T \underline{X} dA$$



$$\underline{f}_b = \begin{Bmatrix} f_{b1x} \\ f_{b1y} \\ f_{b2x} \\ f_{b2y} \\ f_{b3x} \\ f_{b3y} \end{Bmatrix} = \begin{Bmatrix} t \int_{A^e} N_1 X_a dA \\ t \int_{A^e} N_1 X_b dA \\ t \int_{A^e} N_2 X_a dA \\ t \int_{A^e} N_2 X_b dA \\ t \int_{A^e} N_3 X_a dA \\ t \int_{A^e} N_3 X_b dA \end{Bmatrix}$$

EXAMPLE:

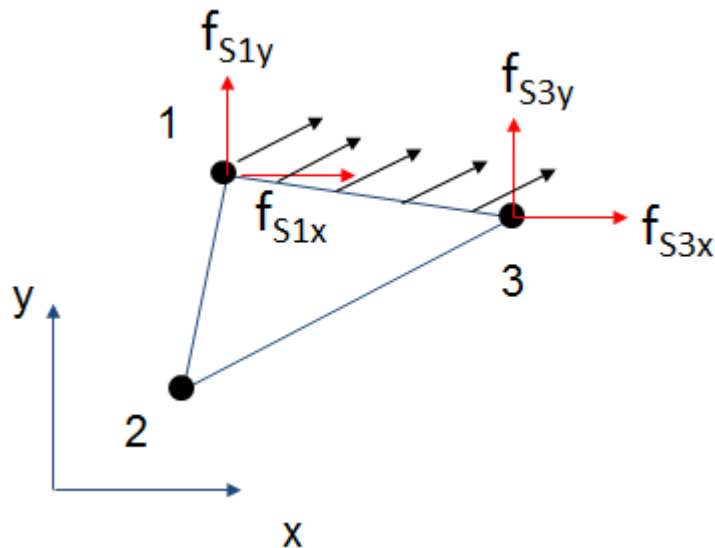
If $X_a=1$ and $X_b=0$

$$\underline{f}_b = \begin{Bmatrix} f_{b1x} \\ f_{b1y} \\ f_{b2x} \\ f_{b2y} \\ f_{b3x} \\ f_{b3y} \end{Bmatrix} = \begin{Bmatrix} t \int_{A^e} N_1 X_a dA \\ t \int_{A^e} N_1 X_b dA \\ t \int_{A^e} N_2 X_a dA \\ t \int_{A^e} N_2 X_b dA \\ t \int_{A^e} N_3 X_a dA \\ t \int_{A^e} N_3 X_b dA \end{Bmatrix} = \begin{Bmatrix} t \int_{A^e} N_1 dA \\ 0 \\ t \int_{A^e} N_2 dA \\ 0 \\ t \int_{A^e} N_3 dA \\ 0 \end{Bmatrix} = \begin{Bmatrix} \frac{tA}{3} \\ 0 \\ \frac{tA}{3} \\ 0 \\ \frac{tA}{3} \\ 0 \end{Bmatrix}$$

Element nodal load vector due to traction

$$\underline{f}_S = \int_{S_T^e} \underline{N}^T \underline{T}_S dS$$

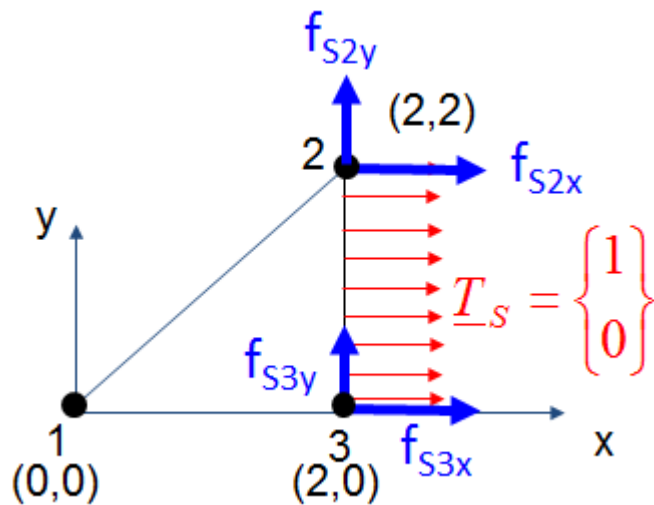
EXAMPLE:



$$\underline{f}_S = t \int_{l_{1-3}^e} \underline{N}^T \Big|_{along\ 1-3} \underline{T}_S dS$$

Element nodal load vector due to traction

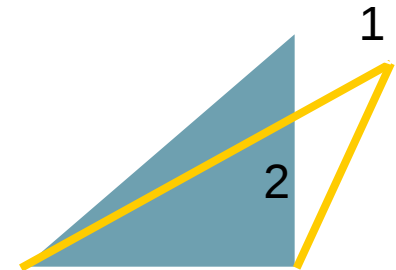
Example



$$\underline{f}_S = t \int_{l_{2-3}^e} \underline{N}^T \Big|_{\text{along } 2-3} \underline{T}_S dS$$

$$\begin{aligned} f_{S_{2x}} &= t \int_{l_{2-3}^e} N_2 \Big|_{\text{along } 2-3} (1) dy \\ &= t \left(\frac{1}{2} \right) \times 2 \times 1 = t \end{aligned}$$

Similarly, compute



$$f_{S_{2y}} = 0$$

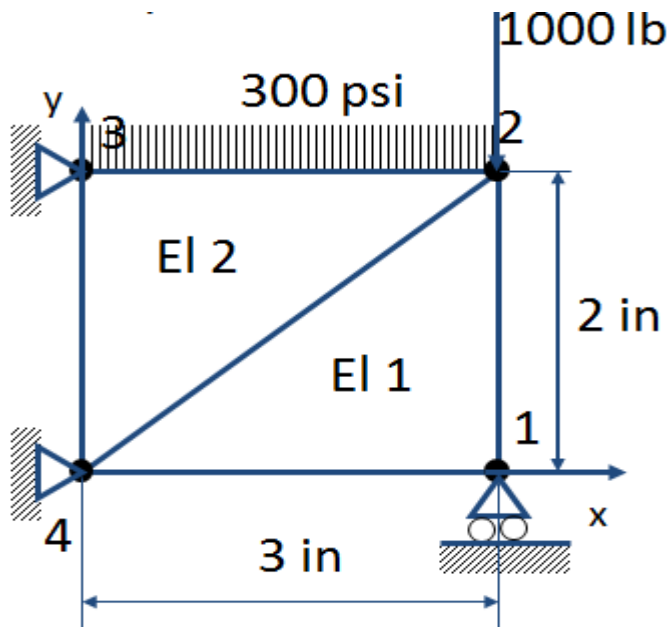
$$f_{S_{3x}} = t$$

$$f_{S_{3y}} = 0$$

Recommendations for use of CST

- Use in areas where strain gradients are small
- Use in mesh transition areas (fine mesh to coarse mesh)
- Avoid CST in critical areas of structures (e.g., stress concentrations, edges of holes, corners)
- In general CSTs are not recommended for general analysis purposes as a very large number of these elements are required for reasonable accuracy.

Example



Thickness (t) = 0.5 in
 $E = 30 \times 10^6$ psi
 $\nu = 0.25$

- Compute the unknown nodal displacements.
- Compute the stresses in the two elements.

Realize that this is a plane stress problem and therefore we need to use

$$\underline{D} = \frac{E}{1-\nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{bmatrix} = \begin{bmatrix} 3.2 & 0.8 & 0 \\ 0.8 & 3.2 & 0 \\ 0 & 0 & 1.2 \end{bmatrix} \times 10^7 \text{ psi}$$

Step 1: Node-element connectivity chart

ELEMEN T	Node 1	Node 2	Node 3	Area (sqin)
1	1	2	4	3
2	3	4	2	3

Node	x	y
1	3	0
2	3	2
3	0	2
4	0	0

Nodal coordinates

Step 2: Compute strain-displacement matrices for the elements

Recall

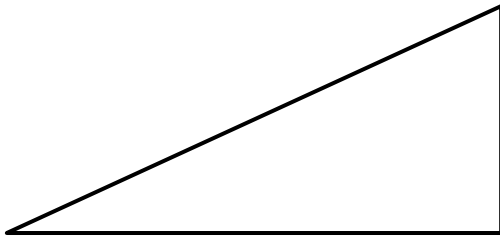
$$\underline{B} = \frac{1}{2A} \begin{bmatrix} b_1 & 0 & b_2 & 0 & b_3 & 0 \\ 0 & c_1 & 0 & c_2 & 0 & c_3 \\ c_1 & b_1 & c_2 & b_2 & c_3 & b_3 \end{bmatrix}$$

with

$$\begin{aligned} b_1 &= y_2 - y_3 & b_2 &= y_3 - y_1 & b_3 &= y_1 - y_2 \\ c_1 &= x_3 - x_2 & c_2 &= x_1 - x_3 & c_3 &= x_2 - x_1 \end{aligned}$$

For Element #1:

2(2)



$$y_1 = 0; y_2 = 2; y_3 = 0$$

$$x_1 = 3; x_2 = 3; x_3 = 0$$

Hence

$$b_1 = 2 \quad b_2 = 0 \quad b_3 = -2$$

$$c_1 = -3 \quad c_2 = 3 \quad c_3 = 0$$

4(3)

1(1)

(local numbers within brackets)

Therefore

$$\underline{B}^{(1)} = \frac{1}{6} \begin{bmatrix} 2 & 0 & 0 & 0 & -2 & 0 \\ 0 & -3 & 0 & 3 & 0 & 0 \\ -3 & 2 & 3 & 0 & 0 & -2 \end{bmatrix}$$

For Element #2:

$$\underline{B}^{(2)} = \frac{1}{6} \begin{bmatrix} -2 & 0 & 0 & 0 & 2 & 0 \\ 0 & 3 & 0 & -3 & 0 & 0 \\ 3 & -2 & -3 & 0 & 0 & 2 \end{bmatrix}$$

Step 3: Compute element stiffness matrices

$$\underline{k}^{(1)} = A t \underline{B}^{(1)T} \underline{D} \underline{B}^{(1)} = (3)(0.5) \underline{B}^{(1)T} \underline{D} \underline{B}^{(1)}$$

$$= \begin{bmatrix} 0.9833 & -0.5 & -0.45 & 0.2 & -0.5333 & 0.3 \\ & 1.4 & 0.3 & -1.2 & 0.2 & -0.2 \\ & & 0.45 & 0 & 0 & -0.3 \\ & & & 1.2 & -0.2 & 0 \\ & & & & 0.5333 & 0 \\ & & & & & 0.2 \end{bmatrix} \times 10^7$$

u_1 v_1 u_2 v_2 u_4 v_4

$$\underline{k}^{(2)} = A t \underline{B}^{(2)T} \underline{D} \underline{B}^{(2)} = (3)(0.5) \underline{B}^{(2)T} \underline{D} \underline{B}^{(2)}$$

$$= \begin{bmatrix} 0.9833 & -0.5 & -0.45 & 0.2 & -0.5333 & 0.3 \\ & 1.4 & 0.3 & -1.2 & 0.2 & -0.2 \\ & & 0.45 & 0 & 0 & -0.3 \\ & & & 1.2 & -0.2 & 0 \\ & & & & 0.5333 & 0 \\ & & & & & 0.2 \end{bmatrix} \times 10^7$$

u_3	v_3	u_4	v_4	u_2	v_2
-------	-------	-------	-------	-------	-------

Step 4: Assemble the global stiffness matrix corresponding to the nonzero degrees of freedom

$$u_3 = v_3 = u_4 = v_4 = v_1 = 0$$

Hence we need to calculate only a small (3x3) stiffness matrix

$$\underline{K} = \begin{bmatrix} 0.983 & -0.45 & 0.2 \\ -0.45 & 0.983 & 0 \\ 0.2 & 0 & 1.4 \end{bmatrix} \times 10^7 \begin{matrix} u \\ u_2 \\ v_2 \end{matrix}$$

u_1
 u_2
 v_2

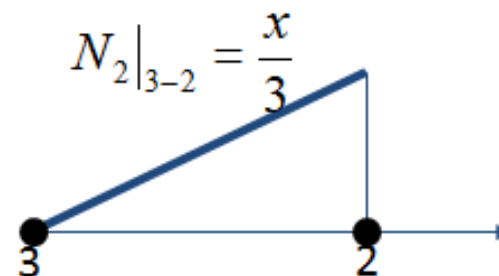
Step 5: Compute consistent nodal loads

$$\underline{f} = \begin{Bmatrix} f_{1x} \\ f_{2x} \\ f_{2y} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ f_{2y} \end{Bmatrix}$$

$$f_{2y} = -1000 + f_{s_{2y}}$$

The consistent nodal load due to traction on the edge 3-2

$$\begin{aligned} f_{s_{2y}} &= \int_{x=0}^3 N_3|_{3-2} (-300) t dx \\ &= (-300)(0.5) \int_{x=0}^3 N_3|_{3-2} dx \\ &= -150 \int_{x=0}^3 \frac{x}{3} dx \\ &= -50 \left[\frac{x^2}{2} \right]_0^3 = -50 \left(\frac{9}{2} \right) = -225 \text{ lb} \end{aligned}$$



Hence

$$\begin{aligned} f_{2y} &= -1000 + f_{s_{2y}} \\ &= -1225 \text{ lb} \end{aligned}$$

Step 6: Solve the system equations to obtain the unknown nodal loads

$$\underline{K} \underline{d} = \underline{f}$$

$$10^7 \times \begin{bmatrix} 0.983 & -0.45 & 0.2 \\ -0.45 & 0.983 & 0 \\ 0.2 & 0 & 1.4 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ v_2 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ -1225 \end{Bmatrix}$$

Solve to get

$$\begin{Bmatrix} u_1 \\ u_2 \\ v_2 \end{Bmatrix} = \begin{Bmatrix} 0.2337 \times 10^{-4} \text{ in} \\ 0.1069 \times 10^{-4} \text{ in} \\ -0.9084 \times 10^{-4} \text{ in} \end{Bmatrix}$$

Step 7: Compute the stresses in the elements

In Element #1

$$\underline{\sigma}^{(1)} = \underline{D} \underline{B}^{(1)} \underline{d}^{(1)}$$

With

$$\begin{aligned} \underline{d}^{(1)T} &= [u_1 \quad v_1 \quad u_2 \quad v_2 \quad u_4 \quad v_4] \\ &= [0.2337 \times 10^{-4} \quad 0 \quad 0.1069 \times 10^{-4} \quad -0.9084 \times 10^{-4} \quad 0 \quad 0] \end{aligned}$$

Calculate

$$\underline{\sigma}^{(1)} = \begin{bmatrix} -114.1 \\ -1391.1 \\ -76.1 \end{bmatrix} \text{ psi}$$

In Element #2

$$\underline{\sigma}^{(2)} = \underline{D} \underline{B}^{(2)} \underline{d}^{(2)}$$

With

$$\begin{aligned} \underline{d}^{(2)T} &= [u_3 \quad v_3 \quad u_4 \quad v_4 \quad u_2 \quad v_2] \\ &= [0 \quad 0 \quad 0 \quad 0 \quad 0.1069 \times 10^{-4} \quad -0.9084 \times 10^{-4}] \end{aligned}$$

Calculate

$$\underline{\sigma}^{(2)} = \begin{bmatrix} 114.1 \\ 28.52 \\ -363.35 \end{bmatrix} \text{ psi}$$

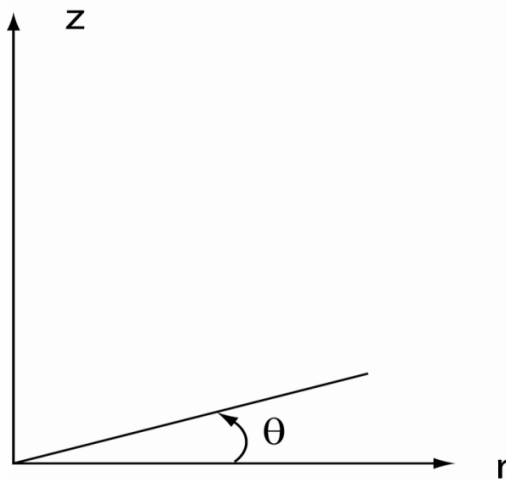
Notice that the stresses are constant in each element

Definition:

A problem in which geometry, loadings, boundary conditions and materials are symmetric about one axis.

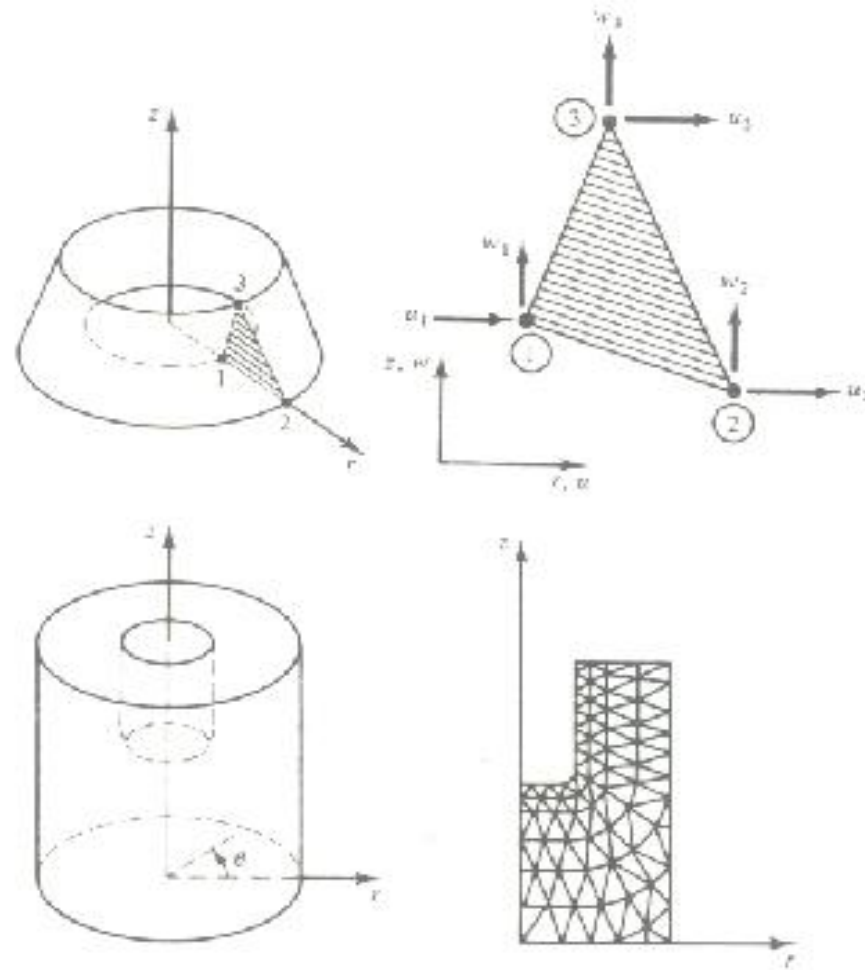
Axi-symmetric Analysis

$$x = r \cos \theta; \quad y = r \sin \theta; \quad z = z$$



- quantities depend on r and z only
- 3-D problem
- 2-D problem

Axi-symmetric Analysis



Axi-symmetric Analysis – Single-Variable Problem

$$-\frac{1}{r} \frac{\partial}{\partial r} \left(r a_{11} \frac{\partial u(r, z)}{\partial r} \right) - \frac{\partial}{\partial z} \left(a_{22} \frac{\partial u(r, z)}{\partial z} \right) + a_{00} u - f(r, z) = 0$$

Weak form:

$$0 = \int_{\Omega_e} \left[\frac{\partial w}{\partial r} \left(a_{11} \frac{\partial u}{\partial r} \right) + \frac{\partial w}{\partial z} \left(a_{22} \frac{\partial u}{\partial z} \right) + a_{00} w u - w f(r, z) \right] r dr dz$$

$$- \oint_{\Gamma_e} w q_n ds$$

where

$$q_n = a_{11} \frac{\partial u(r, z)}{\partial r} n_r + a_{22} \frac{\partial u(r, z)}{\partial z} n_z$$

Finite Element Model – Single-Variable Problem

$$u = \sum_j u_j \phi_j \quad \text{where} \quad \phi_j(r, z) = \phi_j(x, y)$$

Ritz method: $w = \phi_i$

Weak form $\longrightarrow \sum_{j=1}^n K_{ij}^e u_j^e = f_i^e + Q_i^e$

where
$$K_{ij}^e = \int_{\Omega_e} \left(a_{11} \frac{\partial \phi_i}{\partial r} \frac{\partial \phi_j}{\partial r} + a_{22} \frac{\partial \phi_i}{\partial z} \frac{\partial \phi_j}{\partial z} + a_{00} \phi_i \phi_j \right) r dr dz$$

$$f_i^e = \int_{\Omega_e} \phi_i f r dr dz$$

$$Q_i^e = \int_{\Gamma_e} \phi_i q_n ds$$

Single-Variable Problem – Heat Transfer

Heat Transfer:

$$-\frac{1}{r} \frac{\partial}{\partial r} \left(rk \frac{\partial T(r, z)}{\partial r} \right) - \frac{\partial}{\partial z} \left(k \frac{\partial T(r, z)}{\partial z} \right) - f(r, z) = 0$$

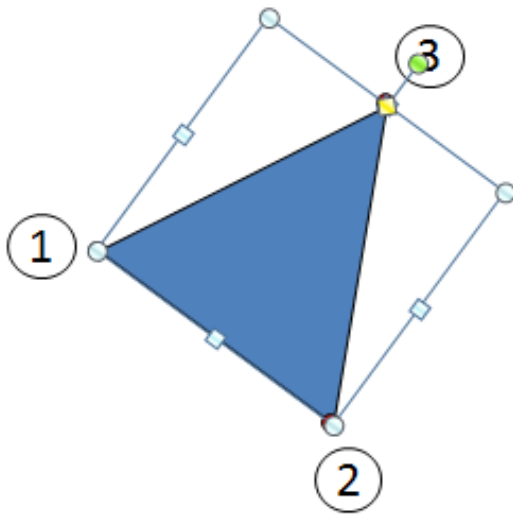
Weak form

$$0 = \int_{\Omega_e} \left[\frac{\partial w}{\partial r} \left(k \frac{\partial T}{\partial r} \right) + \frac{\partial w}{\partial z} \left(k \frac{\partial T}{\partial z} \right) - wf(r, z) \right] r dr dz$$
$$- \oint_{\Gamma_e} w q_n ds$$

where

$$q_n = k \frac{\partial T(r, z)}{\partial r} n_r + k \frac{\partial T(r, z)}{\partial z} n_z$$

3-Node Axi-symmetric Element



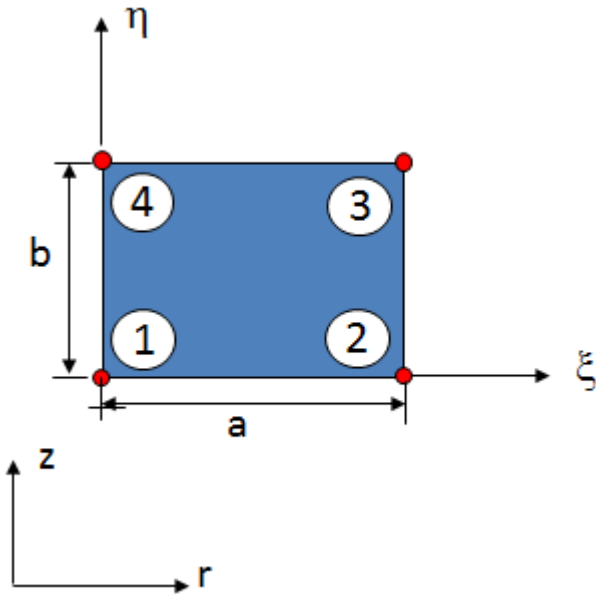
$$T(r, z) = T_1\phi_1 + T_2\phi_2 + T_3\phi_3$$

$$\phi_1 = \frac{\begin{pmatrix} 1 & r & z \end{pmatrix}}{2A_e} \begin{Bmatrix} r_2z_3 - r_3z_2 \\ z_2 - z_3 \\ r_3 - r_2 \end{Bmatrix}$$

$$\phi_2 = \frac{\begin{pmatrix} 1 & r & z \end{pmatrix}}{2A_e} \begin{Bmatrix} r_3z_1 - r_1z_3 \\ z_3 - z_1 \\ r_1 - r_3 \end{Bmatrix}$$

$$\phi_3 = \frac{\begin{pmatrix} 1 & r & z \end{pmatrix}}{2A_e} \begin{Bmatrix} r_1z_2 - r_2z_1 \\ z_1 - z_2 \\ r_2 - r_1 \end{Bmatrix}$$

4-Node Axi-symmetric Element

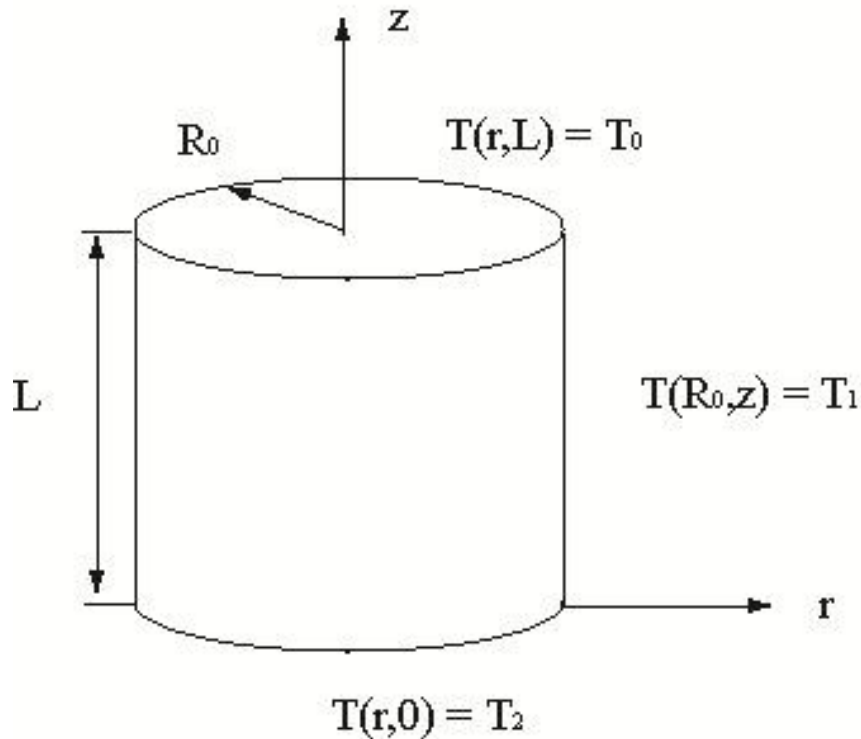


$$T(r, z) = T_1\phi_1 + T_2\phi_2 + T_3\phi_3 + T_4\phi_4$$

$$\phi_1 = \left(1 - \frac{\xi}{a}\right) \left(1 - \frac{\eta}{b}\right) \quad \phi_2 = \frac{\xi}{a} \left(1 - \frac{\eta}{b}\right)$$

$$\phi_3 = \frac{\xi}{a} \frac{\eta}{b} \quad \phi_4 = \left(1 - \frac{\xi}{a}\right) \frac{\eta}{b}$$

Single-Variable Problem – Example



Step 1: Discretization

Step 2: Element equation

$$K_{ij}^e = \int_{\Omega_e} \left(\kappa \frac{\partial \phi_i}{\partial r} \frac{\partial \phi_j}{\partial r} + \kappa \frac{\partial \phi_i}{\partial z} \frac{\partial \phi_j}{\partial z} \right) r dr dz$$

$$f_i^e = \int_{\Omega_e} \phi_i f r dr dz \quad Q_i^e = \int_{\Gamma_e} \phi_i q_n ds$$

Review of CST Element

Constant Strain Triangle (CST) - easiest and simplest finite element

- Displacement field in terms of generalized coordinates

$$u = \beta_1 + \beta_2 x + \beta_3 y$$

$$v = \beta_4 + \beta_5 x + \beta_6 y$$

- Resulting strain field is

$$\epsilon_x = \beta_2 \quad \epsilon_y = \beta_6 \quad \gamma_{xy} = \beta_3 + \beta_5$$

- Strains do not vary within the element. Hence, the name constant strain triangle (CST)
 - Other elements are not so lucky.
 - Can also be called linear triangle because displacement field is linear in x and y - sides remain straight.

Constant Strain Triangle

- The strain field from the shape functions looks like:

$$\begin{Bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{Bmatrix} = \frac{1}{2A} \begin{bmatrix} y_{23} & 0 & y_{31} & 0 & y_{12} & 0 \\ 0 & x_{32} & 0 & x_{13} & 0 & x_{21} \\ x_{32} & y_{23} & x_{13} & y_{31} & x_{21} & y_{12} \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}$$

- Where, x_i and y_i are nodal coordinates ($i=1, 2, 3$)
 - $x_{ij} = x_i - x_j$ and $y_{ij} = y_i - y_j$
 - $2A$ is twice the area of the triangle, $2A = x_{21}y_{31} - x_{31}y_{21}$
- Node numbering is arbitrary except that the sequence 123 must go clockwise around the element if A is to be positive.

Constant Strain Triangle

- Stiffness matrix for element $k = B^T E B \, tA$
- The CST gives good results in regions of the FE model where there is little strain gradient
 - Otherwise it does not work well.

Linear Strain Triangle

Changes the shape functions and results in quadratic displacement distributions and linear strain distributions within the element.

$$u = \beta_1 + \beta_2 x + \beta_3 y + \beta_4 x^2 + \beta_5 xy + \beta_6 y^2$$

$$v = \beta_7 + \beta_8 x + \beta_9 y + \beta_{10} x^2 + \beta_{11} xy + \beta_{12} y^2$$

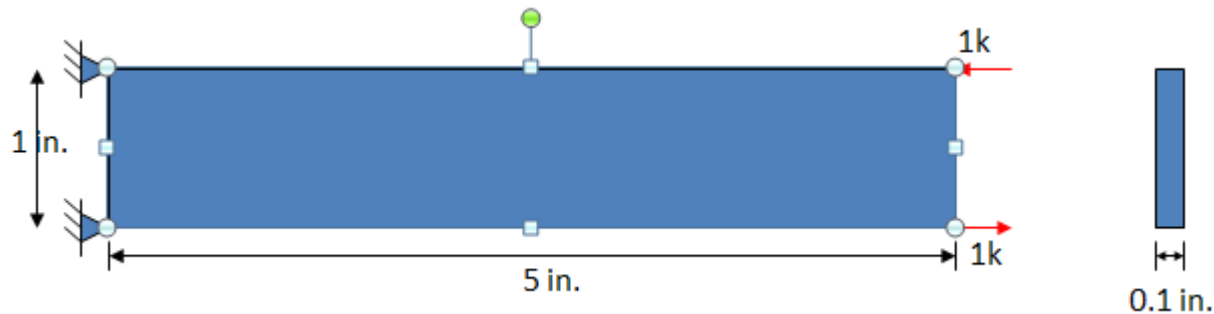
$$\varepsilon_x = \beta_2 + 2\beta_4 x + \beta_5 y$$

$$\varepsilon_y = \beta_9 + \beta_{11} x + 2\beta_{12} y$$

$$\gamma_{xy} = (\beta_3 + \beta_8) + (\beta_5 + 2\beta_{10})x + (2\beta_6 + \beta_{11})y$$

Example Problem

Consider the problem we were looking at:



$$I = 0.1 \times 1^3 / 12 = 0.008333 \text{ in}^4$$

$$\sigma = \frac{M \times c}{I} = \frac{1 \times 0.5}{0.008333} = 60 \text{ ksi}$$

$$\varepsilon = \frac{\sigma}{E} = 0.00207$$

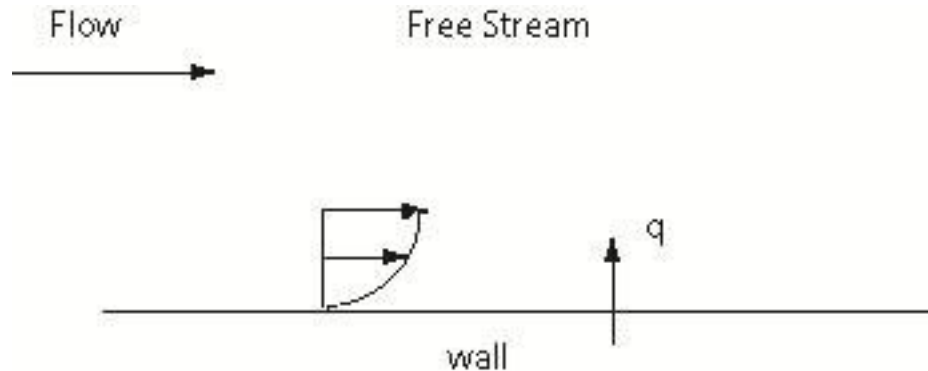
$$\delta = \frac{ML^2}{2EI} = \frac{25}{2 \times 29000 \times 0.008333} = 0.0517 \text{ in.}$$

UNIT-4

STEADY STATE HEAT TRANSFER ANALYSIS

CLOs	Course Learning Outcomes
CLO 1	Understand the concepts of steady state heat transfer analysis for one dimensional slab, fin and thin plate.
CLO 2	Derive the stiffness matrix for fin element.
CLO 3	Solve the steady state heat transfer problems for fin and composite slab.

Thermal Convection



Newton's Law of Cooling

$$q = h(T_s - T_\infty)$$

h : convective heat transfer coefficient ($W/m^2 \cdot C^\circ$)

Boundary conditions:

- Dirichlet BC
- Natural BC
- Mixed BC

Weak Formulation of 1-D Heat Conduction (Steady State Analysis)

Governing Equation of 1-D Heat Conduction

$$-\frac{d}{dx}\left(\kappa(x)A(x)\frac{dT(x)}{dx}\right) - AQ(x) = 0 \quad 0 < x < L$$

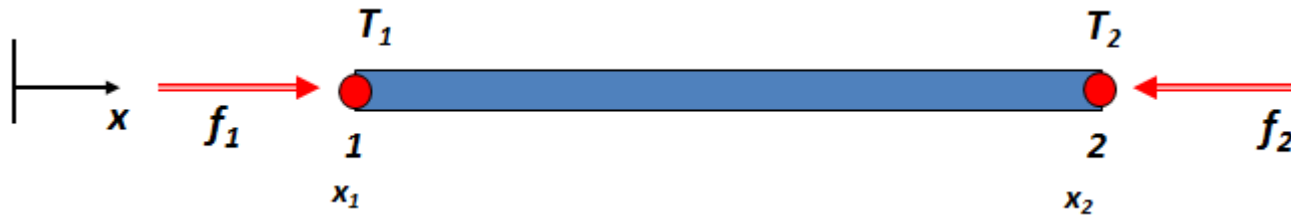
Weighted Integral Formulation

$$0 = \int_0^L w(x) \left[-\frac{d}{dx}\left(\kappa(x)A(x)\frac{dT(x)}{dx}\right) - AQ(x) \right] dx$$

Weak Form from Integration-by-Parts

$$0 = \int_0^L \left[\frac{dw}{dx} \left(\kappa A \frac{dT}{dx} \right) - w A Q \right] dx - w \left(\kappa A \frac{dT}{dx} \right) \Big|_0^L$$

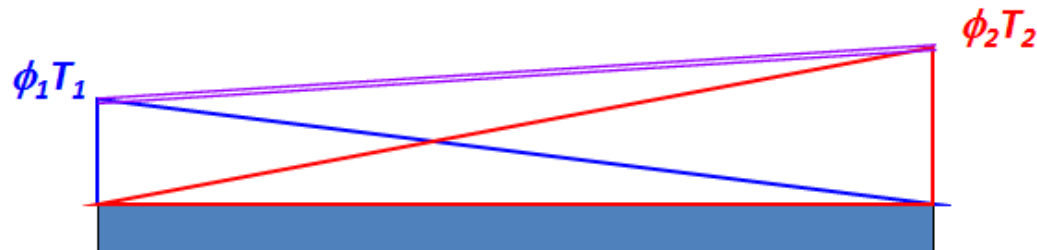
Formulation for 1-D Linear Element



$$f_1(x) = -\kappa A \frac{\partial T}{\partial x} \bigg|_1, \quad f_2(x) = \kappa A \frac{\partial T}{\partial x} \bigg|_2$$

$$T(x) = T_1 \phi_1(x) + T_2 \phi_2(x)$$

$$\phi_1(x) = \frac{x_2 - x}{l}, \quad \phi_2(x) = \frac{x - x_1}{l}$$



Formulation for 1-D Linear Element

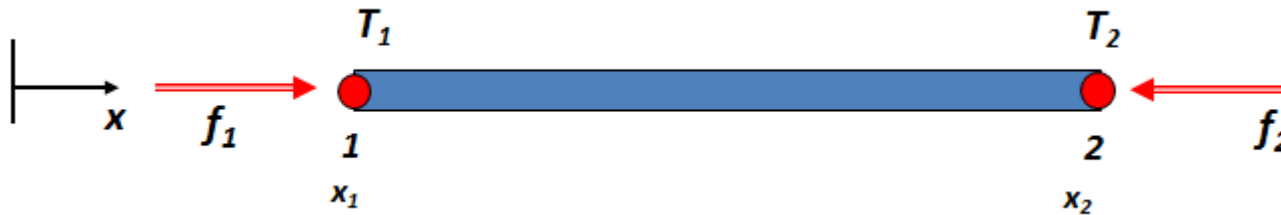
- Let $w(x) = f_i(x)$, $i = 1, 2$

$$\begin{aligned} 0 &= \sum_{j=1}^2 T_j \left[\int_{x_1}^{x_2} \kappa A \left(\frac{d\phi_i}{dx} \frac{d\phi_j}{dx} \right) dx \right] - \int_{x_1}^{x_2} (\phi_i A Q) dx - [\phi_i(x_2) f_2 + \phi_i(x_1) f_1] \\ &= \sum_{j=1}^2 K_{ij} T_j - Q_i - [\phi_i(x_2) f_2 + \phi_i(x_1) f_1] \end{aligned}$$

$$\begin{Bmatrix} f_1 \\ f_2 \end{Bmatrix} + \begin{Bmatrix} Q_1 \\ Q_2 \end{Bmatrix} = \begin{bmatrix} K_{11} & K_{12} \\ K_{12} & K_{22} \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix}$$

$$\text{where } K_{ij} = \int_{x_1}^{x_2} \kappa A \left(\frac{d\phi_i}{dx} \frac{d\phi_j}{dx} \right) dx, \quad Q_i = \int_{x_1}^{x_2} (\phi_i A Q) dx, \quad f_1 = -\kappa A \frac{dT}{dx} \Big|_{x_1}, \quad f_2 = \kappa A \frac{dT}{dx} \Big|_{x_2}$$

Element Equations of 1-D Linear Element

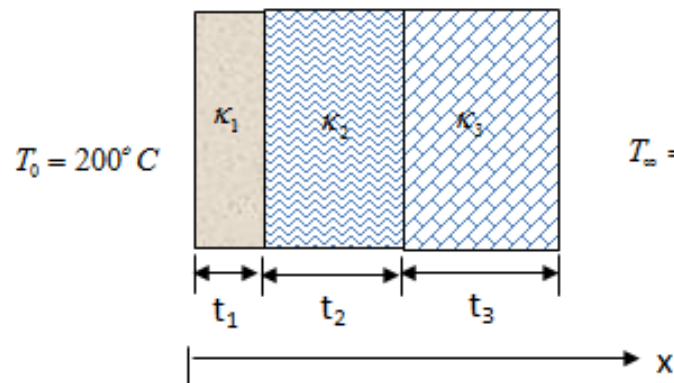


$$\begin{Bmatrix} f_1 \\ f_2 \end{Bmatrix} + \begin{Bmatrix} Q_1 \\ Q_2 \end{Bmatrix} = \frac{\kappa A}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix}$$

$$\text{where } Q_i = \int_{x_1}^{x_2} (\phi_i A Q) dx, \quad f_1 = -\kappa A \left. \frac{dT}{dx} \right|_{x=x_1}, \quad f_2 = \kappa A \left. \frac{dT}{dx} \right|_{x=x_2}$$

1-D Heat Conduction - Example

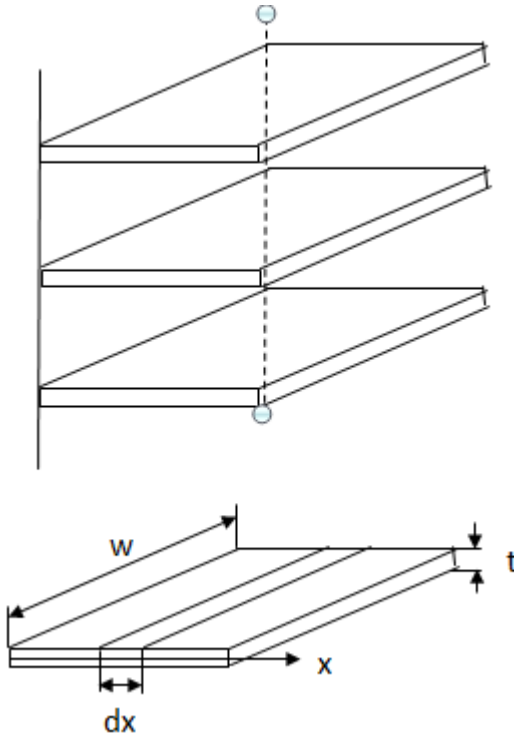
A composite wall consists of three materials, as shown in the figure below. The inside wall temperature is 200°C and the outside air temperature is 50°C with a convection coefficient of $h = 10 \text{ W}(\text{m}^2.\text{K})$. Find the temperature along the composite wall.



$$\kappa_1 = 70 \text{ W}/(\text{m} \cdot \text{K}), \quad \kappa_2 = 40 \text{ W}/(\text{m} \cdot \text{K}), \quad \kappa_3 = 20 \text{ W}/(\text{m} \cdot \text{K})$$

$$t_1 = 2 \text{ cm}, \quad t_2 = 2.5 \text{ cm}, \quad t_3 = 4 \text{ cm}$$

Thermal Conduction and Convection- Fin



Objective: to enhance heat transfer

Governing equation for 1-D heat transfer in thin fin

$$\frac{d}{dx} \left(\kappa A_c \frac{dT}{dx} \right) + A_c Q = 0$$

$$Q_{loss} = \frac{2h(T - T_\infty) \cdot dx \cdot w + 2h(T - T_\infty) \cdot dx \cdot t}{A_c \cdot dx} = \frac{2h(T - T_\infty) \cdot (w + t)}{A_c}$$

$$\frac{d}{dx} \left(\kappa A_c \frac{dT}{dx} \right) - Ph(T - T_\infty) + A_c Q = 0$$

where $P = 2(w + t)$

Fin (*Steady State Analysis*)

Governing Equation of 1-D Heat Conduction

$$-\frac{d}{dx}\left(\kappa(x)A(x)\frac{dT(x)}{dx}\right) + Ph(T - T_{\infty}) - AQ = 0 \quad 0 < x < L$$

Weighted Integral Formulation

$$0 = \int_0^L w(x) \left[-\frac{d}{dx}\left(\kappa(x)A(x)\frac{dT(x)}{dx}\right) + Ph(T - T_{\infty}) - AQ(x) \right] dx$$

Weak Form from Integration-by-Parts

$$0 = \int_0^L \left[\frac{dw}{dx} \left(\kappa A \frac{dT}{dx} \right) + wPh(T - T_{\infty}) - wAQ \right] dx - w \left(\kappa A \frac{dT}{dx} \right) \Big|_0^L$$

Formulation for 1-D Linear Element

Let $w(x) = f_i(x)$, $i = 1, 2$

$$0 = \sum_{j=1}^2 T_j \left[\int_{x_1}^{x_2} \left(\kappa A \frac{d\phi_i}{dx} \frac{d\phi_j}{dx} + Ph\phi_i\phi_j \right) dx \right] - \int_{x_1}^{x_2} \phi_i (AQ + PhT_\infty) dx$$

$$- [\phi_i(x_2) f_2 + \phi_i(x_1) f_1]$$

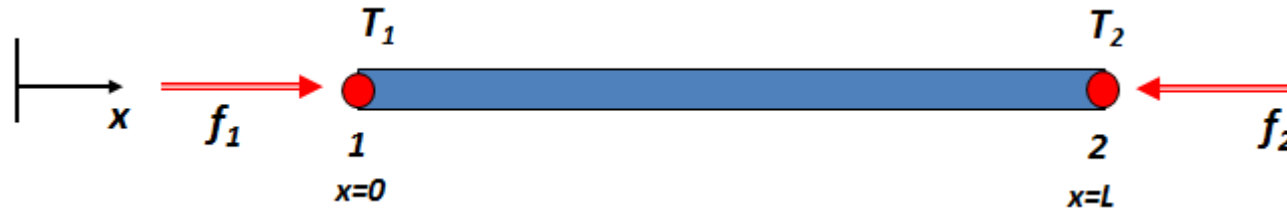
$$= \sum_{j=1}^2 K_{ij} T_j - Q_i - [\phi_i(x_2) f_2 + \phi_i(x_1) f_1]$$

$$\begin{Bmatrix} f_1 \\ f_2 \end{Bmatrix} + \begin{Bmatrix} Q_1 \\ Q_2 \end{Bmatrix} = \begin{bmatrix} K_{11} & K_{12} \\ K_{12} & K_{22} \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix}$$

where $K_{ij} = \int_{x_1}^{x_2} \left(\kappa A \frac{d\phi_i}{dx} \frac{d\phi_j}{dx} + Ph\phi_i\phi_j \right) dx$, $Q_i = \int_{x_1}^{x_2} \phi_i (AQ + PhT_\infty) dx$,

$$f_1 = -\kappa A \frac{dT}{dx} \Big|_{x=x_1}, \quad f_2 = \kappa A \frac{dT}{dx} \Big|_{x=x_2}$$

Element Equations of 1-D Linear Element



$$\begin{Bmatrix} f_1 \\ f_2 \end{Bmatrix} + \begin{Bmatrix} Q_1 \\ Q_2 \end{Bmatrix} = \left\{ \frac{\kappa A}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} + \frac{Phl}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \right\} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix}$$

where $Q_i = \int_{x_1}^{x_2} \phi_i (AQ + PhT_\infty) dx$, $f_1 = -\kappa A \frac{dT}{dx} \Big|_{x=x_1}$, $f_2 = \kappa A \frac{dT}{dx} \Big|_{x=x_2}$

Time-Dependent Problems

In general

$$u(x, t)$$

Two approaches:

$$u(x, t) = \sum u_j \phi_j(x, t)$$

$$u(x, t) = \sum u_j(t) \phi_j(x)$$

Model Problem I – Transient Heat Conduction

$$c \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(a \frac{\partial u}{\partial x} \right) + f(x, t)$$

Weak form:

$$0 = \int_{x_1}^{x_2} \left(a \frac{\partial w}{\partial x} \frac{\partial u}{\partial x} + cw \frac{\partial u}{\partial t} - wf \right) dx - Q_1 w(x_1) - Q_2 w(x_2)$$

$$Q_1 = - \left[a \frac{du}{dx} \right]_{x_1} ; \quad Q_2 = \left[a \frac{du}{dx} \right]_{x_2}$$

Transient Heat Conduction

let: $u(x, t) = \sum_{j=1}^n u_j(t) \phi_j(x)$ and $w = \phi_i(x)$



$$0 = \int_{x_1}^{x_2} \left(a \frac{\partial w}{\partial x} \frac{\partial u}{\partial x} + cw \frac{\partial u}{\partial t} - wf \right) dx - Q_1 w(x_1) - Q_2 w(x_2)$$



$$[K]\{u\} + [M]\{\dot{u}\} = \{F\}$$

$$K_{ij} = \int_{x_1}^{x_2} a \frac{\partial \phi_i}{\partial x} \frac{\partial \phi_j}{\partial x} dx$$

$$M_{ij} = \int_{x_1}^{x_2} c \phi_i \phi_j dx$$

$$F_i = \int_{x_1}^{x_2} \phi_i f dx + Q_i$$

Time Approximation – First Order ODE

$$a \frac{du}{dt} + bu = f(t) \quad 0 < t < T \quad u(0) = u_0$$

Forward difference approximation – explicit

$$u_{k+1} = u_k + \frac{\Delta t}{a} [f_k - bu_k]$$

Backward difference approximation - implicit

$$u_{k+1} = u_k + \frac{\Delta t}{a + b\Delta t} [f_k - bu_k]$$

Stability Requirement

$$\Delta t \leq \Delta t_{cri} = \frac{2}{(1-2\alpha)\lambda_{\max}}$$

where
$$([K] - \lambda[M])\{u\} = \{Q\}$$

Note: One must use the same discretization for solving the eigenvalue problem.

Transient Heat Conduction - Example

$$\frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = 0 \quad 0 < x < 1$$

$$u(0, t) = 0 \quad \frac{\partial u}{\partial t}(1, t) = 0 \quad t > 0$$

$$u(x, 0) = 1.0$$

UNIT-5

DYNAMIC ANALYSIS

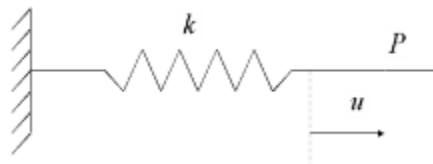
CLOs	Course Learning Outcomes
CLO 1	Understand the concepts of mass and spring system and derive the equations for various structural problems
CLO 2	Understand the concept of dynamic analysis for all types of elements.
CLO 3	Calculate the mass matrices, Eigen values, Eigen vectors, natural frequency and mode shapes for dynamic problems.

For many structural system, the approximation of linear structural behavior is made in order to convert the physical equilibrium statement, to the following set of second order linear differential equation

$$M\ddot{u}(t)_a + C\dot{u}(t)_a + Ku(t)_a = \mathbf{P}(t)$$

Stiffness and flexibility stiffness matrix

Consider a uniform elastic spring subjected to a load P . This Structure obeys Hook's law. If a force P is applied to a spring fixed. At one end, to produce a displacement then the linear force displacement is u .



$$P = ku$$

$$K = fP$$

Stiffness and flexibility stiffness matrix

$$P = ku$$

$$K = fP$$

- K is called the stiffness of the spring
- f is called the flexibility of spring

Suppose the uniform elastic spring has nodal points 1 and 2 at its ends, and that the forces at these points are P_1 and P_2 with corresponding displacements u_1 and u_2 .

➤ Elemental spring



➤ From equilibrium considerations

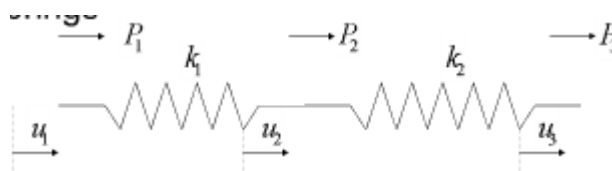
$$P_1 = k(u_1 - u_2)$$

$$P_2 = -P_1 = k(u_2 - u_1)$$

➤ It is convenient to show the above in matrix form as follows

$$\begin{Bmatrix} P_1 \\ P_2 \end{Bmatrix} = \begin{bmatrix} k & -k \\ -k & k \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix}$$

➤ Simple system consisting of just two springs



The system is in equilibrium

$$P_1 + P_2 + P_3 = 0$$

$$P_1 = k_1(u_1 - u_2)$$

$$P_2 = k_2(u_3 - u_2)$$

$$P_2 = -k_1 u_1 + (k_1 + k_2) u_2 - k_2 u_3$$

The equations written in matrix form

$$\begin{Bmatrix} P_1 \\ P_2 \\ P_3 \end{Bmatrix} = \begin{bmatrix} k_1 & -k_1 & 0 \\ -k_1 & (k_1 + k_2) & -k_2 \\ 0 & -k_2 & k_2 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix}$$

$$\{P\} = [K] \{u\}$$

- P vector of external nodal loads acting on the structure
- K structural stiffness matrix
- u – overall nodal displacement vector

Mass Matrices

- The elemental mass matrix which is always symmetrical, is a matrix of equivalent nodal masses that dynamically represent the actual distributed mass of the element.
- The element mass matrix is defined as

$$[M] = \int_v \rho [N]^T [N] dV$$

Dynamic equations

- The force equilibrium of a multi degree of freedom lumped mass system

$$P(t)_i + P(t)_D + P(t)_s = P(t)$$

- Vector of inertia forces acting on the node masses $P(t)_i$
 - Vector of viscous damping or energy dissipation forces $P(t)_D$
 - A vector of internal forces carried by the structure $P(t)_s$
 - Vector of externally applied loads $P(t)$
- For many structural systems the approximation of linear structural behavior is made in order to convert the physical equilibrium statement, to the following set of second order linear differential equation

$$M\ddot{u}(t)_a + C\dot{u}(t)_a + Ku(t)_a = P(t)$$

When loads are suddenly applied or when the loads are of a variable nature, the mass and acceleration effects come into the picture. If a solid such as an engineering structure is deformed elastically and suddenly released. It tends to vibrate about its equilibrium position. This periodic motion due to the restoring strain energy is called free vibration. The number of cycles per unit time is called frequency. The maximum displacement from the equilibrium position is the amplitude.

- Equation for damped forced vibration

$$M\ddot{u}(t) + C\dot{u}(t) + Ku(t) = P(t)$$

- If there is no damping the equation become

$$M\ddot{u}(t) + Ku(t) = P(t)$$

- Free undamped vibration equation $M\ddot{u}(t) + Ku(t) = 0$

- The free undamped vibration equation is linear and homogeneous. Its general solution is a linear combination of exponentials. Under matrix definiteness conditions the exponentials can be expressed as a combination of trigonometric functions: sines and cosines of argument ωt .
- A compact representation of such functions is obtained by using the exponential form $e^{j\omega t}$

$$u(t) = \sum v_i e^{j\omega t}$$

$$\text{Replace } u(t) = v_i e^{j\omega t}$$

$$M\ddot{u}(t) + Ku(t) = 0$$

The Vibration Eigen problem

- The time dependence to the exponential is segregated

$$(-\omega^2 M + K)v e^{j\omega t} = 0$$

- Since is not identically zero, it can be dropped leaving the algebraic condition

$$(-\omega^2 M + K)v = 0$$

- Because v cannot be the null vector this equation is an algebraic Eigen value problem in ω^2 . The Eigen values $\lambda_i = \omega_i^2$ are the roots of the characteristic polynomial be index by I

$$\det(K - \omega_i^2 M) = 0$$

- Dropping the index I this Eigen problem is usually written as

$$Kv = \omega^2 Mv$$