ELECTRICAL CIRCUITS

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How electricity flows?

Electricity Basics

Electricity starts with **electrons.** Every atom contains one or more electrons. Electrons have a **negative charge**.

Simplest model of an atom



Electrical Circuits

Whether you are using a battery, a fuel cell or a solar cell to produce electricity, there are three things that are always the same:

•The source of electricity will have **two terminals**: a positive terminal and a negative terminal.



Basic Electrical Circuits







Voltage

(Pressure) (Electromotive Force)

Voltage Pushes the electrons

How does electricity flow? What causes electrons to move from atom to atom?





How does electricity flow? What causes electrons to move from atom to atom?

Electron Flow is measured in Amps

The flow of the electrons is referred to as Current










































































































Loop (Mesh) Analysis

Loop Analysis

- Nodal analysis was developed by applying KCL at each non-reference node.
- Loop analysis is developed by applying KVL around loops in the circuit.
- Loop (mesh) analysis results in a system of linear equations which must be solved for unknown currents.

Another Summing Circuit

• The output voltage V of this circuit is proportional to the sum of the two input voltages V_1 and V_2 .



Steps of Mesh Analysis

1. Identify mesh (loops).

- 2. Assign a current to each mesh.
- 3. Apply KVL around each loop to get an equation in terms of the loop currents.
- 4. Solve the resulting system of linear equations.

1. Identifying the Meshes



Steps of Mesh Analysis

- 1. Identify mesh (loops).
- 2. Assign a current to each mesh.
- 3. Apply KVL around each loop to get an equation in terms of the loop currents.
- Solve the resulting system of linear equations.

2. Assigning Mesh Currents



Steps of Mesh Analysis

- 1. Identify mesh (loops).
- 2. Assign a current to each mesh.
- Apply KVL around each loop to get an equation in terms of the loop currents.
- 4. Solve the resulting system of linear equations.

Voltages from Mesh Currents





 $V_R = I_1 R$



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3. KVL Around Mesh 1



 $-V_1 + I_1 \ 1k\Omega + (I_1 - I_2) \ 1k\Omega = 0$ $I_1 \ 1k\Omega + (I_1 - I_2) \ 1k\Omega = V_1$
3. KVL Around Mesh 2



 $(I_2 - I_1) \ 1k\Omega + I_2 \ 1k\Omega + V_2 = 0$ $(I_2 - I_1) \ 1k\Omega + I_2 \ 1k\Omega = -V_2$

Steps of Mesh Analysis

- 1. Identify mesh (loops).
- 2. Assign a current to each mesh.
- 3. Apply KVL around each loop to get an equation in terms of the loop currents.
- 4. Solve the resulting system of linear equations.

Matrix Notation

The two equations can be combined into a single matrix/vector equation.

$$\begin{bmatrix} 1k\Omega + 1k\Omega & -1k\Omega \\ -1k\Omega & 1k\Omega + 1k\Omega \end{bmatrix} \begin{bmatrix} I_1 \\ I_2 \end{bmatrix} = \begin{bmatrix} V_1 \\ -V_2 \end{bmatrix}$$

Cramer's Rule

$$\begin{bmatrix} E_1 \\ -E_2 \end{bmatrix} = \begin{bmatrix} I_1 \\ I_2 \end{bmatrix} \begin{bmatrix} (R_1 + R_2) & (-R_3) \\ (-R_3) & (R_2 + R_3) \end{bmatrix}$$

$$I_{1} = \frac{\begin{bmatrix} E_{1} & (-R_{3}) \\ -E_{2}(R_{2}+R_{3}) \end{bmatrix}}{\begin{bmatrix} (R_{1}+R_{2}) & (-R_{3}) \\ (-R_{3}) & (R_{2}+R_{3}) \end{bmatrix}}$$

$$I_{2} = \frac{\begin{bmatrix} (R_{1} + R_{2}) & E_{1} \\ (-R_{3}) & -E_{2} \end{bmatrix}}{\begin{bmatrix} (R_{1} + R_{2}) & (-R_{3}) \\ (-R_{3}) & (R_{2} + R_{3}) \end{bmatrix}}$$

4. Solving the Equations

Let: $V_1 = 7V$ and $V_2 = 4V$ Results: $I_1 = 3.33 \text{ mA}$ $I_2 = -0.33 \text{ mA}$ Finally $V_{out} = (I_1 - I_2) 1 \text{k}\Omega = 3.66V$







Current Sources

- The current sources in this circuit will have whatever voltage is necessary to make the current correct.
- We can't use KVL around any mesh because we don't know the voltage for the current sources.
- What to do?

Current Sources

- The 4mA current source sets I_2 : $I_2 = -4 \text{ mA}$
- The 2mA current source sets a *constraint* on I_1 and I_3 :

 $I_1 - I_3 = 2 \,\mathrm{mA}$

We have two equations and three unknowns.
Where is the third equation?

Supermesh



3. KVL Around the Supermesh

$-12V + I_3 2k\Omega + (I_3 - I_2)1k\Omega + (I_1 - I_2)2k\Omega = 0$

$I_3 2k\Omega + (I_3 - I_2)1k\Omega + (I_1 - I_2)2k\Omega = 12V$

Matrix Notation

The three equations can be combined into a single matrix/vector equation.

$$\begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & -1 \\ 2k\Omega & -1k\Omega - 2k\Omega & 2k\Omega + 1k\Omega \end{bmatrix}\begin{bmatrix} I_1 \\ I_2 \\ I_3 \end{bmatrix} = \begin{bmatrix} -4mA \\ 2mA \\ 12V \end{bmatrix}$$

Advantages of Nodal Analysis

- Solves directly for node voltages.
- Current sources are easy.
- Voltage sources are either very easy or somewhat difficult.
- Works best for circuits with few nodes.
- Works for any circuit.

Advantages of Loop Analysis

- Solves directly for some currents.
- Voltage sources are easy.
- Current sources are either very easy or somewhat difficult.
- Works best for circuits with few loops.

Disadvantages of Loop Analysis

- Some currents must be computed from loop currents.
- Does not work with non-planar circuits.
- Choosing the supermesh may be difficult.

YI: Spice uses a nodal analysis approach

Magnetically Coupled Networks

- A new four-terminal element, the *transformer*, is introduced in this chapter
- A transformer is composed of two closely spaced inductors, that is, two or more magnetically coupled coils
 - *primary side* is connected to the source
 - secondary side is connected to the load

Dot Convention

- dot convention: dots are placed beside each coil (inductor) so that if the currents are entering (or leaving) both dotted terminals, then the fluxes add
- right hand rule says that curling the fingers (of the right hand) around the coil in the direction of the current gives the direction of the magnetic flux based on the direction of the thumb
- We need dots on the schematic to know how the coils are physically oriented wrt one

Mutually Coupled Coils

The following equations define the coupling between the two inductors assuming that each respective current enters the dot side which is also the positive voltagdeiside di_2 $v_1(t) = L_1 \frac{deiside}{dt} \frac{di_2}{dt}$

$$v_2(t) = M \frac{d i_1}{dt} + L_2 \frac{d i_2}{dt}$$

where L_1 and L_2 are the *self-inductances* of the coils (inductors), and M is the *mutual inductance* between the two coils

Mutually Coupled Coils



$$v_1(t) = L_1 \frac{d i_1}{dt} + M \frac{d i_2}{dt}$$
$$v_2(t) = M \frac{d i_1}{dt} + L_2 \frac{d i_2}{dt}$$

Lecture 11

Class Example

• Extension Exercise E11.1

Lecture 11

Mutually Coupled Coils (AC)

 The frequency domain model of the coupled circuit is essentially identical to that of the time domain

$$\mathbf{V}_1 = j\omega L_1 \mathbf{I}_1 + j\omega M \mathbf{I}_2$$
$$\mathbf{V}_2 = j\omega M \mathbf{I}_1 + j\omega L_2 \mathbf{I}_2$$

Class Examples

- Extension Exercise E11.2
- Extension Exercise E11.3

Source Input Impedance

The source sees an *input impedance*, Z_i , that is the sum of the *primary impedance*, and a *reflected impedance*, Z_R , due to the secondary (load) side V_S

$$\mathbf{Z}_{i} \equiv \frac{\mathbf{V}_{S}}{\mathbf{I}_{P}} = \mathbf{Z}_{P} + \mathbf{Z}_{R} = \mathbf{Z}_{P} + f(\mathbf{Z}_{L})$$



Class Example

• Extension Exercise E11.4

Lecture 11

Energy Analysis

- An energy analysis of the mutually coupled inductors provides an expression for the instantaneous stored energy $w(t) = \frac{1}{2} L_1 [i_1(t)]^2 + \frac{1}{2} L_2 [i_2(t)]^2 \pm M i_1(t) i_2(t)$
- The sign is positive (+) if currents are both entering (or leaving) the dots; sign is negative (-) if currents are otherwise

Quantifying the Coupling

- The *mutual inductance*, *M*, is in the range $0 \le M \le \sqrt{L_1 L_2}$
- The *coefficient of coupling* (*k*) between two inductors is defined as

$$0 \le \left[k \equiv \frac{M}{\sqrt{L_1 L_2}} \right] \le 1$$

- for k > 0.5, inductors are said to be tightly coupled

– for $k \le 0.5$, coils are considered to be loosely

coupled

Lecture 11

Class Example

• Extension Exercise E11.5

Lecture 11

Electromagnetic Induction

Chapter 31

Faraday's Law Induced currents Lenz's Law Induced EMF Magnetic Flux Induced Electric fields



In a closed electric circuit, a changing magnetic field will produce an electric current

Electromagnetic Induction Faraday's Law

The induced emf in a circuit is proportional to the rate of change of magnetic flux, through any surface bounded by that circuit.

$$\mathcal{E}$$
 = - d $\Phi_{\rm B}$ / dt

Faraday's Experiments



- Michael Faraday discovered induction in 1831.
- Moving the magnet induces a current I.
- Reversing the direction reverses the current.
- Moving the loop induces a current.
- The induced current is set up by an *induced FME*

Faraday's Experiments



• Changing the current in the right-hand coil induces

a current in the left-hand coil.

The induced current does not depend on the size of

the current in the right-hand coil.

The induced current depends on dI/dt.

Magnetic Flux



In the easiest case, with a constant magnetic field **B**, and a flat surface of area **A**, the magnetic flux is

• Units : 1 tesla x m² = 1 weber

Magnetic Flux



• When B is not constant, or the surface is not flat, one must

do an integral.

• Break the surface into bits <u>dA</u>. The flux through one bit is

$d \mathbf{\Phi}_{B} = \mathbf{E} \int \mathbf{B} \cdot d\mathbf{A} d\mathbf{A} \cdot \mathbf{G} \cdot \mathbf{B} \cos \theta dA$

• Add the bits:
Faraday's Law





Moving the magnet changes the flux Φ_B (1).
 Changing the current changes the flux Φ_B (2).
 Faraday: changing the flux induces an emf.



Lenz's Law

- Faraday's law gives the direction of the induced emf and therefore the direction of any induced current.
- Lenz's law is a simple way to get the directions straight, with less effort.

Lenz's Law:

The induced emf is directed so that any induced current flow

will *oppose* the *change* in magnetic flux (which causes the induced emf).

• This is easier to use than to say ...

Decreasing magnetic flux \Rightarrow emf creates additional magnetic field





If we move the magnet towards the
 loop
the flux of B will increase.
Lenz's Law ⇒ the current induced in
 the

loop will generate a field B* opposed

Example of Faraday's Law

Consider a coil of radius 5 cm with N = 250 turns. A magnetic field B, passing through it, changes at the rate of dB/dt = 0.6 T/s. The total resistance of the coil is 8 Ω . What is the induced current ?



Use Lenz's law to determine the direction of the induced current.

Apply Faraday's law to find the emf and then the current.



Example of Faraday's Law



Lenz's law: dB/dt > 0

The change in B is increasing the

upward flux through the coil.

So the induced current will have a magnetic field whose flux (and therefore field) is *down*.

Hence the induced current must be *clockwise* when looked at from

above.

Use Faraday's law to get the magnitude of the induced emf and current.

 $\Phi_{\mathbf{B}} = \int \underline{B} \, d\underline{A}$ $\mathcal{E} = - \, d\Phi_{\mathbf{B}} \, / dt$

Induced **B**

The induced EMF is $\mathcal{E} = - d\Phi_B / dt$ In terms of B: $\Phi_B = N(BA) = NB(\pi r^2)$ Therefore $\mathcal{E} = - N(\pi r^2) dB / dt$ $\mathcal{E} = - (250) (\pi 0.005^2)(0.6T/s) = -1.18 V$ $(1V=1Tm^2/s)$ Current I = $\mathcal{E} / R = (-1.18V) / (8 \Omega) = -0.147$ A

Magnetic Flux in a Nonuniform Field

A long, straight wire carries a current I. A rectangular loop (w by I) lies at a distance a, as shown in the figure. What is the magnetic flux through the loop?.

W

W

а

Induced emf Due to Changing Curren

A long, straight wire carries a current $I = I_0 + i t$. A rectangular loop (w by I) lies at a distance a, as shown in the figure. What is the induced emf in the loop?. What is the direction of the induced current and field?

W

Motional EMF

Up until now we have considered fixed loops. The flux through them changed because the magnetic field changed with time.

Now try moving the loop in a uniform and constant magnetic field. This changes the flux, too.



Motional EMF - Use Faraday's Law



This changes in time:

Motional EMF - Use Faraday's Law



This changes in time:

 $d\Phi_B/dt = d(BDx)/dt = BDdx/dt = -BDv$

Hence by Faraday's law there is an induced emf and current. What is the direction of the current?

Motional EMF - Use Faraday's Law



The flux is $\Phi_B = \underline{B} \cdot \underline{A} = BDx$ This changes in time:

 $d\Phi_B/dt = d(BDx)/dt = BDdx/dt = -BDv$ Hence by Faraday's law there is an induced emf and current. What is the direction of the current? **Lenz's law**: there is less inward flux through the loop. Hence the induced current gives inward flux. \Rightarrow So the induced current is clockwise.



Now Faraday's Law $d\Phi_B/dt = -\mathcal{E}$ gives the EMF $\Rightarrow \mathcal{E} = BDv$

In a circuit with a resistor, this gives $\mathcal{E} = BDv = IR \implies I = BDv/R$

Thus moving a circuit in a magnetic field produces an emf exactly like a battery. This is the principle of an electric generator.

Rotating Loop - The Electricity Generator

Consider a loop of area A in a region of space in which there is a uniform magnetic field B. Rotate the loop with an angular frequency ω .

The flux changes because angle θ changes with time: $\theta = \omega t$.

Hence:

 $d\Phi_{B}/dt = d(\underline{B},\underline{A})/dt$ $= d(BA\cos\theta)/dt$ $= B A d(\cos(\omega t))/dt$ $= - BA\omega \sin(\omega t)$



Rotating Loop - The Electricity Genera



$d\Phi_{\rm B}/dt = - BA\omega \sin(\omega t)$

•Then by Faraday's Law this motion causes an emf

$$\mathcal{E} = - d\Phi_{\rm B}/dt = BA\omega \sin(\omega t)$$

This is an <u>AC (alternating current)</u>

A New Source of EMF

- If we have a conducting loop in a magnetic field, we can create an EMF (like a battery) by changing the value of <u>B</u>. <u>A</u>.
- This can be done by changing the area, by changing the magnetic field, or the angle between them.
- We can use this source of EMF in electrical circuits in the same way we used batteries.
- Remember we have to do work to move the loop or to change B, to generate the EMF (Nothing is for free!).

Example: a 120 turn coil (r= 1.8 cm, R = 5.3Ω) is placed outside a solenoid (r=1.6cm, n=220/cm, i=1.5A). The current in the solenoid is reduced to 0 in 0.16s. What current appears in the coil? **Current induced** in coil: $i_{C} = \frac{EMF}{R} = (\frac{N}{R}) \frac{d\Phi_{B}}{dt}$ $\Phi_{R} = \vec{B} \cdot \vec{A} = (\mu_{0} n i_{s}) A_{s}$ Only field in coil is inside solenoid

Example: a 120 turn coil (r = 1.8 cm, R = 5.3Ω) is placed outside a solenoid (r=1.6cm, n=220/cm, i=1.5A). The current in the solenoid is reduced to 0 in 0.16s. What current appears in the coil? **Current induced** in coil: $i_{\mathcal{C}} = \frac{EMF}{R} = (\frac{N}{R}) \frac{d\Phi_{B}}{dt}$ $\Phi_{R} = \vec{B} \cdot \vec{A} = (\mu_{0} n i_{s}) A_{s}$ Only field in coil is inside solenoid $i_{c} = \left(\frac{N}{R}\right) \frac{d(\mu_{0}ni_{s}A_{s})}{dt} = \left(\frac{N}{R}\right) \mu_{0}nA_{s} \frac{di_{s}}{dt}$ $= \left(\frac{N}{R}\right) \mu_0 nA \frac{i_0}{t} = 4.72 mA$

Induced Electric Fields

Consider a stationary conductor Χ B in a time-varying magnetic the electrons must feel a force F. It is not $\underline{F} = \underline{qvxB}$, because the wire is stationary. Instead: we know that $\varepsilon = - d\Phi_{\rm B}/dt$ This is equivalent to an induced electric field E, such that: <u>E</u> = q<u>E</u> and $\varepsilon = \int \underline{E} \cdot \underline{dI}$

a time-varying magnetic field B causes an

Induced Electric Fields $\mathcal{E} = \mathcal{Q}\underline{E}$. $\frac{dI}{\mathcal{E}} = - d\Phi_{\rm B}/dt$ and B X Then: $\mathbf{O} \subseteq \mathbf{E}$. $\mathbf{d} = - \mathbf{d} \Phi_{\mathbf{B}} / \mathbf{d} \mathbf{t}$ **Faraday's Law**

The induced electric field E is NOT a conservative field

We can NOT write E = - dV/dI or $\underline{E} = -\nabla V$

The electrostatic field E_e is conservative $\int E_e dl = 0$



Induced Electric Field

<u>E</u>= q <u>E</u>_e

 $\Delta Vab = -\int \underline{E}_{e} \cdot \underline{dI}$

 $\int \underline{E}_{e} \cdot \underline{dI} = 0$ and $\underline{E}_{e} = \nabla V$

Conservativ

e Work or energy difference does NOT depend on path

Caused by stationary charges or emf sources <u>E</u> = q <u>E</u>

 $\int \underline{E} \cdot \underline{dl} = - d\Phi_{\mathbf{B}}/dt$ $\int \underline{E} \cdot \underline{dl} \neq 0$

Nonconservative

Work or energy difference DOES depend on path

Caused by changing magnetic fields

Induced Electric Fields

$$\oint E d = - d\Phi_B/dt$$

Faraday's Law



B

Now suppose there is no conductor: Is there still an electric field?

YES!, the field does not depend on the presence of the conductor.

For a magnetic field with axial or cylindrical symmetry, the field lines of E are circles.

Self-Inductance

- When the switch is closed, the current does not immediately reach its maximum value
- Faraday's law can be used to describe the effect



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Self-induced emf



- A current in the coil produces a magnetic field directed toward the left (a)
- If the current increases, the increasing flux creates an induced emf of the polarity shown (b)
- The polarity of the induced emf reverses if the current decreases (c)



Inductance of a Solenoid

• The magnetic flux through each turn is $\Phi_{B} = BA = \left(\mu_{o}\frac{N}{\ell}I\right)A$ • Therefore, the inductance is $L = \frac{N\Phi_B}{M} = \frac{\mu_o N^2 A}{M}$ I l • This shows that *L* depends on the geometry of the object

Inductance Units



$$L = \left[\frac{V}{A/s} \right] = \left[\Omega - s \right] = \left[\text{Henry} \right] = \left[H \right]$$





(b)

3. A 2.00-H inductor carries a steady current of 0.500 A. When the switch in the circuit is opened, the current is effectively zero after 10.0 ms. What is the average induced emf in the inductor during this time?

5. A 10.0-mH inductor carries a current $I = I \max \sin \omega t$, with $I \max = 5.00$ A and $\omega/2\pi = 60.0$ Hz. What is the back emf as a function of time?

7. An inductor in the form of a solenoid contains 420 turns, is 16.0 cm in length, and has a cross-sectional area of 3.00 cm2. What uniform rate of decrease of current through the inductor induces an emf of 175 μ N?

LR Circuits



Charging

Kirchhoff Loop Equation:

$$V_{o} - RI - L\frac{dI}{dt} = 0$$

Solution:



R

LR Circuits



Discharging

Kirchhoff Loop Equation:

 $RI + L\frac{dI}{dt} = 0$

Solution:

 $I = I_o e^{-t/\tau}$

R

Active Figure 32.3





(SLIDESHOW MODE ONLY)

14. Calculate the resistance in an *RL* circuit in which L = 2.50 H and the current increases to 90.0% of its final value in 3.00 s.

20. A 12.0-V battery is connected in series with a resistor and an inductor. The circuit has a time constant of 500 μ s, and the maximum current is 200 mA. What is the value of the inductance?

24. A series *RL* circuit with L = 3.00 H and a series *RC* circuit with *C* = 3.00 μ F have equal time constants. If the two circuits contain the same resistance *R*, (a) what is the value of *R* and (b) what is the time constant?

Energy in a coil

$$\mathbf{P} = \mathbf{V}\mathbf{I} = \left(\mathbf{L}\frac{\mathbf{d}\mathbf{I}}{\mathbf{d}t}\right)\mathbf{I}$$

PE in an Inductor

PE in an Capacitor

 $U = \frac{1}{2}LI^{2}$ $U = \frac{1}{2}CV^{2}$



31. An air-core solenoid with 68 turns is 8.00 cm long and has a diameter of 1.20 cm. How much energy is stored in its magnetic field when it carries a current of 0.770 A?

33. On a clear day at a certain location, a 100-V/m vertical electric field exists near the Earth's surface. At the same place, the Earth's magnetic field has a magnitude of 0.500×10^{-4} T. Compute the energy densities of the two fields.

36. A 10.0-V battery, a 5.00- Ω resistor, and a 10.0-H inductor are connected in series. After the current in the circuit has reached its maximum value, calculate (a) the power being supplied by the battery, (b) the power being delivered to the resistor, (c) the power being delivered to the inductor, and (d) the energy stored in the magnetic field of the inductor.

Example 32-5: The Coaxial Cable

- Calculate *L* for the cable
- The total flux is

$$\Phi_B = \int B \, dA = \int_a^b \frac{\mu_o I}{2\pi r} \ell \, dr = \frac{\mu_o I \, \ell}{2\pi} \ln\left(\frac{b}{a}\right)$$

• Therefore, L is

$$L = \frac{\Phi_B}{I} = \frac{\mu_o \ell}{2\pi} \ln\left(\frac{b}{a}\right)$$

The total energy is

$$U = \frac{1}{2} L I^2 = \frac{\mu_o \ell I^2}{4\pi} \ln\left(\frac{b}{a}\right)$$


Mutual Inductance





LC Circuits



Kirchhoff Loop Equation:

$$\frac{Q}{C} - L\frac{dI}{dt} = 0$$





$$Q = Q_{max} \cos(\omega t + \phi)$$
$$\omega = \sqrt{\frac{1}{LC}}$$

I(t=0)=0

 $Q(t=0) = Q_{max}$



Active Figure 32.17





(SLIDESHOW MODE ONLY)

LRC Circuits

Kirchhoff Loop Equation:



Damped *RLC*Circuit

• The maximum value of *Q* decreases after each oscillation

 $-R < R_C$

 This is analogous to the amplitude of a damped spring-mass system



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Active Figure 32.21





(SLIDESHOW MODE ONLY)

LRC Circuits



 $Q = Q_{o}e^{\frac{R}{2L}t} \cos(\omega't + \phi)$

$\omega' = \sqrt{1 + 1}$	1	\mathbb{R}^2
	LC	$\overline{4L^2}$

 $\frac{4L}{C} > R^2$

 $\frac{4L}{2} = R^2$

 $\frac{4L}{-}$ < R²

С

 $\frac{1}{\mathrm{LC}} > \frac{\mathrm{R}^2}{4\mathrm{L}^2}$

 $\frac{1}{\mathrm{LC}} = \frac{\mathrm{R}^2}{4\mathrm{L}^2}$

 \mathbb{R}^2

- Underdamped
- Critically Damped
- Overdamped

Table 32.1

Analogies Between Electrical and Mechanical Systems

Electric Circuit				One-Dimensional Mechanical System
Charge	Q	\leftrightarrow	x	Position
Current	Ι	\leftrightarrow	v_x	Velocity
Potential difference	ΔV	\leftrightarrow	$F_{\mathbf{x}}$	Force
Resistance	R	\leftrightarrow	b	Viscous damping coefficient
Capacitance	C	\leftrightarrow	1/k	(k = spring constant)
Inductance	L	\leftrightarrow	m	Mass
Current = time derivative of charge	$I = \frac{dQ}{dt}$	↔	$v_x = \frac{dx}{dt}$	Velocity = time derivative of position
Rate of change of current = second time derivative of charge	$\frac{dI}{dt} = \frac{d^2Q}{dt^2}$	↔	$a_x = \frac{dv_x}{dt} = \frac{d^2x}{dt^2}$	Acceleration = second time derivative of position
Energy in inductor	$U_L = \frac{1}{2} L I^2$	\leftrightarrow	$K = \frac{1}{2} m v^2$	Kinetic energy of moving object
Energy in capacitor	$U_C = \frac{1}{2} \frac{Q^2}{C}$	\leftrightarrow	$U = \frac{1}{2}kx^2$	Potential energy stored in a spring
Rate of energy loss due to resistance	I^2R	↔	<i>b v</i> ²	Rate of energy loss due to friction
RLC circuit	$L \frac{d^2Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0$	↔	$m \ \frac{d^2x}{dt^2} + b \ \frac{dx}{dt} + kx = 0$	Damped object on a spring

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41. An emf of 96.0 mV is induced in the windings of a coil when the current in a nearby coil is increasing at the rate of 1.20 A/s. What is the mutual inductance of the two coils?

49. A fixed inductance $L = 1.05 \ \mu$ H is used in series with a variable capacitor in the tuning section of a radiotelephone on a ship. What capacitance tunes the circuit to the signal from a transmitter broadcasting at 6.30 MHz?

55. Consider an *LC* circuit in which L = 500 mH and $C = 0.100 \ \mu$ F. (a) What is the resonance frequency ω_0 ? (b) If a resistance of 1.00 k Ω is introduced into this circuit, what is the frequency of the (damped) oscillations? (c) What is the percent difference between the two frequencies?

LC Demo



R = 10 W C = 2.5 mF L = 850 mH

- 1. Calculate period
- 2. What if we change C = 10 mF
- 3. Underdamped?
- 4. How can we change damping?

Reactance and resistance which are in series cannot simply be added because they differ in the phase relationship between current and voltage. With reactance, current and voltage are out of phase ("ELI the ICE man") while with resistance they are in phase. A method called arrow is triangulation is used to combine the two and to determine the resulting inductive phase relationship. reactance

The impedance of the reactance and resistance in series is found by completing the triangle whose legs are R and X:

$$Z = \sqrt{R^2 + X^2}$$

The phase angle between current and voltage (the angle by which the reactance voltage leads the current is given by:

$$\theta = \arctan \frac{X}{R}$$

Make sure that X is given the correct sign (positive or negative) when finding the phase angle.

X

 X_{c}

This

arrow is

θ

inductive reactance

is positive; capacitive

reactance is negative

For a parallel combination of resistance and reactance, we use a different formula for

For a parallel combination of resistance and reactance, we use a different formula for finding the resulting impedance:

$$Z = \frac{RX}{\sqrt{R^2 + X^2}}$$

In this case, the phase angle by which voltage leads current is given by

$$\theta = \arctan \frac{R}{X}$$

If the reactance is capacitive, the phase angle is negative. If the reactance is inductive, the phase angle is positive.

Ohm's Law for Impedance

$$E = IZ$$
 $I = \frac{E}{Z}$ $Z = \frac{E}{I}$

Page 157

AC Circuits



Lecture Outline

Driven Series LCR Circuit: General solution Resonance condition Power considerations **Transformers** Voltage changes Faraday's Law in action gives induced primary current.

Power considerations

Phasors

- R: V in phase with i
- C: V lags i by 90°
- V leads i by 90°
- $V_{R} = Ri_{R} = \varepsilon_{m} \sin\omega t \implies i_{R} = \frac{\varepsilon_{m}}{R} \sin\omega t$ $V_{C} = \frac{Q}{C} = \varepsilon_{m} \sin\omega t \implies i_{C} = \omega C \varepsilon_{m} \cos\omega t$ $V_{L} = L \frac{di_{L}}{dt} = \varepsilon_{m} \sin\omega t \implies i_{L} = -\frac{\varepsilon_{m}}{\omega L} \cos\omega t$
- A phasor is a vector whose magnitude is the maximum value of a quantity (eg V or I) and which rotates counterclockwise in a 2-d plane with angular velocity ω . Recall uniform circular motion:





Series LCR

AC Circuit

Back to the original problem: the loop equation gives:

$$L\frac{d^{2}Q}{dt^{2}} + \frac{Q}{C} + R\frac{dQ}{dt} = \varepsilon_{m}\sin\omega t$$

Assume a solution of the form:

$$i = i_m \sin(\omega t - \phi)$$

Here all unknowns, (i_m, ϕ) , must be found from the loop eqn; the initial conditions have been taken care of by taking the emf to be: $\varepsilon = \varepsilon_m \sin \omega t$.

 To solve this problem graphically, first write down expressions for the voltages across R,C, and L and then plot the appropriate phasor diagram.

Phasors: LCR

- Given: $\varepsilon = \varepsilon_m \sin \omega t$
- Assume:
 - $i = i_m \sin(\omega t \phi) \implies$

$$Q = -\frac{i_{m}}{\omega}\cos(\omega t - \phi)$$
$$\frac{di}{dt} = i_{m}\omega\cos(\omega t - \phi)$$

$$V_{\rm R} = {\rm Ri} = {\rm Ri}_{\rm m} \sin(\omega t - \phi)$$
$$V_{\rm C} = \frac{{\rm Q}}{{\rm C}} = -\frac{1}{\omega {\rm C}} {\rm i}_{\rm m} \cos(\omega t - \phi)$$
$$V_{\rm L} = {\rm L}\frac{{\rm di}}{{\rm dt}} = \omega {\rm Li}_{\rm m} \cos(\omega t - \phi)$$

From these equations, we can draw the phasor diagram to the right.

This picture corresponds to a snapshot at t=0. The projections of these phasors along the vertical axis are the actual values of the voltages at the given time.





The phasor diagram has been relabeled in terms of the reactances defined from:



The unknowns (i_m, ϕ) can now be solved for graphically since the vector sum of the voltages

 $V_L + V_C + V_R$ must sum to the driving emf ε .

Lecture 20, ACT 3

A driven RLC circuit is connected as shown.

- For what frequencies ω of the voltage source is the current through the resistor largest? ω small (b) ω large (c) $\omega = \frac{1}{\sqrt{LC}}$



Conceptual Question

A driven RLC circuit is connected as shown.

- For what frequencies ω of the voltage source is the current through the resistor largest? ω small (b) ω large (c) $\omega = \frac{1}{\sqrt{LC}}$



•This is NOT a series RLC circuit. We cannot blindly apply our techniques for solving the circuit. We must think a little bit.

 However, we can use the frequency dependence of the impedances (reactances) to answer this question.

- The reactance of an inductor = $X_L = \omega L$.
- The reactance of a capacitor = $X_C = 1/(\omega C)$.
- Therefore,
 - \bullet in the low frequency limit, $X_L \rightarrow 0$ and $X_C \rightarrow \infty$.

•Therefore, as $\omega \rightarrow 0$, the current will flow mostly through the inductor; the current through the capacitor approaches 0.

- in the high frequency limit, $X_L \to \infty$ and $X_C \to 0$.

•Therefore, as $\omega \to \infty$, the current will flow mostly through the capacitor, approaching a maximum $i_{max} = \epsilon/R$.

Phasors:LCR



Phasors:Tips

 This phasor diagram was drawn as a snapshot of time t=0 with the voltages being given as the projections along the y-axis.

•Sometimes, in working problems, it is easier to draw the diagram at a time when the current is along the x-axis (when i=0).



"Full Phasor Diagram"

From this diagram, we can also create a triangle which allows us to calculate the impedance Z:

 $|X_L - X_C|$ φ

ε_m

R

Х

" Impedance Triangle"

R

Phasors:LCR

We have found the general solution for the driven LCR circuit:



Lagging & Leading

The phase ϕ between the current and the driving emf depends on the relative magnitudes of the inductive and capacitive reactances.



Conceptual Question

- The series LCR circuit shown is driven by a generator with voltage $\varepsilon = \varepsilon_{m} \sin \omega t$. The time dependence of the current i which flows in the circuit is shown in the plot.
 - How should ω be changed to bring the current and driving voltage into phase?

1A

(a)



(c) impossible

) increase ω (b) decrease ω

• Which of the following phasors represents the current i at t=0?

Lecture 21, ACT 1

- The series LCR circuit shown is driven by a generator with voltage $\varepsilon = \varepsilon_{m} \sin \omega t$. The time dependence of the current i which flows in the circuit is shown in the plot.
 - How should ω be changed to bring the current and driving voltage into phase?



(a) increase ω (b) decrease ω (c) impossible

• From the plot, it is clear that the current is LEADING the applied voltage.

Therefore, the phasor diagram must look like this:

1A



Therefore, $X_C > X_L$

•To bring the current into phase with the applied voltage, we need to increase X_{c} .

• Increasing () will do both!!

Lecture 21, ACT 1



Resonance

• For fixed R,C,L the current i_m will be a maximum at the resonant frequency ω_0 which makes the impedance Z purely resistive.

ie:
$$i_m = \frac{\varepsilon_m}{Z} = \frac{\varepsilon_m}{\sqrt{R^2 + (X_L - X_C)^2}}$$

reaches a maximum when:

$$X_L = X_C$$

the frequency at which this condition is obtained is given from:

$$\omega_{0}L = \frac{1}{\omega_{0}C} \implies \omega_{0} = \frac{1}{\sqrt{LC}}$$

- Note that this resonant frequency is identical to the natural frequency of the LC circuit by itself!
- At this frequency, the current and the driving voltage are in phase! $tan\phi = \frac{X_L X_C}{R} = 0$

Resonance

The current in an LCR circuit depends on the values of the elements and on the driving frequency through the relation

$$i_{m} = \frac{\varepsilon_{m}}{Z} = \frac{\varepsilon_{m}}{\sqrt{R^{2} + (X_{L} - X_{C})^{2}}}$$

$$i_{\rm m} = \frac{\varepsilon_{\rm m}}{R} \frac{1}{\sqrt{1 + \tan^2 \phi}} = \frac{\varepsilon_{\rm m}}{R} \cos \phi$$

Suppose you plot the current versus ω , the source voltage frequency, you would get:



Power in LCR Circuit

The power supplied by the *emf* in a series LCR circuit depends on the frequency ω_0 . It will turn out that the maximum power is supplied at the resonant frequency ω_0 .

• The instantaneous power (for some frequency, ω) delivered at time t is given by:

$$P(t) = \varepsilon(t)i(t) = (\varepsilon_{m} \sin \omega t)(i_{m} \sin(\omega t - \phi))$$

Remember what this stands for

The most useful quantity to consider here is not the instantaneous power but rather the average power delivered in a cycle.

 $\langle P(t) \rangle = \varepsilon_m i_m \langle \sin \omega t \sin(\omega t - \phi) \rangle$

To evaluate the average on the right, we first expand the $sin(\omega t - \phi)$ term.

Power in LCR Circuit Expanding, $\sin \omega t \sin(\omega t - \phi) = \sin \omega t (\sin \omega t \cos \phi - \cos \omega t \sin \phi)$ Taking the averages, $\langle \sin \omega t \cos \omega t \rangle = 0$ sinotcosot (Product of even and odd function = 0) Generally: n $\langle \sin^2 x \rangle = \frac{1}{2\pi} \int_{0}^{2\pi} \sin^2 x dx = \frac{1}{2}$ 2π ωt Putting it all back together again, $\langle P(t) \rangle = \varepsilon_m i_m \left\{ \cos\phi \langle \sin^2 \omega t \rangle - \sin\phi \langle \sin\omega t \cos\omega t \rangle \right\}$ sin²∞t 0 $\langle P(t) \rangle = \frac{1}{2} \varepsilon_m i_m \cos \phi$ -1 2π ωt

Power in LCR Circuit

• This result is often rewritten in terms of rms values:

Power delivered depends on the phase, ϕ , the "power factor"

phase depends on the values of L, C, R, and ω

therefore...



 $\langle P(t) \rangle = \varepsilon_{rms} i_{rms} \cos \phi$

- Power, as well as current, peaks at $\omega = \omega_0$. The *sharpness* of the resonance depends on the values of the components.
- Recall:

$$i_{\rm m} = \frac{\varepsilon_{\rm m}}{R} \frac{1}{\sqrt{1 + \tan^2 \phi}} = \frac{\varepsilon_{\rm m}}{R} \cos \phi$$

• Therefore,

$$\langle P(t) \rangle = \frac{\varepsilon^2_{rms}}{R} \cos^2 \phi$$

We can write this in the following manner (you can do the algebra):

$$\langle P(t) \rangle = \frac{\varepsilon^2 rms}{R} \frac{x^2}{x^2 + Q^2 (x^2 - 1)^2} \qquad x = \frac{\omega}{\omega_o}$$

... introducing the curious factors Q and x...

$$Q = \frac{L\omega_o}{R}$$
The Q factor

A parameter "Q" is often defined to describe the *sharpness* of resonance peaks in both mechanical and electrical oscillating systems. "Q" is defined as

$$Q \equiv 2\pi \frac{U_{max}}{\Delta U}$$

where U_{max} is max energy stored in the system and ΔU is the energy dissipated in one cycle

norio

For RLC circuit,
$$U_{max}$$
 is (e.g.)
And losses only in R, namely

$$U_{max} = \frac{1}{2}LI_{max}^{2}$$

$$\Delta U = \frac{1}{2}I_{max}^{2}RT = \frac{1}{2}I_{max}^{2}R\left(\frac{2\pi}{\omega_{res}}\right)$$
This gives

$$Q = \frac{\omega_{res}L}{R}$$
And for completeness, note

$$X = \frac{\omega}{\omega_{res}}$$

Power in RLC





Conceptual Question 2

Consider the two circuits shown where $C_{II} = 2 C_{I}$. — What is the relation

2A

2B

between the quality factors, Q_I and Q_{II} , of the two (a) $Q_{II} = Q_I$

•What is the relation between P_I and P_{II} , the power delivered by the generator to the circuit when each circuit is operated at its resonant frequency?

a) $P_{II} < P_{I}$ (b) $P_{II} = P_{I}$ (c) $P_{II} > P_{I}$

(c) $Q_{II} > Q_{I}$

Lecture 21, ACT 2

• Consider the two circuits shown where $C_{II} = 2 C_{I}$.

what is the relation
 between the quality factors,
 Q_I and Q_{II}, of the two

2A

• We know the definition of Q:

• At first glance, it looks like Q is independent of C.

• At second glance, we see this cannot be true, since the resonant frequency ω_0 depends on C!

(b) $Q_{II} = Q_{I}$

 $Q \equiv \frac{\omega_0 L}{\omega_0 L}$

$$\omega_0 = \frac{1}{\sqrt{LC}} \quad \blacksquare$$

Doubling C decreases
$$\omega_0$$
 by sqrt(2)!

Doubling C decreases Q by sqrt(2)!

Doubling C increases the width of the resonance!

3

(C) $Q_{II} > Q_{I}$

Lecture 21, ACT 2

• Consider the two circuits shown where $C_{II} = 2 C_{I}$. - What is the relation 2A between the quality factors, Q_{I} and Q_{II} , of the two

2B

•What is the relation between $\mathsf{P}_{\rm I}$ and $\mathsf{P}_{\rm II}$, the power delivered by the generator to the circuit when each circuit is operated at its resonant frequency?

(a) $P_{II} < P_{I}$ (b) $P_{II} = P_{I}$ (c) $P_{II} > P_{I}$

(b) $Q_{II} = Q_{I}$

(c) $Q_{II} > Q_{I}$

At the resonant frequency, the impedance of the circuit is purely resistive.
 Since the resistances in each circuit are the same, the impedances at the resonant frequency for each circuit are equal.

• Therefore, The power delivered by the generator to each circuit is identical.

Power Transmission

How can we transport power from power stations to homes? Why do we use "high tension" lin e

- At home, the AC voltage obtained from outlets in this country is 120V at 60Hz.
- Transmission of power is typically at very high voltages (eg ~500 kV)
- Transformers are used to raise the voltage for transmission and lower the voltage for use. We'll describe these next.

But why?

- Calculate ohmic losses in the transmission lines:
- Define *efficiency* of transmission:

$$\varepsilon \equiv \frac{P_{out}}{P_{in}} = \frac{iV_{in} - i^2R}{iV_{in}} = 1 - \frac{iR}{V_{in}} \left(\frac{V_{in}}{V_{in}}\right) = 1 - \frac{P_{in}R}{V_{in}^2} Make V_{in}$$

– Note for fixed input power and line resistance, the inefficiency $\propto 1/V^2$

Example: Quebec to Montreal $1000 \text{ km} \Rightarrow \text{R} = 220\Omega$ T

With V_{in} = 735kV, ε = 80%. The efficiency goes to zero quickly if V_{in} were lowered!



Transformers



The iron is used to maximize the mutual inductance. We assume that the entire flux produced by each turn of the primary is trapped in the iron.



•Note: "no load" means no current in secondary. The primary current, termed "the magnetizing current" is small!

Ideal Transformers

with a Load

- What happens when we connect a resistive load to the secondary coil?
 - Flux produced by primary coil induces an *emf* in secondary
 - c *emf* in secondary produces current i₂
 - $i_2 = \frac{V_2}{R}$ This current produces a flux in the secondary coil \propto N_2i_2, which opposes the original flux -- Lenz's law
 - **c** This changing flux appears in the primary circuit as well; the sense of it is to reduce the *emf* in the primary...
 - **ç** However, V_1 is a voltage source.
 - **ç** Therefore, there must be an increased current i_1 (supplied by the voltage source) in the primary which produces a flux $\propto N_1 i_1$ which exactly cancels the flux produced by i_2 .







Transformers with a Load

• With a resistive load in the secondary, the primary current is given by:

$$\dot{i}_1 = \frac{N_2}{N_1}\dot{i}_2 = \frac{N_2}{N_1}\frac{V_2}{R} = \frac{V_1}{(N_1/N_2)^2R}$$

Tron

(primary) (secondary)



This is the equivalent resistance seen by the source.

Lecture 21, ACT 3

The primary coil of an ideal transformer is connected to an AC voltage source as shown. There are 50 turns in the primary and 200 turns in the secondary.

If V₁ = 120 V, what is the potential drop across the resistor R ?
(a) 30 V
(b) 120 V

(primary) (secondary) (C) 480 V

iron

If 960 W are dissipated in the R, what is the current in the

3A

3B

(a) 8 A (b) 16 A (c) 32 A

Lecture 21, ACT 3

The primary coil of an ideal transformer is connected to an AC voltage source as shown. There are 50 turns in the primary and 200 turns in the secondary.

3A

If V₁ = 120 V, what is the potential drop across the resistor R ?
(a) 30 V
(b) 120 V

(primary) (secondary (C) 480 V

Tron

The ratio of turns, $(N_2/N_1) = (200/50) = 4$

The ratio of secondary voltage to primary voltage is equal to the ratio of turns, $(V_2/V_1) = (N_2/N_1)$

Therefore, $(V_2/V_1) = 480 V$



The primary coil of an ideal transformer is connected to an AC voltage source as shown. There are 50 turns in the primary and 200 turns in the secondary.

- If $V_1 = 120$ V, what is the potential drop across the resistor

(a) 30 V (b) 120 V The ratio of turns, $(N_2/N_1) = (200/50) = 4$ The ratio $(V_2/V_1) = (N_2/N_1)$. Therefore, $(V_2/V_1) = 480$ V



V₁

Ç If 960 W are dissipated in the resistor R, what is the current in the primary ?

(a) 8 A (b) 16 A

(c) 32 A

Gee, we didn't talk about power yet....

3A

3B

But, let's assume energy is conserved...since it usually is around here

Therefore, 960 W should be produced in the primary

 $P_1 = V_1 I_1$ implies that 960W/120V = 8 A

Transformers with a Load

• To get that last ACT, you had to use a general philosophy -- energy conservation.

An expression for the RMS power flow looks like this:



$$P_{rms} = V_{1rms}i_{1rms} = \frac{N_1}{N_2}V_{2rms}\frac{N_2}{N_1}i_{2rms} = V_{2rms}i_{2rms}$$

Note: This equation simply says that all power delivered by the generator is dissipated in the resistor ! Energy conservation!!

Chap. 4 Circuit Theorems

- Introduction
- Linearity property
- Superposition
- Source transformations
- Thevenin's theorem
- Norton's theorem
- Maximum power transfer



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17 Linearity Dronerty

Homogeneity property (Scaling) $i \rightarrow v = iR$ $ki \rightarrow kv = kiR$

Additivity property $i_1 \rightarrow v_1 = i_1 R$ $i_2 \rightarrow v_2 = i_2 R$ $i_1 + i_2 \rightarrow (i_1 + i_2) R = i_1 R + i_2 R = v_1 + v_2$

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• Linear circuit consist of – linear elements – linear dependent sources – indepenvd²ent sources $p = i^2 R = \frac{1}{R} : nonlinear$

 $v_s = 10V \rightarrow i = 2A$ $v_s = 1V \rightarrow i = 0.2A$ $v_s = 5mV \leftarrow i = 1mA$

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• KVL
$$12i_1 - 4i_2 + v_s = 0$$
 (4.1.1)
 $-4i_1 + 16i_2 - 3v_x - v_s = 0$ (4.1.2)
 $v_x = 2i_1$
(4.1.2) becomes
 $-10i_1 + 16i_2 - v_s = 0$ (4.1.3)

Eqs(4.1.1) and (4.1.3) we get
$$2i_1+12i_2=0 \rightarrow i_1=-6i_2$$

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Eq(4.1.1), we gEetxample 4.1

$$-76i_2 = v_s = 0 \implies i_2 = \frac{v_s}{76}$$

When_{*V_s*} = 12V $I_0 = i_2 = \frac{12}{76}$ A

When $v_s = 24$ V $I_0 = i_2 = \frac{24}{16}$ A Showing that when the source value is doubled, I_0 doubles.

• Assume $I_0 = 1$ A and use linearity to find the actual value of I_0 in the circuit in fig 4.4. 6Ω $2 V_2 I_2 2 \Omega$ $1 V_1$ 3Ω I_4 I_3 I_1 10 $I_{s} = 15 \, \text{A}$ 7Ω 4Ω 5Ω Eastern Mediterranean **Circuit Theorems** 203 University

Example 4.2
If
$$I_0 = 1A$$
, then $v_1 = (3+5)I_0 = 8V$
 $I_1 = v_1/4 = 2A$, $I_2 = I_1 + I_0 = 3A$
 $V_2 = V_1 + 2I_2 = 8 + 6 = 14V$, $I_3 = \frac{V_2}{7} = 2A$
 $I_4 = I_3 + I_2 = 5A \Rightarrow I_S = 5A$
 $I_0 = 1A \rightarrow I_S = 5A$
 $I_0 = 3A \leftarrow I_S = 15A$

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4.3 Superposition

 The superposition principle states that the voltage across (or current through) an element in a linear circuit is the algebraic sum of the voltages across (or currents through) that element due to each independent source acting alone.

Turn off, killed, inactive source:

 independent voltage source: 0 V (short circuit)

- independent current source: 0 A (open circuit) Eastern Mediterranean Circuit Theorems Pependent sources are left intact.

Steps to apply superposition principle:

- 1. Turn off all independent sources except one source. Find the output (voltage or current) due to that active source using nodal or mesh analysis.
- 2. Repeat step 1 for each of the other independent sources.
- 3. Find the total contribution by adding algebraically all the contributions due to the independent courses

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How to turn off independent sources

Turn off voltages sources = short voltage sources; make it equal to zero voltage
Turn off current sources = open current sources; make it equal to zero current

- Superposition involves more work but simpler circuits.
- Superposition is not applicable to the effect on power.

Use the superposition theorem to find in the circuit in Eig 4.6



Since there are Every putes 4?
let
$$V = V_1 + V_2$$

Voltage division to get
 $V = \frac{4}{48} (6) = 2V$
Hence $i_3 = \frac{8}{4+8} (3) = 2A$
 $v_2 = 4i_3 = 8V$
And we find
 $v = v_1 + v_2 = 2 + 8 = 10V$
(a)
 $v = v_1 + v_2 = 2 + 8 = 10V$







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4.5 Source Transformation

 A source transformation is the process of replacing a voltage source V_s in series with a resistor R by a current source i_s in parallel with a resistor R, or vice versa



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And work of the current source positive terminal of voltage source
Impossible source Transformation

ideal voltage source (R = 0)
ideal current source (R=∞)

• Use source transformation to find V_o in the circuit in Fig 4.17.





we use current divi2sion in Fig.4.18(c) to get $i = \frac{1}{2+8}(2) = 0.4A$

and

 $v_o = 8i = 8(0.4) = 3.2V$

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Applying KVL around the loop in Fig 4.21(b) gives $-3+5i+v_x+18=0$

(4.7.1) 1Ω Appling KVL to the loop containing only the 3V voltage sar-Ge₁t₁he₁v_x = 0 stop an div_x yields

Substitutin1g5 $+155i+i30-E0q \rightarrow (47-1-)45wAeobtain$

Alterna_tiv_vel_x y₄i + v_x + 18 = 0 \Rightarrow i = -4.5A

thus $v_x = 3 - i = 7.5 V$

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4.5 Thevenin's Theorem

 Thevenin's theorem states that a linear twoterminal circuit can be replaced by an equivalent circuit consisting of a voltage source V_{Th} in series with a resistor R_{Th} where V_{Th} is the open circuit voltage at the terminals and R_{Th} is the input or equivalent resistance at the terminals when the independent source are turn off.





(a)



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How to Find Thevenin's Voltage

• Equivalent circuit: same voltage-current $Ve_{II}a=bv_{a}at$ the performing invalided by a - b



$$V_{\rm Th} = v_{oo}$$

(a)

How to Find Thevenin's Resistance $R_{\text{Th}} = R_{\text{in}}$:

- input resistance of the dead circuit at a b.
 - a b open circuited
 - Turn off allindependent sources



Circuit Theorems

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CASE 1

- If the network has no dependent sources:
 - Turn off all independent source.
 - $-R_{TH}$: can be obtained via simplification of either parallel or series connection seen from a-b



The Thevenin's resistance may be negative, indicating that the circuit has ability providing power



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 Find the Thevenin's equivalent circuit of the circuit shown in Fig 4.27, to the left of the te gh 4Ω 1Ω a R 12Ω 32 V R_I 2 A



Find V_{th} $V_{\rm Th}$: (1) Mesh analysis $-32+4i_1+12(i_1-i_2)=0$, $i_2=-2A$ $:: i_1 = 0.5 A$ $V_{\rm Th} = 12(i_1 - i_2) = 12(0.5 + 2.0) = 30V$ (2) Alternatively, Nodal Analysis $(32 - V_{\rm Th})/4 + 2 = V_{\rm Th}/12$ $\therefore V_{\text{Th}} = 30 \text{V}$

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(3) Alternatively, sourcetransform



Toget i_L :

$$i_{L} = \frac{V_{\text{Th}}}{R_{\text{Th}} + R_{L}} = \frac{30}{4 + R_{L}}$$

$$R_L = 6 \rightarrow I_L = 30/10 = 3A$$

 $R_L = 16 \rightarrow I_L = 30/20 = 1.5A$
 $R_L = 36 \rightarrow I_L = 30/40 = 0.75A$

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• Find the Thevenin's equivalent of the circuit in Fig. 2Ω 2Ω 2Ω 2Ω

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5 A

Circuit Theorems

 v_x

6Ω

4Ω

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b

0

• Tofind R_{Th} : Fig(a) ependent source case) independent source $\rightarrow 0$ dependent source \rightarrow intact $w_{-1}V = R = \frac{v_o}{v_o} = \frac{1}{1}$

$$v_o = 1 \text{V}, \quad R_{\text{Th}} = \frac{v_o}{i_o} = \frac{1}{i_o}$$

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Example 4.9
Toget
$$V_{Th}$$
 :Fig(b) Mesh analysis
 $i_1 = 5$
 $-2v_x + 2(i_3 - i_2) = 0 \Rightarrow v_x = i_3 - i_2$
 $4(i_2 - i_1) + 2(i_2 - i_1) + 6i_2 = 0 \Rightarrow 12i_2 - 4i_1 - 2i_3 = 0$
But $4(i_1 - i_2) = v_x$
 $\therefore i_2 = 10/3$.
 $V_{Th} = v_{oc} = 6i_2 = 20V$ 20 V

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But
$$i_x = \frac{0 - v_o}{2} = -\frac{v_o}{2}$$

 $i_o = i_x + \frac{v_o}{4} = -\frac{v_o}{2} + \frac{v_o}{4} = -\frac{v_o}{4}$ or $v_o = -4i_o$

Thus
$$R_{\text{Th}} = \frac{v_o}{i_o} = -4\Omega$$
: Supplying power

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4.6 Norton's Theorem

Norton's theorem states that a linear twoterminal circuit can be replaced by equivalent circuit consisting of a current source I_N in parallel with a resistor R_N where I_N is the short-circuit current through the terminals and R_N is the input or equivalent resistance at the terminals when the independent source are turn off.



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• Thevenin and Norton resistances are $R_N = R_{Th}$ a

• Short circuit current from a to b $I_N = i_{sc} = \frac{V_{\text{Th}}}{R_{\text{Th}}}$



Thevenin or Norton equivalent circuit :

- The open circuit voltage V_{oc} across terminals a and b
- The short circuit current *i_{sc}* at terminals *a* and *b*
- The equivalent or input resistance R_{in} at terminals a and brwhere all independent source are turn of f_Nf.= isc

$$R_{Th} = \frac{V_{Th}}{R_{Th}} = R_N$$

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 Find the Norton equivalent circuit of the circ 8 Ω




Example 4.11
Tofind
$$i_N$$
 (Fig.4.40(*b*))
short – circuit terminals *a* and *b*.
Mesh : $i_1 = 2A$, $20i_2 - 4i_1 - i_2 = 0$
 $i_2 = 1A = i_{sc} = I_N$

$$u = i_{sc} = i_{s$$

Alternative method for
$$I_N$$
 and $I_N = \frac{V_{Th}}{R_{Th}}$

 V_{Th} : open – circuit voltage across terminals *a* and *b*

(Fig 4.40(c)):





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Example 4.12

• Using Norton's theorem, find R_N and I_N of the circuit





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4.8 Maximum Power Trandfer





$$\frac{dp}{dR_L} = V_{TH}^2 \left[\frac{(R_{TH} + R_L)^2 - 2R_L(R_{TH} + R_L)}{(R_{TH} + R_L)^4} \right]$$
$$= V_{TH}^2 \left[\frac{(R_{TH} + R_L - 2R_L)}{(R_{TH} + R_L)^3} \right] = 0$$
$$0 = (R_{TH} + R_L - 2R_L) = (R_{TH} - R_L)$$
$$R_L = R_{TH}$$
$$P_{max} = \frac{V_{TH}^2}{4R_{TH}}$$

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Example 4.13

• Find the value of R_L for maximum power transfer in the circuit of Fig. 4.50. Find the maximum power



Example 4.13 $R_{TH} = 2 + 3 + 6 || 2 = 5 + \frac{6 \times 12}{18} = 9\Omega$



